



New analytical wave structures for the (3 + 1)-dimensional Kadomtsev-Petviashvili and the generalized Boussinesq models and their applications



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ABSTRACT

Different types of soliton wave solutions for the (3 + 1)-dimensional Kadomtsev-Petviashvili and the generalized Boussinesq equations are investigated via the solitary wave ansatz method. These solutions are classified into three categories, namely solitary wave, shock wave, and singular wave solutions. The corresponding integrability criteria, termed as constraint conditions, obviously arise from the study. Moreover, the influences of the free parameters and interaction properties in these solutions are discussed graphically for physical interests and possible applications.

Introduction

Traveling wave solutions of the nonlinear evolution equations (NLEEs) are of utmost important through the wave phenomena since they act as a bridge between mathematics and its applications in different branches of science [1–13].

In the last decay, soliton wave solutions and its characteristics have been investigated and applied in many fields, such as ocean engineering [14,15], optical fibers [16,17], materials [18,19], fluid dynamics [20], and so on. This kind of wave solutions has various forms such as solitary waves, shock waves, singular waves, cnoidal waves, snoidal waves, cuspons, and peakons.

The most appropriate way to comprehend the dynamics of the NLEEs is to find their exact solutions [21–32]. Different approaches are used in literature for calculating the exact solutions for the NLEEs. Among these method; the improved fractional sub-equation method [33,34], Kudryashov method and its extended form [35–38], the unified method [39–41] and its generalized scheme [42–47], the homotopy perturbation method [48,49], and the new extended trial equation method [50,51].

The main purpose of this paper is to find the solitary wave (which is sufficiently short in duration and locally irregular given in space disturbances), shock wave (it is a type of propagating disturbance that

moves faster than the other waves in the medium), and singular wave (this is a type of traveling wave solutions has blow up phenomenon) solutions for the (3 + 1)-dimensional Kadomtsev-Petviashvili [52–54] and the generalized Boussinesq [54,55] equations via the solitary wave ansatz method [56,57].

The governing equations are:

The (3 + 1)-dimensional Kadomtsev-Petviashvili equation (3D-KPE) [52–54]

The 3D-KPE was first introduced in 1970 by Boris B. Kadomtsev and Vladimir I. Petviashvili [58]. The 3D-KPE describes the water waves of long wavelength with weakly nonlinear restoring forces, waves in ferromagnetic media, and two-dimensional matter-wave pulses in Bose-Einstein condensates. Due to its significance, it have been studied extensively in the literature [52–54].

$$u_{tx} + \nu_1(uu_x)_x + \nu_2 u_{xxx} + \nu_3 u_{yy} + \nu_4 u_{zz} = 0, \quad (1)$$

where $u = u(x, y, z, t)$ is a real valued function in its arguments and the coefficients $\nu_1 = 6$, $\nu_2 = 1$, and $\nu_3 = \nu_4 = \pm 3$. The coefficients $\nu_3 = \nu_4 = \pm 3$ are used for weak surface tension and strong surface tension, respectively [59–61].

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The generalized Boussinesq equation (GBE) [54,55]

The GBE was used in coastal engineering as a computer model for the simulation of water waves in shallow seas and seaports [62]. Furthermore, the GBE arises in the study of water waves [63], anharmonic lattice waves [64], and dense lattices [65].

$$v_{tt} - \delta_1(v^2)_{xx} - \delta_2 v_{xx} - \delta_3 v_{xxxx} - \delta_4 v_{yy} - \delta_5 v_{zz} = 0, \tag{2}$$

where $v = v(x, y, z, t)$ is an elevation on the free surface of fluid, the coefficients $\delta_1, \delta_2, \delta_3, \delta_4,$ and δ_5 are real constants depend on the depth of fluid and characteristic speed of long waves.

This article has been arranged as follows: in Section “Solution to the 3D-KPE” and “Solution to the GBE”, various soliton solutions for the above two equations have been investigated and the properties for these solutions are described with some figures. In Section “Conclusions”, the conclusions have been drawn.

Solution to the 3D-KPE

In this section, the solitary wave, shock wave and singular wave solutions for (1) are calculated.

Solitary wave solutions

In order to calculate the solitary wave solution, suppose

$$u(x, y, z, t) = \frac{A}{\cosh^{\lambda} \psi} \quad \text{where} \quad \psi = \alpha x + \beta y + \gamma z - vt, \tag{3}$$

where α, β, γ are the inverse widths, A is the amplitude and v is the velocity of the solitary wave, λ is a constant to be determined later. By using (3)

$$\begin{aligned} u_{tx} &= \frac{A\lambda(\lambda+1)\alpha v}{\cosh^{\lambda+2} \psi} - \frac{A\lambda^2 v}{\cosh^{\lambda} \psi}, \\ (uu_x)_x &= \frac{2A^2\lambda^2\alpha^2}{\cosh^{2\lambda} \psi} - \frac{A^2\lambda(2\lambda+1)\alpha^2}{\cosh^{2\lambda+2} \psi}, \\ u_{xxxx} &= \frac{A\lambda^4\alpha^4}{\cosh^{\lambda} \psi} - \frac{2A\lambda(\lambda+1)(\lambda^2+2\lambda+2)\alpha^4}{\cosh^{\lambda+2} \psi} + \frac{A\lambda(\lambda+1)(\lambda+2)(\lambda+3)\alpha^4}{\cosh^{\lambda+4} \psi}, \\ u_{yy} &= \frac{A\lambda^2\beta^2}{\cosh^{\lambda} \psi} - \frac{A\lambda(\lambda+1)\beta^2}{\cosh^{\lambda+2} \psi}, \\ u_{zz} &= \frac{A\lambda^2\gamma^2}{\cosh^{\lambda} \psi} - \frac{A\lambda(\lambda+1)\gamma^2}{\cosh^{\lambda+2} \psi}, \end{aligned}$$

substituting above values into (1)

$$\begin{aligned} &\frac{A\lambda(\lambda+1)\alpha v}{\cosh^{\lambda+2} \psi} - \frac{A\lambda^2 v}{\cosh^{\lambda} \psi} + v_1 \frac{2A^2\lambda^2\alpha^2}{\cosh^{2\lambda} \psi} - v_1 \frac{A^2\lambda(2\lambda+1)\alpha^2}{\cosh^{2\lambda+2} \psi} - v_2 \\ &\frac{2A\lambda(\lambda+1)(\lambda^2+2\lambda+2)\alpha^4}{\cosh^{\lambda+2} \psi} + v_2 \frac{A\lambda(\lambda+1)(\lambda+2)(\lambda+3)\alpha^4}{\cosh^{\lambda+4} \psi} + v_2 \\ &\frac{A\lambda^4\alpha^4}{\cosh^{\lambda} \psi} + v_3 \frac{A\lambda^2\beta^2}{\cosh^{\lambda} \psi} - v_3 \frac{A\lambda(\lambda+1)\beta^2}{\cosh^{\lambda+2} \psi} + v_4 \frac{A\lambda^2\gamma^2}{\cosh^{\lambda} \psi} - v_4 \\ &\frac{A\lambda(\lambda+1)\gamma^2}{\cosh^{\lambda+2} \psi} = 0. \end{aligned}$$

By comparing the powers $2\lambda, \lambda+2$ and $\lambda+4, 2\lambda+2$

$$\begin{aligned} A\lambda(\lambda+1)\alpha v + 2v_1 A^2\lambda^2\alpha^2 - 2v_2 A\lambda(\lambda+1)(\lambda^2+2\lambda+2)\alpha^4 - v_3 A\lambda \\ (\lambda+1)\beta^2 - v_4 A\lambda(\lambda+1)\gamma^2 = 0, \\ -v_1 A^2\lambda(2\lambda+1)\alpha^2 + v_2 A\lambda(\lambda+1)(\lambda+2)(\lambda+3)\alpha^4 = 0, \end{aligned}$$

set $\lambda = 2$

$$A = \frac{12\alpha^2 v_2}{v_1}, \quad v = \frac{4\alpha^4 v_2 + \beta^2 v_3 + \gamma^2 v_4}{\alpha}.$$

Thus

$$u_1(x, y, z, t) = \frac{A}{\cosh^2(\alpha x + \beta y + \gamma z - vt)}. \tag{4}$$

