

# Importance of multiple slips on bioconvection flow of cross nanofluid past a wedge with gyrotactic motile microorganisms

Ali Saleh Alshomrani<sup>a,\*</sup>, Malik Zaka Ullah<sup>a</sup>, Dumitru Baleanu<sup>b,c</sup>

<sup>a</sup> Department of Mathematics, Faculty of Science, King Abdulaziz University, Jeddah, Saudi Arabia

<sup>b</sup> Department of Mathematics, Cankaya University, Turkey

<sup>c</sup> Institute of Space Sciences, Magurele-Bucharest, Romania



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## ABSTRACT

In the current article, a mathematical model is developed to visualize the flow of non-Newtonian magneto cross nanofluid with mass and heat transport rates having activation energy, motile microorganisms and bioconvection over the wedge. The phenomena of microorganisms is implemented to control the suspension of nanomaterials. The results of hydromagnetic are also integrated into the momentum expression. Nanofluid is developed by dispersing the nanosized particles in the regular fluid. Nanosized solid materials like carbides, ceramics, graphene, metal, alloyed CNTs etc. have been utilized for the preparation of nanofluid. Physically regular fluids have low thermal efficiency. Therefore, the nanosize particles can be utilized to enhance the thermal efficiency of the regular fluids. Nanofluids have many features in hybrid power engine, heat transfer and can be useful in cancer therapy and medicine. The constructed system is first simplified into nonlinear form by introducing similarity variables. Then obtained ordinary differential equations (ODEs) which are evaluated for numerical solution. Further, for numerical approximation, the popular bvp4c scheme built-in function in MATLAB is utilized. Reliable outcomes are achieved for the temperature, velocity, concentration and motile microorganism density profiles. Results for numerous essential flow parameters are shown via numerical outcomes and graphs. It is revealed that velocity upsurges with enhancement in mixed convection parameter while reduces for bioconvection Rayleigh and buoyancy ratio parameters. Furthermore, the volumetric concentration of nanoparticles boost up for growing estimations of activation energy parameter. The microorganisms field upsurges with larger microorganism slip parameters while reduces with the augmentation in magnitude of bioconvection Lewis number and Peclet number. The obtained numerical results are compared with the available data and found good agreement.

## 1. Introduction

Nanofluids are the suspension of nanomaterials in liquid that produce a significant change in their characteristics at modest concentration of nanomaterials. Currently, many researchers are working on nanofluids to modify their behavior so that it can be useful in different fields of life, where thermal efficiency improvement is prominent. Moreover, nanofluids are commonly used in nuclear power plants, transportation, telecommunications, electronics, bioengineering, agricultural engineering and nutrition. The

\* Corresponding author.

E-mail address: [aszaalshomrani@kau.edu.sa](mailto:aszaalshomrani@kau.edu.sa) (A.S. Alshomrani).