Shock wave solutions

To calculate the shock wave soliton, suppose

$$u(x, y, z, t) = A \tanh^{\lambda} \psi \quad \text{where} \quad \psi = \alpha x + \beta y + \gamma z - vt \quad \text{and} \quad \lambda > 0, \tag{5}$$

from (5)

$$\begin{aligned} u_{tx} &= -A\lambda(\lambda-1)\alpha v \tanh^{\lambda-2} \psi + 2A\lambda^2\alpha v \tanh^{\lambda} \psi \\ &\quad - A\lambda(\lambda+1)\alpha v \tanh^{\lambda+2} \psi, \\ (uu_x)_x &= A^2\lambda(2\lambda-1)\alpha^2 \tanh^{2\lambda-2} \psi - 4A^2\lambda^2\alpha^2 \tanh^{2\lambda} \psi \\ &\quad + A^2\lambda(2\lambda+1)\alpha^2 \tanh^{2\lambda+2} \psi, \\ u_{xxxx} &= A\lambda(\lambda-1)(\lambda-2)(\lambda-3)\alpha^4 \tanh^{\lambda-4} \psi - 4A\lambda(\lambda-1)(\lambda^2-2\lambda \\ &\quad + 2)\alpha^4 \tanh^{\lambda-2} \psi - 4A\lambda(\lambda+1)(\lambda^2+2\lambda+2)\alpha^4 \tanh^{\lambda+2} \psi \\ &\quad + 2A\lambda^2(3\lambda^2+5)\alpha^4 \tanh^{\lambda} \psi + A\lambda(\lambda+1)(\lambda+2)(\lambda \\ &\quad + 3)\alpha^4 \tanh^{\lambda+4} \psi, \\ u_{yy} &= A\lambda(\lambda-1)\beta^2 \tanh^{\lambda-2} \psi - 2A\lambda^2\beta^2 \tanh^{\lambda} \psi + A\lambda(\lambda+1)\beta^2 \tanh^{\lambda+2} \psi, \\ u_{zz} &= A\lambda(\lambda-1)\gamma^2 \tanh^{\lambda-2} \psi - 2A\lambda^2\gamma^2 \tanh^{\lambda} \psi + A\lambda(\lambda+1)\gamma^2 \tanh^{\lambda+2} \psi, \end{aligned}$$

substituting above values into (1)

$$\begin{aligned} &-A\lambda(\lambda-1)\alpha v \tanh^{\lambda-2} \psi + 2A\lambda^2\alpha v \tanh^{\lambda} \psi - A\lambda(\lambda+1) \\ &\quad \alpha v \tanh^{\lambda+2} \psi + v_1 A^2\lambda(2\lambda-1)\alpha^2 \tanh^{2\lambda-2} \psi - 4v_1 A^2\lambda^2\alpha^2 \\ &\quad \tanh^{2\lambda} \psi + v_1 A^2\lambda(2\lambda+1)\alpha^2 \tanh^{2\lambda+2} \psi + v_2 A\lambda(\lambda-1)(\lambda-2)(\lambda-3) \\ &\quad \alpha^4 \tanh^{\lambda-4} \psi - 4v_2 A\lambda(\lambda-1)(\lambda^2-2\lambda+2)\alpha^4 \tanh^{\lambda-2} \psi - 4v_2 A\lambda \\ &\quad (\lambda+1)(\lambda^2+2\lambda+2)\alpha^4 \tanh^{\lambda+2} \psi + 2v_2 A\lambda^2(3\lambda^2+5) \\ &\quad \alpha^4 \tanh^{\lambda} \psi + v_2 A\lambda(\lambda+1)(\lambda+2)(\lambda+3)\alpha^4 \tanh^{\lambda+4} \psi + v_3 A\lambda(\lambda-1) \\ &\quad \beta^2 \tanh^{\lambda-2} \psi - 2v_3 A\lambda^2\beta^2 \tanh^{\lambda} \psi + v_3 A\lambda(\lambda+1)\beta^2 \tanh^{\lambda+2} \psi + v_4 A\lambda \\ &\quad (\lambda-1)\gamma^2 \tanh^{\lambda-2} \psi - 2v_4 A\lambda^2\gamma^2 \tanh^{\lambda} \psi + v_4 A\lambda(\lambda+1)\gamma^2 \tanh^{\lambda+2} \psi \\ &= 0. \end{aligned}$$

By comparing the powers $2\lambda, \lambda+2$ and $\lambda+4, 2\lambda+2$

$$\begin{aligned} &-A\lambda(\lambda+1)\alpha v - 4v_1 A^2\lambda^2\alpha^2 - 4v_2 A\lambda(\lambda+1)(\lambda^2+2\lambda+2)\alpha^4 + v_3 A\lambda \\ &\quad (\lambda+1)\beta^2 + v_4 A\lambda(\lambda+1)\gamma^2 = 0, \\ &v_1 A^2\lambda(2\lambda+1)\alpha^2 + v_2 A\lambda(\lambda+1)(\lambda+2)(\lambda+3)\alpha^4 = 0, \end{aligned}$$

set $\lambda = 2$

$$A = -\frac{12\alpha^2 v_2}{v_1}, \quad v = \frac{-8\alpha^4 v_2 + \beta^2 v_3 + \gamma^2 v_4}{\alpha}.$$

Thus

$$u_2(x, y, z, t) = A \tanh^2(\alpha x + \beta y + \gamma z - vt). \tag{6}$$

Singular wave Form-I

For the singular wave Form I solution, suppose

$$u(x, y, z, t) = A \coth^{\lambda} \psi \quad \text{where} \quad \psi = \alpha x + \beta y + \gamma z - vt \quad \text{and} \quad \lambda > 0, \tag{7}$$

from (7)

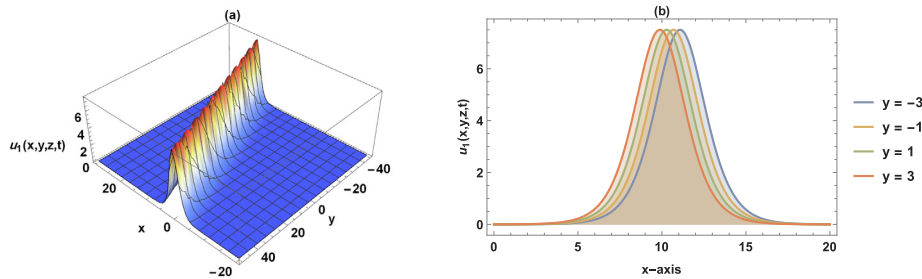


Fig. 1. $u_1(x, y, z, t)$ when $z = 0, t = 2$ in 3D- and 2D-plots.

$$\begin{aligned}
 u_{tx} &= -A\lambda(\lambda - 1)\alpha\nu\coth^{\lambda-2}\psi + 2A\lambda^2\alpha\nu\coth^{\lambda-1}\psi - A\lambda(\lambda + 1)\alpha\nu\coth^{\lambda+2}\psi, \\
 (uu_x)_x &= A^2\lambda(2\lambda - 1)\alpha^2\coth^{2\lambda-2}\psi - 4A^2\lambda^2\alpha^2\coth^{2\lambda}\psi \\
 &\quad + A^2\lambda(2\lambda + 1)\alpha^2\coth^{2\lambda+2}\psi, \\
 u_{xxxx} &= A\lambda(\lambda - 1)(\lambda - 2)(\lambda - 3)\alpha^4\coth^{\lambda-4}\psi - 4A\lambda(\lambda - 1)(\lambda^2 - 2\lambda \\
 &\quad + 2)\alpha^4\coth^{\lambda-2}\psi - 4A\lambda(\lambda + 1)(\lambda^2 + 2\lambda + 2)\alpha^4\coth^{\lambda+2}\psi \\
 &\quad + 2A\lambda^2(3\lambda^2 + 5)\alpha^4\coth^{\lambda}\psi + A\lambda(\lambda + 1)(\lambda + 2)(\lambda \\
 &\quad + 3)\alpha^4\coth^{\lambda+4}\psi, \\
 u_{yy} &= A\lambda(\lambda - 1)\beta^2\coth^{\lambda-2}\psi - 2A\lambda^2\beta^2\coth^{\lambda-1}\psi + A\lambda(\lambda + 1)\beta^2\coth^{\lambda+2}\psi, \\
 u_{zz} &= A\lambda(\lambda - 1)\gamma^2\coth^{\lambda-2}\psi - 2A\lambda^2\gamma^2\coth^{\lambda-1}\psi + A\lambda(\lambda + 1)\gamma^2\coth^{\lambda+2}\psi,
 \end{aligned}$$

substituting above values into (1)