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## Nomenclature

$u, v$	Components of velocity [ $m.s^{-1}$ ]
$x, y$	Space coordinates [ $m$ ]
$Ha$	Hartmann number
$B_0$	Magnetic field strength [ $N.m^{-1}.A^{-1}$ ]
$U_0$	Reference velocity [ $m.s^{-1}$ ]
$u_w(x)$	Stretching velocity [ $m.s^{-1}$ ]
$u_e(x)$	Free stream velocity [ $m.s^{-1}$ ]
$(\rho c)_p$	Nanoparticle specific heat [ $J.kg^{-3}.K^{-1}$ ]
$(\rho c)_f$	Heat capacity of fluid [ $J.kg^{-3}.K^{-1}$ ]
$u_{slip}$	Velocity slip [ $m.s^{-1}$ ]
$T_{slip}$	Thermal slip [ $K$ ]
$C_{slip}$	Concentration slip
$N_{slip}$	Microorganism slip
$Q(x)$	heat generation/absorption coefficient
Pr	Prandtl number
$L, P, Q_1, R$	Velocity slip factor [s/m], temperature, concentration and microorganism slip factors [ $m$ ]
$Rc$	Bio-convection Rayleigh parameter
$Rb$	Buoyancy ratio number
$Nb$	Brownian motion parameter
$S$	Mixed convection parameter
$Nt$	Thermophoresis parameter
$E_a$	Activation energy
$Le$	Lewis number
$We$	Weissenberg number
$Pe$	Peclet number
$Lb$	Bioconvection Lewis number
$B_1$	Thermal slip parameter
$B_2$	Solutal slip parameter
$B_3$	Microorganism slip parameter
$Re_x$	Reynolds number
$Nu_x$	Local Nusselt number
$Sh_x$	Local Sherwood number
$Sh_n$	Local microorganism density number
$N_\infty$	Ambient fluid microorganisms
$T_\infty$	Ambient fluid temperature [ $K$ ]
$C_\infty$	Ambient fluid concentration
$D_B$	Brownian diffusion coefficient [ $m^2.s^{-1}$ ]
$D_m$	Microorganisms diffusion coefficient [ $m^2.s^{-1}$ ]
$D_T$	Thermophoresis diffusion coefficient [ $m^2.s^{-1}$ ]
$k_\infty, D_\infty$	Ambient thermal conductivity
$b$	Chemotaxis constant [ $m$ ]
$m$	Fitted rate constant
$W_c$	Cell swimming speed [ $m.s^{-1}$ ]
$C_w$	Concentration at the surface
$N_w$	Microorganisms at the surface
$T_w$	Temperature at the surface [ $K$ ]
$f$	Dimensionless velocity
$N$	Microorganism field
$Kr$	Chemical reaction constant
<i>Greek symbols</i>	
$\omega_1$	Thermal conductivity parameter
$\omega_2$	Concentration conductivity parameter
$\omega_3$	Microorganism conductivity parameter
$\alpha_1$	Velocity slip parameter
$\sigma$	Electrical conductivity [ $S.m^{-1}$ ]

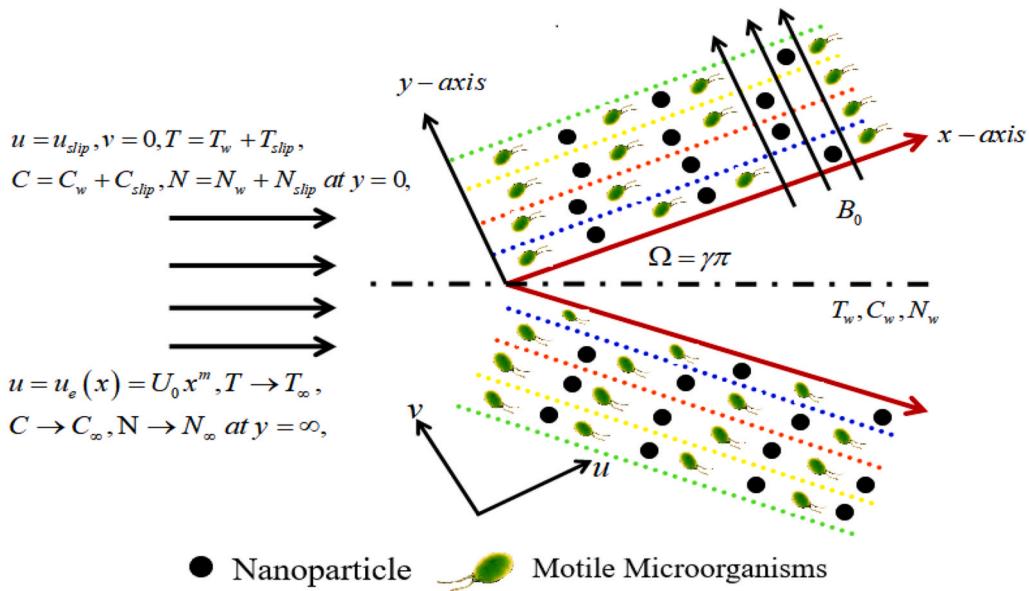
$\sigma^*$	Chemical reaction parameter
$\alpha^*$	Thermal diffusivity [ $m^2.s^{-1}$ ]; $\omega$ Temperature difference parameter
$\delta_1$	Microorganism concentration difference parameter
$\delta_0$	Heat generation parameter
$\tau$	Heat capacity ratio
$q_w$	Thermal mass flux
$q_m$	Solutal mass flux
$q_n$	Microorganism mass flux
$\varphi$	Volumetric concentration
$\theta$	Temperature distribution
$\chi$	Microorganism concentration
$g$	Gravity [ $m.s^{-2}$ ]
$\rho_f$	Fluid density [ $kg.m^{-3}$ ]
$\rho_m$	Microorganisms density [ $kg.m^{-3}$ ]
$\rho_p$	Nanoparticles density [ $kg.m^{-3}$ ]
$\nu$	Kinematic viscosity [ $m^2.s^{-1}$ ]
$\psi(x,y)$	Stream function
$\zeta$	Dimensionless variable
$\gamma$	Wedge angle parameter
$\Omega$	Angle of the wedge
$\Lambda$	Cross time constant