$$\begin{aligned}
 &-A\lambda(\lambda - 1)\alpha\nu\coth^{\lambda-2}\psi + 2A\lambda^2\alpha\nu\coth^{\lambda-1}\psi - A\lambda(\lambda + 1) \\
 &\quad \alpha\nu\coth^{\lambda+2}\psi + \nu_1 A^2\lambda(2\lambda - 1)\alpha^2\coth^{2\lambda-2}\psi - 4\nu_1 A^2\lambda^2\alpha^2 \\
 &\quad \coth^{2\lambda}\psi + \nu_1 A^2\lambda(2\lambda + 1)\alpha^2\coth^{2\lambda+2}\psi + \nu_2 A\lambda(\lambda - 1)(\lambda - 2)(\lambda - 3) \\
 &\quad \alpha^4\coth^{\lambda-4}\psi - 4\nu_2 A\lambda(\lambda - 1)(\lambda^2 - 2\lambda + 2)\alpha^4\coth^{\lambda-2}\psi - 4\nu_2 A\lambda \\
 &\quad (\lambda + 1)(\lambda^2 + 2\lambda + 2)\alpha^4\coth^{\lambda+2}\psi + 2\nu_2 A\lambda^2(3\lambda^2 + 5)\alpha^4\coth^{\lambda}\psi + \nu_2 A \\
 &\quad \lambda(\lambda + 1)(\lambda + 2)(\lambda + 3)\alpha^4\coth^{\lambda+4}\psi + \nu_3 A\lambda(\lambda - 1)\beta^2\coth^{\lambda-2}\psi - 2 \\
 &\quad \nu_3 A\lambda^2\beta^2\coth^{\lambda-1}\psi + \nu_3 A\lambda(\lambda + 1)\beta^2\coth^{\lambda+2}\psi + \nu_4 A\lambda(\lambda - 1) \\
 &\quad \gamma^2\coth^{\lambda-2}\psi - 2\nu_4 A\lambda^2\gamma^2\coth^{\lambda-1}\psi + \nu_4 A\lambda(\lambda + 1)\gamma^2\coth^{\lambda+2}\psi = 0.
 \end{aligned}$$

By comparing the powers $2\lambda, \lambda + 2$ and $\lambda + 4, 2\lambda + 2$

$$\begin{aligned}
 &-A\lambda(\lambda + 1)\alpha\nu - 4\nu_1 A^2\lambda^2\alpha^2 - 4\nu_2 A\lambda(\lambda + 1)(\lambda^2 + 2\lambda + 2)\alpha^4 + \nu_3 A\lambda \\
 &\quad (\lambda + 1)\beta^2 + \nu_4 A\lambda(\lambda + 1)\gamma^2 = 0, \\
 &\nu_1 A^2\lambda(2\lambda + 1)\alpha^2 + \nu_2 A\lambda(\lambda + 1)(\lambda + 2)(\lambda + 3)\alpha^4 = 0,
 \end{aligned}$$

set $\lambda = 2$

$$A = -\frac{12\alpha^2\nu_2}{\nu_1}, \quad \nu = \frac{-8\alpha^4\nu_2 + \beta^2\nu_3 + \gamma^2\nu_4}{\alpha}.$$

Thus

$$u_3(x, y, z, t) = A\coth^2(\alpha x + \beta y + \gamma z - \nu t).$$

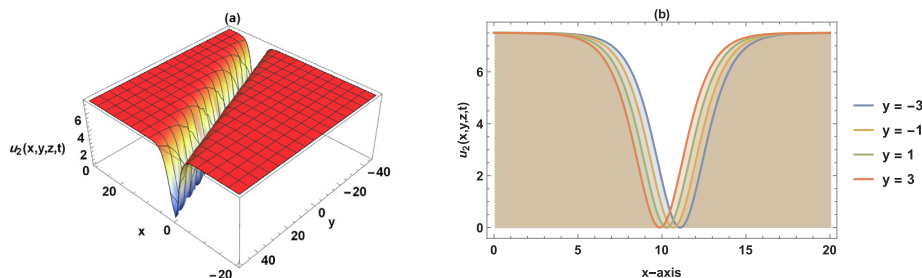


Fig. 2. $u_2(x, y, z, t)$ when $z = 0, t = 2$ in 3D- and 2D-plots.

Singular wave Form-II

For the singular wave Form II solution, suppose

$$u(x, y, z, t) = A\text{csch}^{\lambda}\psi \quad \text{where } \psi = \alpha x + \beta y + \gamma z - \nu t \quad \text{and } \lambda > 0, \tag{9}$$

from (9)

$$\begin{aligned}
 u_{tx} &= -A\lambda(\lambda + 1)\alpha\nu\text{csch}^{\lambda+2}\psi - A\lambda^2\alpha\nu\text{csch}^{\lambda}\psi, \\
 (uu_x)_x &= 2A^2\lambda^2\alpha^2\text{csch}^{2\lambda}\psi + A^2\lambda(2\lambda + 1)\alpha^2\text{csch}^{2\lambda+2}\psi, \\
 u_{xxxx} &= A\lambda^4\alpha^4\text{csch}^{\lambda}\psi + 2A\lambda(\lambda + 1)(\lambda^2 + 2\lambda + 2)\alpha^4\text{csch}^{\lambda+2}\psi \\
 &\quad + A\lambda(\lambda + 1)(\lambda + 2)(\lambda + 3)\alpha^4\text{csch}^{\lambda+4}\psi, \\
 u_{yy} &= A\lambda(\lambda + 1)\beta^2\text{csch}^{\lambda+2}\psi + A\lambda^2\beta^2\text{csch}^{\lambda}\psi, \\
 u_{zz} &= A\lambda(\lambda + 1)\gamma^2\text{csch}^{\lambda+2}\psi + A\lambda^2\gamma^2\text{csch}^{\lambda}\psi,
 \end{aligned}$$

substituting above values into (1)

$$\begin{aligned}
 &-A\lambda(\lambda + 1)\alpha\nu\text{csch}^{\lambda+2}\psi - A\lambda^2\alpha\nu\text{csch}^{\lambda}\psi + 2\nu_1 A^2\lambda^2\alpha^2\text{csch}^{2\lambda}\psi + \nu_1 A^2\lambda \\
 &\quad (2\lambda + 1)\alpha^2\text{csch}^{2\lambda+2}\psi + \nu_2 A\lambda^4\alpha^4\text{csch}^{\lambda}\psi + 2\nu_2 A\lambda(\lambda + 1)(\lambda^2 + 2\lambda + 2) \\
 &\quad \alpha^4\text{csch}^{\lambda+2}\psi + \nu_2 A\lambda(\lambda + 1)(\lambda + 2)(\lambda + 3)\alpha^4\text{csch}^{\lambda+4}\psi + \nu_3 A\lambda(\lambda + 1) \\
 &\quad \beta^2\text{csch}^{\lambda+2}\psi + \nu_3 A\lambda^2\beta^2\text{csch}^{\lambda}\psi + \nu_4 A\lambda(\lambda + 1)\gamma^2\text{csch}^{\lambda+2}\psi + \nu_4 A\lambda^2\gamma^2 \\
 &\quad \text{csch}^{\lambda}\psi = 0.
 \end{aligned}$$

By comparing the powers $2\lambda, \lambda + 2$ and $\lambda + 4, 2\lambda + 2$

$$\begin{aligned}
 &-A\lambda(\lambda + 1)\alpha\nu + 2\nu_1 A^2\lambda^2\alpha^2 + 2\nu_2 A\lambda(\lambda + 1)(\lambda^2 + 2\lambda + 2)\alpha^4 + \nu_3 A\lambda \\
 &\quad (\lambda + 1)\beta^2 + \nu_4 A\lambda(\lambda + 1)\gamma^2 \\
 &= 0. \quad \nu_1 A^2\lambda(2\lambda + 1)\alpha^2 + \nu_2 A\lambda(\lambda + 1)(\lambda + 2)(\lambda + 3)\alpha^4 = 0,
 \end{aligned}$$

set $\lambda = 2$

$$A = -\frac{12\alpha^2\nu_2}{\nu_1}, \quad \nu = \frac{4\alpha^4\nu_2 + \beta^2\nu_3 + \gamma^2\nu_4}{\alpha}.$$

Thus

$$u_4(x, y, z, t) = A\text{csch}^2(\alpha x + \beta y + \gamma z - \nu t). \tag{10}$$

Figs. 1–4 depicts the 3D and 2D charts of the solution given by $u_i(x, y, z, t), i = 1, 2, 3, 4$ respectively with the parameters $\alpha = 0.5, \beta = 0.1, \gamma = 0.1, \nu_1 = 2, \nu_2 = 5,$ and $\nu_3 = \nu_4 = 3$.