current article is proposed to examine the behavior of nanofluids in different fields and also discuss the enhanced heat transfer characteristics of nanofluids. Initially, it was investigated in the Argonne National Laboratory (ANL) of the USA in 1905 by Choi and Eastman [1]. Such fluids offer the possibility of stirring because of their enhanced convection as compare to regular fluids. The main objective of nanofluids is the higher thermal conductivity potential according to base fluids and higher stability according to micro-particles or nano-particles. Buongiorno [2] has discussed the Brownian motion and thermophoresis features of nanofluid. Raju et al. [3] studied improvement of convection by using nanoliquid in the case of heat absorption/generation. Sheikholeslami and Ganji [4] analyzed the method of heat transfer by utilizing nanofluid. Sheikholeslami and Bhatti [5] described the influences of fluid objects forms in the background of a strong gravitational field. Hsiao [6] examined the magnetohydrodynamic Carreau nanoliquid flow and suggested that investigated findings are beneficial for improvement of different thermal conductivity processes. Turkyilmazoglu [7] inspected the nanoliquid flow via Buongiorno's relation in asymmetric tube. Khan and Shehzad [8] investigated Brownian motion and thermophoresis characteristics of nanomaterials flow over a moving configuration. Waqas et al. [9] used multiple slip characteristics in the flow of nanomaterials across a rotating disk in the appearance of microorganisms. Hassan et al. [10] assessed the characteristics of the transfer of heat over the wave layer through the nanofluid flow. Many researchers admit the significance of nanofluid and have suggested ultimate results [11–25].

Raju et al. [26] studied aspects of the cross-diffusion consequences on the MHD flow of Carreau fluid across a moving sheet with non-uniform temperature source/sink. Sultan et al. [27] explored the numerical treatment of mixed convective Carreau fluid flow with multiple wedge slips. Mahanthesh et al. [28] studied the efficacy of Hall's current and incremental source of heat in the non-linear heat transfer rate of dusty TiO<sub>2</sub>-EO nanofluids with nonlinear radiative effects. Mahanthesh et al. [29] also investigated the effects of MHD flow of SWCNT and MWCNT nanofluids on such a rotating stretching disk with the thermal and exponential space-dependent heat source.

The activation energy initially proposed by Svante Arrhenius in 1889. Activation energy is the basic energy given by the reagents for conversion into materials for different chemical responses. Potential and kinetic energies related with molecules are useful to snapping of bonds or to the stretching and twisting of bonds. Maleque [30] examined the consequences of endothermic/exothermic chemical processes on flat disk subject to activation energy and permeable medium. Awad et al. [31] conducted another relative analysis to discuss activation energy impact utilizing Newtonian model. More work on activation energy can be seen via attempts [32, 33].

The macroscopic flow of convective liquid particles resulting from a transition in the presence of a density gradient is termed as bioconvection. Macroscopic motion is caused by floating microorganisms which changed the composition of the nanofluid. Bioconvection has a wide range of applications in the natural sciences, biotechnology, microsystems, bioinformatics fields, nanomaterials and microfluidic techniques. Bioconvection also plays a major role in engineering in which the electromagnetic field is utilized to set up a bioconvection process for the manufacturing of mechanical energy and power resources. A further important characteristic of bioconvection is the aggregation of nanotechnology with motile microorganisms that enhances the stability, heat and mass transport of nanomaterials. In addition, to obtain substantial bio-convection, microbial thermal transfer is managed by a variety of carbon, light, biochemical inputs, inertia and gravitational rotational velocity that are often difficult to achieve in nanofluids. Virtually, this microbial movement in liquids can be associated with the purification of microorganism's cells and differentiation of various strains. Kuznetsov [34] addressed the bioconvection of nanomaterials across the convectively heated plate subject to gyrotactic micro-organisms. Kuznetsov and Avramenko [35] discussed the idea of bioconvection with the suspension of gyrotactic



**Fig. 1.** A schematic of the wedge in a magneto-cross nanofluid.

microorganisms. Khan et al. [36] observed influences of gyrotactic micro-organisms in magneto Burgers nanoliquid flow. Bio-convective nanofluids flow subject to different thermo-physical characteristics is achieved by Begum et al. [37]. Waqas et al. [38] examined heterogeneous movement of nanoparticles in presence of gyrotactic micro-organisms. Zohra et al. [39] investigated bio-convection slip flow of viscous liquid through a spinning rotating cone. Uddin et al. [40] developed a well-known numerical approach namely the Chebyshev collocation procedure for the analysis of the rheological existence of the Newtonian fluid in the presence of sliding microorganisms. Naz et al. [41] analyzed the dynamics of magnetohydrodynamic (MHD) cross nanofluid with gyrotactic motile microorganisms, bio-convection and entropy generation. Mohamed et al. [42] examined the experimental analysis of micro-bioconvective flow utilizing the Adomian process of decomposition. Ansari et al. [43] analyzed the magnetohydrodynamic bioconvective Casson nanofluid flow via finite element analysis through paired quasilinearisation. Hassan et al. [44] investigated the bioconvection flow of modified second-grade nanofluids subject to nanotubes and gyrotactic motile microorganisms. Wang et al. [45] scrutinized effective Prandtl impacts on bio-convective thermally proposed magnetized tangent hyperbolic nanofluids with micro-organisms and second-order velocity slip. Khan et al. [46] analyzed the bioconvection of couple-stress nanoliquids subject to activation energy and Wu's slip. More work on this topic is seen via studies [47–49].

The main survey in this article examines the impacts of numerous slips in the cross-magneto-nanofluid bioconvection flow past the wedge. Embedded ordinary differential equations are resolved mathematically by using built-in solver bvp4c. Furthermore, Lobatto-IIIa formula is implemented with tolerance  $10^{-6}$ . The outcomes of interesting variables against the flow field are discussed and elaborated through graphical and tabular data. The presented work plays a major role in engineering in which the electromagnetic field is utilized to set up a bioconvection process for the manufacturing of mechanical energy and power resources. A further important characteristic of bioconvection is the aggregation of nanotechnology with motile microorganisms that enhances the stability, heat and mass transport of nanomaterials. In addition, to obtain substantial bio-convection, microbial thermal transfer is managed by a variety of carbon, light, biochemical inputs, inertia and gravitational rotational velocity that are often difficult to achieve in nanofluids.

## 2. Mathematical formulation of the flow

Two-dimensional MHD cross nanofluid flow across a wedge geometry has been formulated. The significance of bioconvection, multiple slips, activation energy, velocity, thermal distribution, the concentration of nanoparticles and the rescaled density of motile microorganisms are also scrutinized in the presence of nano-shaped particles. Here coordinates system is chosen as the x-axis is along with the flow geometry and magnetic strength is parallel to the y-axis direction. Geometry of the flow model is captured through Fig. 1.