Fig. (1)(a) and (b) represent a bright soliton wave which is a stable

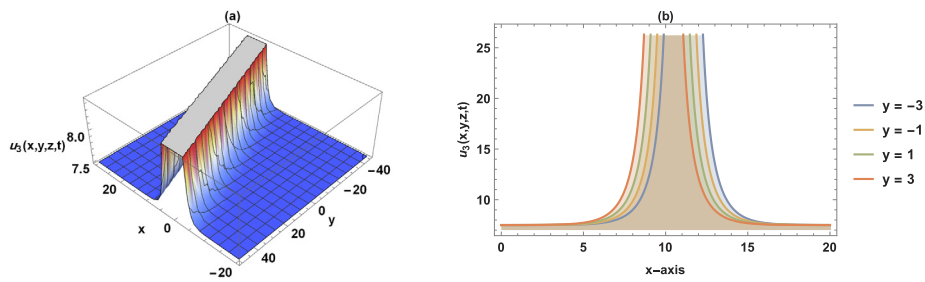


Fig. 3. $u_3(x, y, z, t)$ when $z = 0, t = 2$ in 3D- and 2D-plots.

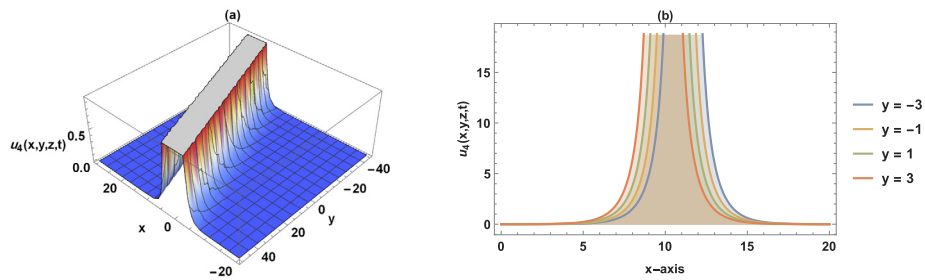


Fig. 4. $u_4(x, y, z, t)$ when $z = 0, t = 2$ in 3D- and 2D-plots.

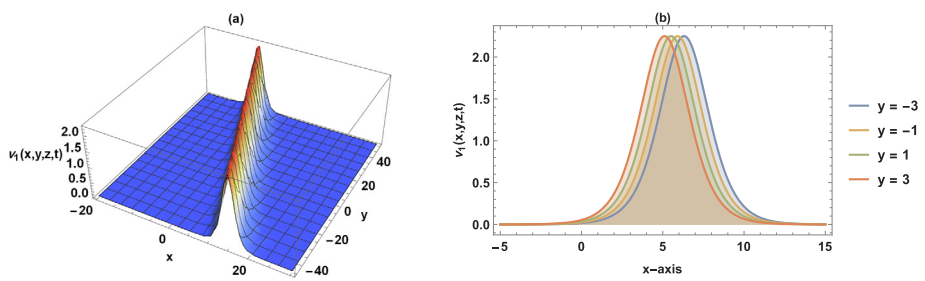


Fig. 5. $v_1(x, y, z, t)$ when $z = 0, t = 2$ in 3D- and 2D-plots.

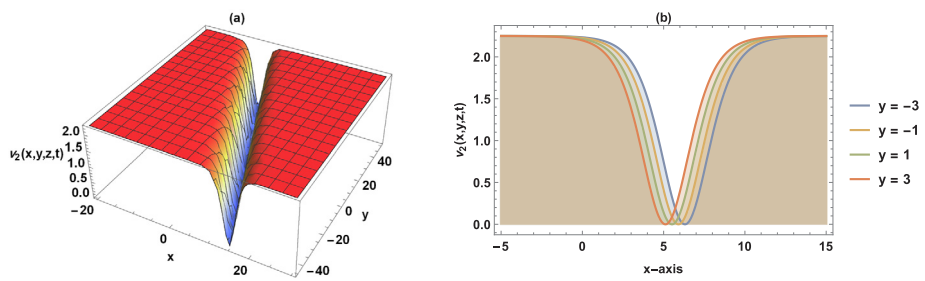


Fig. 6. $v_2(x, y, z, t)$ when $z = 0, t = 2$ in 3D- and 2D-plots.

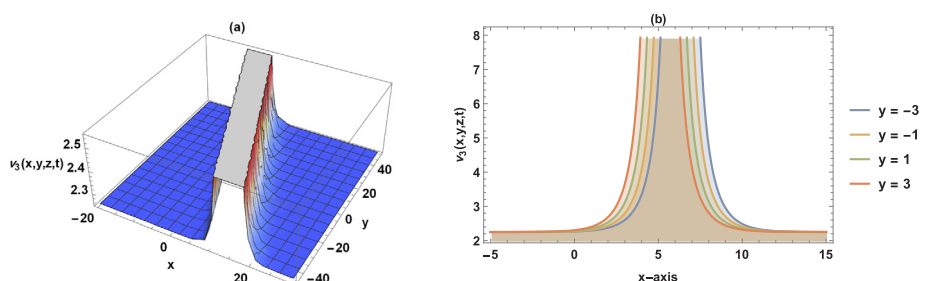


Fig. 7. $v_3(x, y, z, t)$ when $z = 0, t = 2$ in 3D- and 2D-plots.

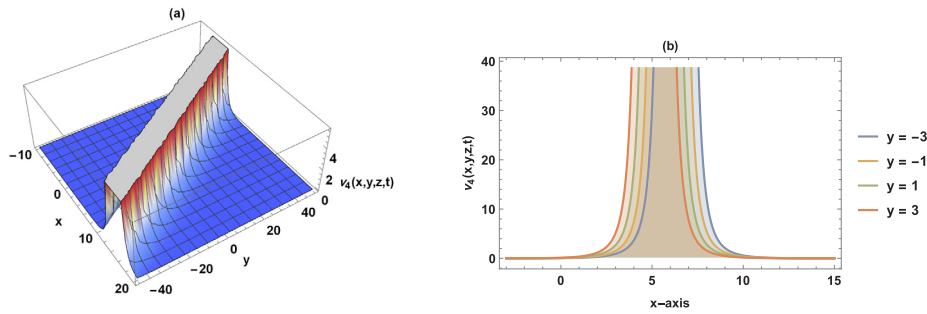


Fig. 8. $v_4(x, y, z, t)$ when $z = 0, t = 2$ in 3D- and 2D-plots.

solution while Fig. (2)(a) and (b) represent a dark soliton wave which is also a stable solution (shock wave solution).

Figs. (3)(a), (b), (4)(a), and (b) represent singular wave solutions which are not stable.

Solution to the GBE

In this section, the solitary wave, shock wave and singular wave solutions for (2) are obtained.

Solitary wave solutions

In order to calculate the solitary wave solution, suppose

$$v(x, y, z, t) = \frac{A}{\cosh^{\lambda} \psi} \quad \text{where} \quad \psi = \alpha x + \beta y + \gamma z - vt, \tag{11}$$

from (11)

$$\begin{aligned} v_{tt} &= \frac{A\lambda^2 v^2}{\cosh^{\lambda} \psi} - \frac{A\lambda(\lambda+1)v^2}{\cosh^{\lambda+2} \psi}, \\ (v^2)_{xx} &= \frac{4A^2\lambda^2\alpha^2}{\cosh^{2\lambda} \psi} - \frac{2A^2\lambda(2\lambda+1)\alpha^2}{\cosh^{2\lambda+2} \psi}, \\ v_{xx} &= \frac{A\lambda^2\alpha^2}{\cosh^{\lambda} \psi} - \frac{A\lambda(\lambda+1)\alpha^2}{\cosh^{\lambda+2} \psi}, \\ v_{xxxx} &= \frac{A\lambda^4\alpha^4}{\cosh^{\lambda} \psi} - \frac{2A\lambda(\lambda+1)(\lambda^2+2\lambda+2)\alpha^4}{\cosh^{\lambda+2} \psi} + \frac{A\lambda(\lambda+1)(\lambda+2)(\lambda+3)\alpha^4}{\cosh^{\lambda+4} \psi}, \\ v_{yy} &= \frac{A\lambda^2\beta^2}{\cosh^{\lambda} \psi} - \frac{A\lambda(\lambda+1)\beta^2}{\cosh^{\lambda+2} \psi}, \\ v_{zz} &= \frac{A\lambda^2\gamma^2}{\cosh^{\lambda} \psi} - \frac{A\lambda(\lambda+1)\gamma^2}{\cosh^{\lambda+2} \psi}, \end{aligned}$$

substituting above values into (2)