The governing equations related to the above assumptions are given below [17,27,50,51]:

### 2.1. Equation of continuity

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0 \quad (1)$$

## 2.2. Equation of momentum

$$u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = u_e \frac{du_e}{dx} + v \frac{\partial}{\partial y} \left[ \frac{\frac{\partial u}{\partial y}}{1 + \left( \Lambda \frac{\partial u}{\partial y} \right)^n} \right] - \frac{\sigma B_0^2}{(\rho c)_f} (u - u_e) \\ + \frac{1}{\rho_f} [(1 - C_f) \rho_f \beta^* g (T - T_\infty) - (\rho_p - \rho_f) g (C - C_\infty) - (N - N_\infty) g \gamma^{**} (\rho_m - \rho_f)], \quad (2)$$

where  $u_e$  depicts free stream velocity,  $\rho_f$  density of fluid,  $n$  power-law index,  $\Lambda$  cross-time constant,  $(\rho c)_f$  heat capacity of fluid,  $B_0$  magnetic strength,  $\sigma$  electric conductivity,  $g$  gravitational acceleration,  $\rho_p$  density of the fluid,  $\rho_m$  density of motile microorganisms and  $\nu$  kinematic viscosity.

## 2.3. Equation of energy

$$u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} = \frac{1}{(\rho c)_f} \frac{\partial}{\partial y} \left[ k(T) \frac{\partial T}{\partial y} \right] + \tau \left\{ D_B \frac{\partial C}{\partial y} \frac{\partial T}{\partial y} + \frac{D_T}{T_\infty} \left( \frac{\partial T}{\partial y} \right)^2 \right\} + \frac{Q}{\rho c_p} (T - T_\infty), \quad (3)$$

where  $k(T)$  is mass diffusivity for cross fluid which can be defined as [66]:

$$k(T) = k_\infty \left[ 1 + \epsilon_1 \frac{T - T_\infty}{T_w - T_\infty} \right], \quad (4)$$

## 2.4. Equation of concentration

$$u \frac{\partial C}{\partial x} + v \frac{\partial C}{\partial y} = \frac{\partial}{\partial y} \left[ D(C) \frac{\partial C}{\partial y} \right] + \frac{D_T}{T_\infty} \frac{\partial^2 T}{\partial y^2} - K r^2 (C - C_\infty) \left( \frac{T}{T_\infty} \right)^m \exp \left( \frac{-E_a}{\kappa T} \right), \quad (5)$$

where  $Kr$  depicts constant chemical reaction,  $D(C)$  mass diffusivity,  $m$  fitted rate constant, and  $E_a$  activation energy. The mass diffusivity is addressed as [66]:

$$D(C) = D_\infty \left[ 1 + \epsilon_2 \frac{C - C_\infty}{C_w - C_\infty} \right], \quad (6)$$

## 2.5. Equation of motile microorganisms

$$u \frac{\partial N}{\partial x} + v \frac{\partial N}{\partial y} = D_m \left( \frac{\partial^2 N}{\partial y^2} \right) - \frac{b W_c}{(C_w - C_\infty)} \left[ N \frac{\partial^2 C}{\partial y^2} + \frac{\partial C}{\partial y} \frac{\partial N}{\partial y} \right], \quad (7)$$

where  $D_m$  shows coefficient of microorganism's diffusion,  $b$  chemotaxis constant and  $W_c$  cell swimming speed.