$$\begin{aligned} \frac{A\lambda^2 v^2}{\cosh^{\lambda} \psi} - \frac{A\lambda(\lambda+1)v^2}{\cosh^{\lambda+2} \psi} - \delta_1 \frac{4A^2\lambda^2\alpha^2}{\cosh^{2\lambda} \psi} + \delta_1 \frac{2A^2\lambda(2\lambda+1)\alpha^2}{\cosh^{2\lambda+2} \psi} - \delta_2 \frac{A\lambda^2\alpha^2}{\cosh^{\lambda} \psi} + \delta_2 \frac{A\lambda(\lambda+1)\alpha^2}{\cosh^{\lambda+2} \psi} - \delta_3 \frac{A\lambda^4\alpha^4}{\cosh^{\lambda} \psi} + \delta_3 \frac{2A\lambda(\lambda+1)(\lambda^2+2\lambda+2)\alpha^4}{\cosh^{\lambda+2} \psi} - \delta_3 \frac{A\lambda(\lambda+1)(\lambda+2)(\lambda+3)\alpha^4}{\cosh^{\lambda+4} \psi} - \delta_4 \frac{A\lambda^2\beta^2}{\cosh^{\lambda} \psi} + \delta_4 \frac{A\lambda(\lambda+1)\beta^2}{\cosh^{\lambda+2} \psi} - \delta_5 \frac{A\lambda^2\gamma^2}{\cosh^{\lambda} \psi} + \delta_5 \frac{A\lambda(\lambda+1)\gamma^2}{\cosh^{\lambda+2} \psi} = 0. \end{aligned}$$

By comparing the powers $2\lambda, \lambda + 2$ and $\lambda + 4, 2\lambda + 2$

$$\begin{aligned} -A\lambda(\lambda+1)v^2 - 4\delta_1 A^2\lambda^2\alpha^2 + \delta_2 A\lambda(\lambda+1)\alpha^2 + 2\delta_3 A\lambda(\lambda+1)(\lambda^2+2\lambda+2)\alpha^4 + \delta_4 A\lambda(\lambda+1)\beta^2 + \delta_5 A\lambda(\lambda+1)\gamma^2 &= 0, \\ 2\delta_1 A^2\lambda(2\lambda+1)\alpha^2 - \delta_3 A\lambda(\lambda+1)(\lambda+2)(\lambda+3)\alpha^4 &= 0, \end{aligned}$$

set $\lambda = 2$

$$A = \frac{6\alpha^2\delta_3}{\delta_1}, \quad v = \pm \sqrt{\alpha^2\delta_2 + 4\alpha^4\delta_3 + \beta^2\delta_4 + \gamma^2\delta_5}.$$

Thus

$$v_1(x, y, z, t) = \frac{A}{\cosh^2(\alpha x + \beta y + \gamma z - vt)}. \tag{12}$$

and the condition for the existence of the solution is $(\alpha^2\delta_2 + 4\alpha^4\delta_3 + \beta^2\delta_4 + \gamma^2\delta_5) > 0$.

Shock wave solitons

In order to calculate the shock wave soliton, suppose

$$v(x, y, z, t) = A \tanh^{\lambda} \psi \quad \text{where} \quad \psi = \alpha x + \beta y + \gamma z - vt \quad \text{and} \quad \lambda > 0, \tag{13}$$

from (13)

$$\begin{aligned} v_{tt} &= A\lambda(\lambda-1)v^2 \tanh^{\lambda-2} \psi - 2A\lambda^2 v^2 \tanh^{\lambda} \psi + A\lambda(\lambda+1)v^2 \tanh^{\lambda+2} \psi, \\ (v^2)_{xx} &= 2A^2\lambda(2\lambda-1)\alpha^2 \tanh^{2\lambda-2} \psi - 8A^2\lambda^2\alpha^2 \tanh^{2\lambda} \psi + 2A^2\lambda(2\lambda+1)\alpha^2 \tanh^{2\lambda+2} \psi, \\ v_{xx} &= A\lambda(\lambda-1)\alpha^2 \tanh^{\lambda-2} \psi - 2A\lambda^2\alpha^2 \tanh^{\lambda} \psi + A\lambda(\lambda+1)\alpha^2 \tanh^{\lambda+2} \psi, \\ v_{xxxx} &= A\lambda(\lambda-1)(\lambda-2)(\lambda-3)\alpha^4 \tanh^{\lambda-4} \psi - 4\lambda(\lambda-1)(\lambda^2-2\lambda+2)\alpha^4 \tanh^{\lambda-2} \psi - 4A\lambda(\lambda+1)(\lambda^2+2\lambda+2)\alpha^4 \tanh^{\lambda+2} \psi + 2A\lambda^2(3\lambda^2+5)\alpha^4 \tanh^{\lambda} \psi + A\lambda(\lambda+1)(\lambda+2)(\lambda+3)\alpha^4 \tanh^{\lambda+4} \psi, \\ v_{yy} &= A\lambda(\lambda-1)\beta^2 \tanh^{\lambda-2} \psi - 2A\lambda^2\beta^2 \tanh^{\lambda} \psi + A\lambda(\lambda+1)\beta^2 \tanh^{\lambda+2} \psi, \\ v_{zz} &= A\lambda(\lambda-1)\gamma^2 \tanh^{\lambda-2} \psi - 2A\lambda^2\gamma^2 \tanh^{\lambda} \psi + A\lambda(\lambda+1)\gamma^2 \tanh^{\lambda+2} \psi, \end{aligned}$$

substituting above values into (2)

$$\begin{aligned} A\lambda(\lambda-1)v^2 \tanh^{\lambda-2} \psi - 2A\lambda^2 v^2 \tanh^{\lambda} \psi + A\lambda(\lambda+1)v^2 \tanh^{\lambda+2} \psi - 2\delta_1 A^2 \lambda(2\lambda-1)\alpha^2 \tanh^{2\lambda-2} \psi + 8\delta_1 A^2 \lambda^2 \alpha^2 \tanh^{2\lambda} \psi - 2\delta_1 A^2 \lambda(2\lambda+1)\alpha^2 \tanh^{2\lambda+2} \psi - \delta_2 A\lambda(\lambda-1)\alpha^2 \tanh^{\lambda-2} \psi + 2\delta_2 A\lambda^2 \alpha^2 \tanh^{\lambda} \psi - \delta_2 A\lambda(\lambda+1)\alpha^2 \tanh^{\lambda+2} \psi - \delta_3 A\lambda(\lambda-1)(\lambda-2)(\lambda-3)\alpha^4 \tanh^{\lambda-4} \psi + 4\delta_3 \lambda(\lambda-1)(\lambda^2-2\lambda+2)\alpha^4 \tanh^{\lambda-2} \psi + 4\delta_3 A\lambda(\lambda+1)(\lambda^2+2\lambda+2)\alpha^4 \tanh^{\lambda+2} \psi - 2\delta_3 A\lambda^2(3\lambda^2+5)\alpha^4 \tanh^{\lambda} \psi - \delta_3 A\lambda(\lambda+1)(\lambda+2)(\lambda+3)\alpha^4 \tanh^{\lambda+4} \psi - \delta_4 A\lambda(\lambda-1)\beta^2 \tanh^{\lambda-2} \psi + 2\delta_4 A\lambda^2 \beta^2 \tanh^{\lambda} \psi - \delta_4 A\lambda(\lambda+1)\beta^2 \tanh^{\lambda+2} \psi - \delta_5 A\lambda(\lambda-1)\gamma^2 \tanh^{\lambda-2} \psi + 2\delta_5 A\lambda^2 \gamma^2 \tanh^{\lambda} \psi - \delta_5 A\lambda(\lambda+1)\gamma^2 \tanh^{\lambda+2} \psi = 0. \end{aligned}$$

By comparing the powers $2\lambda, \lambda + 2$ and $\lambda + 4, 2\lambda + 2$

$$\begin{aligned} A\lambda(\lambda+1)v^2 + 8\delta_1 A^2\lambda^2\alpha^2 - \delta_2 A\lambda(\lambda+1)\alpha^2 + 4\delta_3 A\lambda(\lambda+1)(\lambda^2+2\lambda+2)\alpha^4 - \delta_4 A\lambda(\lambda+1)\beta^2 - \delta_5 A\lambda(\lambda+1)\gamma^2 &= 0, \\ +2\delta_1 A^2\lambda(2\lambda+1)\alpha^2 + \delta_3 A\lambda(\lambda+1)(\lambda+2)(\lambda+3)\alpha^4 &= 0, \end{aligned}$$

set $\lambda = 2$

$$A = -\frac{6\alpha^2\delta_3}{\delta_1}, \quad v = \pm \sqrt{\alpha^2\delta_2 - 8\alpha^4\delta_3 + \beta^2\delta_4 + \gamma^2\delta_5}.$$

Thus

$$v_2(x, y, z, t) = \text{Atanh}^2(\alpha x + \beta y + \gamma z - \nu t), \tag{14}$$

and the condition for the existence of the solution is $(\alpha^2\delta_2 - 8\alpha^4\delta_3 + \beta^2\delta_4 + \gamma^2\delta_5) > 0$.

Singular wave Form-I

To calculate the singular wave Form I solution, suppose

$$v(x, y, z, t) = \text{Acoth}^\lambda \psi \quad \text{where} \quad \psi = \alpha x + \beta y + \gamma z - \nu t \quad \text{and} \quad \lambda > 0, \tag{15}$$

from (15)

$$\begin{aligned} v_{tt} &= A\lambda(\lambda - 1)\nu^2 \text{coth}^{\lambda-2} \psi - 2A\lambda^2\nu^2 \text{coth}^\lambda \psi + A\lambda(\lambda + 1)\nu^2 \text{coth}^{\lambda+2} \psi, \\ (v^2)_{xx} &= 2A^2\lambda(2\lambda - 1)\alpha^2 \text{coth}^{2\lambda-2} \psi - 8A^2\lambda^2\alpha^2 \text{coth}^{2\lambda} \psi \\ &\quad + 2A^2\lambda(2\lambda + 1)\alpha^2 \text{coth}^{2\lambda+2} \psi, \\ v_{xxx} &= A\lambda(\lambda - 1)\alpha^2 \text{coth}^{\lambda-2} \psi - 2A\lambda^2\alpha^2 \text{coth}^\lambda \psi + A\lambda(\lambda + 1)\alpha^2 \text{coth}^{\lambda+2} \psi, \\ v_{xxxx} &= A\lambda(\lambda - 1)(\lambda - 2)(\lambda - 3)\alpha^4 \text{coth}^{\lambda-4} \psi - 4\lambda(\lambda - 1)(\lambda^2 - 2\lambda \\ &\quad + 2)\alpha^4 \text{coth}^{\lambda-2} \psi - 4A\lambda(\lambda + 1)(\lambda^2 + 2\lambda + 2)\alpha^4 \text{coth}^{\lambda+2} \psi \\ &\quad + 2A\lambda^2(3\lambda^2 + 5)\alpha^4 \text{coth}^\lambda \psi + A\lambda(\lambda + 1)(\lambda + 2)(\lambda \\ &\quad + 3)\alpha^4 \text{coth}^{\lambda+4} \psi, \\ v_{yy} &= A\lambda(\lambda - 1)\beta^2 \text{coth}^{\lambda-2} \psi - 2A\lambda^2\beta^2 \text{coth}^\lambda \psi + A\lambda(\lambda + 1)\beta^2 \text{coth}^{\lambda+2} \psi, \\ v_{zz} &= A\lambda(\lambda - 1)\gamma^2 \text{coth}^{\lambda-2} \psi - 2A\lambda^2\gamma^2 \text{coth}^\lambda \psi + A\lambda(\lambda + 1)\gamma^2 \text{coth}^{\lambda+2} \psi, \end{aligned}$$

substituting above values into (2)

$$\begin{aligned} &A\lambda(\lambda - 1)\nu^2 \text{coth}^{\lambda-2} \psi - 2A\lambda^2\nu^2 \text{coth}^\lambda \psi + A\lambda(\lambda + 1)\nu^2 \text{coth}^{\lambda+2} \psi - 2\delta_1 A^2 \\ &\lambda(2\lambda - 1)\alpha^2 \text{coth}^{2\lambda-2} \psi + 8\delta_1 A^2 \lambda^2 \alpha^2 \text{coth}^{2\lambda} \psi - 2\delta_1 A^2 \lambda(2\lambda + 1) \\ &\alpha^2 \text{coth}^{2\lambda+2} \psi - \delta_2 A\lambda(\lambda - 1)\alpha^2 \text{coth}^{\lambda-2} \psi + 2\delta_2 A\lambda^2 \alpha^2 \text{coth}^\lambda \psi - \delta_2 A\lambda \\ &(\lambda + 1)\alpha^2 \text{coth}^{\lambda+2} \psi - \delta_3 A\lambda(\lambda - 1)(\lambda - 2)(\lambda - 3)\alpha^4 \text{coth}^{\lambda-4} \psi + 4\delta_3 A\lambda \\ &(\lambda - 1)(\lambda^2 - 2\lambda + 2)\alpha^4 \text{coth}^{\lambda-2} \psi + 4\delta_3 A\lambda(\lambda + 1)(\lambda^2 + 2\lambda + 2) \\ &\alpha^4 \text{coth}^{\lambda+2} \psi - 2\delta_3 A\lambda^2(3\lambda^2 + 5)\alpha^4 \text{coth}^\lambda \psi - \delta_3 A\lambda(\lambda + 1)(\lambda + 2) \\ &(\lambda + 3)\alpha^4 \text{coth}^{\lambda+4} \psi - \delta_4 A\lambda(\lambda - 1)\beta^2 \text{coth}^{\lambda-2} \psi + 2\delta_4 A\lambda^2 \beta^2 \\ &\text{coth}^\lambda \psi - \delta_4 A\lambda(\lambda + 1)\beta^2 \text{coth}^{\lambda+2} \psi - \delta_5 A\lambda(\lambda - 1)\gamma^2 \text{coth}^{\lambda-2} \psi + 2\delta_5 A \\ &\lambda^2 \gamma^2 \text{coth}^\lambda \psi - \delta_5 A\lambda(\lambda + 1)\gamma^2 \text{coth}^{\lambda+2} \psi = 0. \end{aligned}$$

By comparing the powers $2\lambda, \lambda + 2$ and $\lambda + 4, 2\lambda + 2$

$$\begin{aligned} &A\lambda(\lambda + 1)\nu^2 + 8\delta_1 A^2 \lambda^2 \alpha^2 - \delta_2 A\lambda(\lambda + 1)\alpha^2 + 4\delta_3 A\lambda(\lambda + 1) \\ &(\lambda^2 + 2\lambda + 2)\alpha^4 - \delta_4 A\lambda(\lambda + 1)\beta^2 - \delta_5 A\lambda(\lambda + 1)\gamma^2 = 0, \\ &+ 2\delta_1 A^2 \lambda(2\lambda + 1)\alpha^2 + \delta_3 A\lambda(\lambda + 1)(\lambda + 2)(\lambda + 3)\alpha^4 = 0, \end{aligned}$$

set $\lambda = 2$

$$A = -\frac{6\alpha^2\delta_3}{\delta_1}, \quad \nu = \pm \sqrt{\alpha^2\delta_2 - 8\alpha^4\delta_3 + \beta^2\delta_4 + \gamma^2\delta_5}$$

Thus

$$v_3(x, y, z, t) = \text{Acoth}^2(\alpha x + \beta y + \gamma z - \nu t). \tag{16}$$

and the condition for the existence of the solution is $(\alpha^2\delta_2 - 8\alpha^4\delta_3 + \beta^2\delta_4 + \gamma^2\delta_5) > 0$.

Singular wave Form-II

To calculate the singular wave Form II solution, suppose

$$v(x, y, z, t) = \text{Acsch}^\lambda \psi \quad \text{where} \quad \psi = \alpha x + \beta y + \gamma z - \nu t \quad \text{and} \quad \lambda > 0, \tag{17}$$

from (17)

$$\begin{aligned} v_{tt} &= A\lambda^2\nu^2 \text{csch}^4 \psi + A\lambda(\lambda + 1)\nu^2 \text{csch}^{\lambda+2} \psi, \\ (v^2)_{xx} &= 4A^2\lambda^2\alpha^2 \text{csch}^{2\lambda} \psi + 2A^2\lambda(2\lambda + 1)\alpha^2 \text{csch}^{2\lambda+2} \psi, \\ v_{xx} &= A\lambda(\lambda + 1)\alpha^2 \text{csch}^{\lambda+2} \psi + A\lambda^2\alpha^2 \text{csch}^\lambda \psi, \\ v_{xxx} &= A\lambda^4\alpha^4 \text{csch}^4 \psi + 2A\lambda(\lambda + 1)(\lambda^2 + 2\lambda + 2)\alpha^4 \text{csch}^{\lambda+2} \psi \\ &\quad + A\lambda(\lambda + 1)(\lambda + 2)(\lambda + 3)\alpha^4 \text{csch}^{\lambda+4} \psi, \\ v_{yy} &= A\lambda(\lambda + 1)\beta^2 \text{csch}^{\lambda+2} \psi + A\lambda^2\beta^2 \text{csch}^\lambda \psi, \\ v_{zz} &= A\lambda(\lambda + 1)\gamma^2 \text{csch}^{\lambda+2} \psi + A\lambda^2\gamma^2 \text{csch}^\lambda \psi, \end{aligned}$$

substituting above values into (2)

$$\begin{aligned} &A\lambda^2\nu^2 \text{csch}^4 \psi + A\lambda(\lambda + 1)\nu^2 \text{csch}^{\lambda+2} \psi - 4\delta_1 A^2 \lambda^2 \alpha^2 \text{csch}^{2\lambda} \psi - 2\delta_1 A^2 \lambda \\ &(2\lambda + 1)\alpha^2 \text{csch}^{2\lambda+2} \psi - \delta_2 A\lambda(\lambda + 1)\alpha^2 \text{csch}^{\lambda+2} \psi - \delta_2 A\lambda^2 \alpha^2 \\ &\text{csch}^\lambda \psi - \delta_3 A\lambda^4 \alpha^4 \text{csch}^4 \psi - 2\delta_3 A\lambda(\lambda + 1)(\lambda^2 + 2\lambda + 2) \\ &\alpha^4 \text{csch}^{\lambda+2} \psi - \delta_3 A\lambda(\lambda + 1)(\lambda + 2)(\lambda + 3)\alpha^4 \text{csch}^{\lambda+4} \psi - \delta_4 A\lambda(\lambda + 1) \\ &\beta^2 \text{csch}^{\lambda+2} \psi - \delta_4 A\lambda^2 \beta^2 \text{csch}^\lambda \psi - \delta_5 A\lambda(\lambda + 1)\gamma^2 \text{csch}^{\lambda+2} \psi - \delta_5 A\lambda^2 \gamma^2 \\ &\text{csch}^\lambda \psi = 0. \end{aligned}$$

By comparing the powers $2\lambda, \lambda + 2$ and $\lambda + 4, 2\lambda + 2$

$$\begin{aligned} &A\lambda(\lambda + 1)\nu^2 - 4\delta_1 A^2 \lambda^2 \alpha^2 - \delta_2 A\lambda(\lambda + 1)\alpha^2 - 2\delta_3 A\lambda(\lambda + 1) \\ &(\lambda^2 + 2\lambda + 2)\alpha^4 - \delta_4 A\lambda(\lambda + 1)\beta^2 - \delta_5 A\lambda(\lambda + 1)\gamma^2 = 0, \\ &2\delta_1 A^2 \lambda(2\lambda + 1)\alpha^2 + \delta_3 A\lambda(\lambda + 1)(\lambda + 2)(\lambda + 3)\alpha^4 = 0, \end{aligned}$$

set $\lambda = 2$

$$A = -\frac{6\alpha^2\delta_3}{\delta_1}, \quad \nu = \pm \sqrt{\alpha^2\delta_2 + 4\alpha^4\delta_3 + \beta^2\delta_4 + \gamma^2\delta_5}.$$

Thus

$$v_4(x, y, z, t) = \text{Acsch}^2(\alpha x + \beta y + \gamma z - \nu t), \tag{18}$$

and the condition for the existence of the solution is $(\alpha^2\delta_2 + 4\alpha^4\delta_3 + \beta^2\delta_4 + \gamma^2\delta_5) > 0$.

Figs. 5–8 depicts the 3D and 2D charts of the solution given by $v_i(x, y, z, t), i = 1, 2, 3, 4$ respectively with the parameters $\alpha = 0.5, \beta = 0.1, \gamma = 0.1, \delta_1 = 2, \delta_2 = 5, \delta_3 = \delta_4 = 3, \text{ and } \delta_5 = 1$.

The same discussion as in Figs. 1–4 can be investigated here.

Conclusions

This paper had investigated the analytical solutions to the (3 + 1)-Dimensional Kadomtsev-Petviashvili and the generalized Boussinesq equations with the help of the solitary wave ansatz method. These solutions included solitary wave, shock wave, and singular wave solutions. The dynamical behavior and the propagation of these solutions are discussed in a graded index waveguide by choosing suitable parameters. To our best of knowledge, the discussion and results in this work, comparing with the other results in literature, are new and had different wave structures. The obtained solutions can be critical to understand attributes of the 3D-KPE and the GBE which are important in different branches of science where these equations are used to describe some physical phenomenon.

References

- [1] Ilhan OA, Manafian J, Shahriari M. Lump wave solutions and the interaction phenomenon for a variable-coefficient Kadomtsev-Petviashvili equation. *Comput Math Appl* 2019. <https://doi.org/10.1016/j.camwa.2019.03.048>.
- [2] Manafian J, Mohammadi-Ivatloo B, Abapour M. Lump-type solutions and interaction phenomenon to the (2 + 1)-dimensional Breaking Soliton equation. *Appl Math Comput* 2019;356:13–41.
- [3] Osman MS. Multi-soliton rational solutions for quantum Zakharov-Kuznetsov equation in quantum magnetoplasmas. *Waves Random Complex* 2016;26(4):434–43.
- [4] Abdel-Gawad HI, Elazab NS, Osman M. Exact solutions of space dependent Korteweg-de Vries equation by the extended unified method. *J Phys Soc Jpn* 2013;82(4):044004.
- [5] Osman MS. New analytical study of water waves described by coupled fractional

- variant Boussinesq equation in fluid dynamics. *Pramana* 2019;93(2):26.
- [6] Dehghan M, Manafian J, Saadatmandi A. Solving nonlinear fractional partial differential equations using the homotopy analysis method. *Numer Meth Part DE* 2010;26(2):448–79.
- [7] Manafian J. On the complex structures of the Biswas-Milovic equation for power, parabolic and dual parabolic law nonlinearities. *Eur Phys J Plus* 2015;130(12):255.
- [8] Manafian J, Lakestani M. Lump-type solutions and interaction phenomenon to the bidirectional Sawada-Kotera equation. *Pramana* 2019;92(3):41.
- [9] Foroutan M, Manafian J, Ranjbaran A. Lump solution and its interaction to (3 + 1)-D potential-YTSE equation. *Nonlinear Dyn* 2018;92(4):2077–92.
- [10] Dehghan M, Manafian J, Saadatmandi A. Analytical treatment of some partial differential equations arising in mathematical physics by using the Exp-function method. *Int J Modern Phys B* 2011;25(22):2965–81.
- [11] Manafian J. Novel solitary wave solutions for the (3 + 1)-dimensional extended Jimbo-Miwa equations. *Comput Math Appl* 2018;76(5):1246–60.
- [12] Tariq KU, Younis M, Rezazadeh H, Rizvi STR, Osman MS. Optical solitons with quadratic-cubic nonlinearity and fractional temporal evolution. *Mod Phys Lett B* 2018;32(26):1850317.
- [13] Osman MS, Lu D, Khater MM. A study of optical wave propagation in the non-autonomous Schrödinger-Hirota equation with power-law nonlinearity. *Results Phys* 2019;13:102157.
- [14] Dysthe K, Krogstad HE, Müller P. Oceanic rogue waves. *Ann Rev Fluid Mech* 2008;40:287–310.
- [15] Farazmand M, Sapsis TP. Reduced-order prediction of rogue waves in two-dimensional deep-water waves. *J Comput Phys* 2017;340:418–34.
- [16] Eberhard M, Savojardo A, Maruta A, Römer RA. Rogue wave generation by inelastic quasi-soliton collisions in optical fibres. *Opt Express* 2017;25(23):28086–99.
- [17] Chen S, Ye Y, Baronio F, Liu Y, Cai XM, Grellu P. Optical Peregrine rogue waves of self-induced transparency in a resonant erbium-doped fiber. *Opt Express* 2017;25(24):29687–98.
- [18] Li ZD, Li QY, Xu TF, He PB. Breathers and rogue waves excited by anti-magnonic spin-transfer torque. *Phys Rev E* 2016;94(4):042220.
- [19] Li BQ, Ma YL. Gaussian rogue waves for a nonlinear variable coefficient Schrödinger system in inhomogeneous optical nanofibers. *J Nanoelectr Optoelectr* 2017;12:1397–401.
- [20] Tasbozan O, Senol M, Kurt A, Özkan O. New solutions of fractional Drinfeld-Sokolov-Wilson system in shallow water waves. *Ocean Eng* 2018;161:62–8.
- [21] Islam MH, Khan K, Akbar MA, Salam MA. Exact traveling wave solutions of modified KdV-Zakharov-Kuznetsov equation and viscous Burgers equation. *SpringerPlus* 2014;3(1):105.
- [22] Mahmud F, Samsuzzoha M, Akbar MA. The generalized Kudryashov method to obtain exact traveling wave solutions of the PHI-four equation and the Fisher equation. *Results Phys* 2017;7:4296–302.
- [23] Khan K, Akbar MA, Koppelaar H. Study of coupled nonlinear partial differential equations for finding exact analytical solutions. *R Soc Open Sci* 2015;2(7):140406.
- [24] Khan K, Akbar MA. Study of analytical method to seek for exact solutions of variant Boussinesq equations. *SpringerPlus* 2014;3(1):324.
- [25] Khan K, Akbar MA. Solving unsteady Korteweg-de Vries equation and its two alternatives. *Math Method Appl Sci* 2016;39(10):2752–60.
- [26] Rahman N, Akbar MA. Traveling waves solutions of nonlinear Klein Gordon equation by extended (G'/G)-expansion method. *Ann Pure Appl Math* 2013;3(1):10–6.
- [27] Hafez MG, Alam MN, Akbar MA. Exact traveling wave solutions to the Klein-Gordon equation using the novel (G'/G)-expansion method. *Results Phys* 2014;4:177–84.
- [28] Akbar MA, Ali NHM. The improved F-expansion method with Riccati equation and its applications in mathematical physics. *Cogent Math. Stat.* 2017;4(1):1282577.
- [29] Miah MM, Ali HS, Akbar MA, Wazwaz AM. Some applications of the (G'/G, 1/G)-expansion method to find new exact solutions of NLEEs. *Eur Phys J Plus* 2017;132(6):252.
- [30] Miah MM, Ali HS, Akbar MA. An investigation of abundant traveling wave solutions of complex nonlinear evolution equations: The perturbed nonlinear Schrödinger equation and the cubic-quintic Ginzburg-Landau equation. *Cogent Math Stat* 2016;3(1):1277506.
- [31] Akbar MA, Ali NHM, Roy R. Closed form solutions of two time fractional nonlinear wave equations. *Results Phys* 2018;9:1031–9.
- [32] Huda MA, Akbar MA, Shanta SS. The new types of wave solutions of the Burger's equation and the Benjamin-Bona-Mahony equation. *J Ocean Eng Sci* 2018;3(1):1–10.
- [33] Sahoo S, Ray SS. Improved fractional sub-equation method for (3 + 1)-dimensional generalized fractional KdV-Zakharov-Kuznetsov equations. *Comput Math Appl* 2015;70(2):158–66.
- [34] Guo S, Mei L, Li Y, Sun Y. The improved fractional sub-equation method and its applications to the space-time fractional differential equations in fluid mechanics. *Phys Lett A* 2012;376(4):407–11.
- [35] Ryabov PN, Sineleshchikov DI, Kochanov MB. Application of the Kudryashov method for finding exact solutions of the high order nonlinear evolution equations. *Appl Math Comput* 2011;218(7):3965–72.
- [36] Foroutan M, Manafian J, Taghipour-Farshi H. Exact solutions for Fitzhugh-Nagumo model of nerve excitation via Kudryashov method. *Opt Quant Electron* 2017;49(11):352.
- [37] Kaplan M, Bekir A, Akbulut A. A generalized Kudryashov method to some nonlinear evolution equations in mathematical physics. *Nonlinear Dyn* 2016;85(4):2843–50.
- [38] Bulut H, Pandir Y, Demiray ST. Exact solutions of time-fractional KdV equations by using generalized Kudryashov method. *Int J Modeling Optimization* 2014;4(4):315.
- [39] Osman MS, Machado JAT, Baleanu D. On nonautonomous complex wave solutions described by the coupled Schrödinger-Boussinesq equation with variable-coefficients. *Opt Quant Electron* 2018;50:73.
- [40] Osman MS. On complex wave solutions governed by the 2D Ginzburg-Landau equation with variable coefficients. *Optik* 2018;156:169–74.
- [41] Abdel-Gawad HI, Tantawy M, Osman MS. Dynamic of DNA's possible impact on its damage. *Math Methods Appl Sci* 2016;39(2):168–76.
- [42] Wazwaz AM, Osman MS. Analyzing the combined multi-waves polynomial solutions in a two-layer-liquid medium. *Comput Math Appl* 2018;76(2):276–83.
- [43] Osman MS. On multi-soliton solutions for the (2 + 1)-dimensional breaking soliton equation with variable coefficients in a graded-index waveguide. *Comput Math Appl* 2018;75(1):1–6.
- [44] Abdel-Gawad HI, Osman M. Exact solutions of the Korteweg-de Vries equation with space and time dependent coefficients by the extended unified method. *Indian J Pure Appl Math* 2014;45(1):1–11.
- [45] Osman MS, Machado JAT. The dynamical behavior of mixed-type soliton solutions described by (2 + 1)-dimensional Bogoyavlensky-Konopelchenko equation with variable coefficients. *J Electromagnet Wave* 2018;32(11):1457–64.
- [46] Osman MS. Multiwave solutions of time-fractional (2 + 1)-dimensional Nizhnik-Novikov-Veselov equations. *Pramana* 2017;88(67):1–9.
- [47] Osman MS, Machado JAT. New nonautonomous combined multi-wave solutions for (2 + 1)-dimensional variable coefficients KdV equation. *Nonlinear Dyn* 2018;93(2):733–40.
- [48] He JH. Homotopy perturbation method: a new nonlinear analytical technique. *Appl Math Comput* 2003;135(1):73–9.
- [49] He JH. Application of homotopy perturbation method to nonlinear wave equations. *Chaos Solitons Fractals* 2005;26(3):695–700.
- [50] Gurefe Y, Misirli E, Sonmezoglu A, Ekici M. Extended trial equation method to generalized nonlinear partial differential equations. *Appl Math Comput* 2003;219(10):5253–60.
- [51] Cheng-Shi LIU. A new trial equation method and its applications. *Commun Theor Phys* 2006;45(3):395–7.
- [52] Ma WX. Comment on the (3 + 1) dimensional Kadomtsev-Petviashvili equations. *Commun Nonlinear Sci Numer Simul* 2011;16(7):2663–6.
- [53] Osman MS. Nonlinear interaction of solitary waves described by multi-rational wave solutions of the (2 + 1)-dimensional Kadomtsev-Petviashvili equation with variable coefficients. *Nonlinear Dyn* 2017;87(2):1209–16.
- [54] Chen Y, Yan Z, Zhang H. New explicit solitary wave solutions for (2 + 1)-dimensional Boussinesq equation and (3 + 1)-dimensional KP equation. *Phys Lett A* 2003;307:107–13.
- [55] Gai XL, Geo YT, Yu X, Sun ZY. Soliton Interactions for the Generalized (3 + 1)-Dimensional Boussinesq Equation. *Int. J. Mod. Phys. B* 2012;26(7):1250062.
- [56] Biswas A. 1-soliton solution of the K(m, n) equation with generalized evolution. *Phys Lett A* 2008;372:4601–2.
- [57] Triki H, Wazwaz AM. Bright and dark soliton solutions for a K(m, n) equation with t-dependent coefficients. *Phys Lett A* 2009;373:2162–5.
- [58] Kadomtsev BB, Petviashvili VI. On the stability of solitary waves in weakly dispersive media. *Sov Phys Dokl* 1970;15:539–41.
- [59] Ablowitz MJ, Clarkson PA. *Solitons, nonlinear evolution equations and inverse scattering*. Cambridge: Cambridge University Press; 1991.
- [60] Hirota R. *The direct method in soliton theory*. Cambridge: Cambridge University Press; 2004.
- [61] Wazwaz AM. Multiple-soliton solutions for the KP equation by Hirota's bilinear method and by the tanh-coth method. *Appl Math Comput* 2007;190(1):633–40.
- [62] Hamm L, Madsen PA, Peregrine DH. Wave transformation in the nearshore zone: a review. *Coast Eng* 1993;21(1–3):5–39.
- [63] Makhankov VG. Dynamics of classical solitons (in non-integrable systems). *Phys Rep* 1987;35:1–128.
- [64] Smith SJ, Chatterjee R. The generalized Korteweg-de Vries equation for interacting phonons in a coherent state basis. *Phys Lett A* 1987;125:129–33.
- [65] Rosenau P. Dynamics of dense lattices. *Phys Rev B* 1987;36:5868–76.