## 2.6. Boundary conditions

The boundary restrictions are

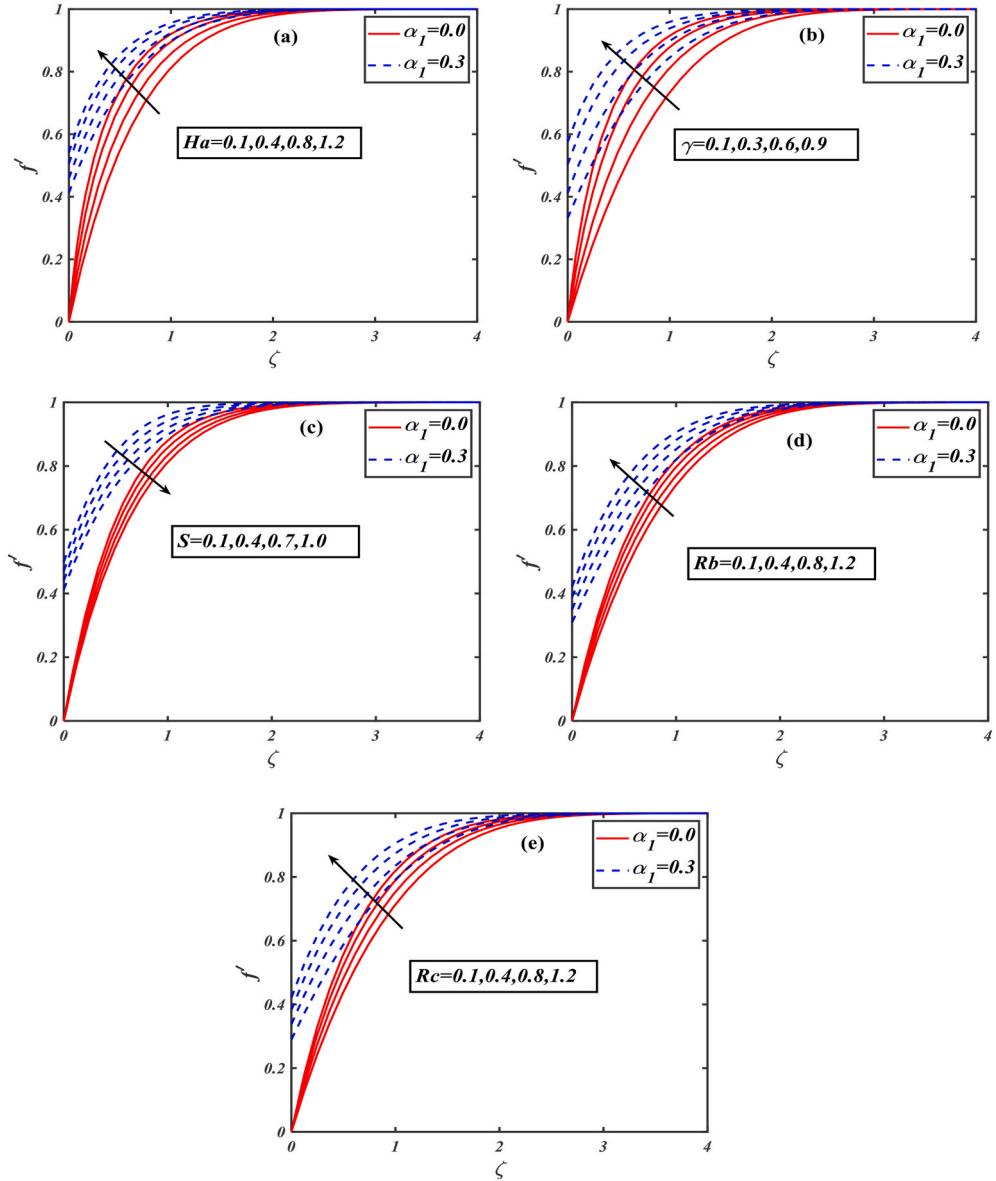
$$u = u_{slip}, \quad v = 0, \quad T = T_w + T_{slip}, \quad C = C_w + C_{slip}, \quad N = N_w + N_{slip} \text{ at } y = 0, \quad (8)$$

$$u = u_e(x) = U_0 x^m, \quad T \rightarrow T_\infty, \quad C \rightarrow C_\infty, \quad N \rightarrow N_\infty \text{ at } y = \infty, \quad (9)$$

where  $T_w$  depicts temperature at the wall,  $C_w$  concentration of nanoparticle at surface and  $N_w$  density of motile microorganisms at the surface and  $u_{slip}$ ,  $T_{slip}$ ,  $C_{slip}$ ,  $N_{slip}$  wall velocity, thermal, concentration of nanoparticles and rescaled density of motile microorganisms slips respectively. Mathematically [52]:

$$u_{slip} = L \frac{\partial u}{\partial y} \left[ \left( 1 + \left( \Lambda \frac{\partial u}{\partial y} \right)^n \right)^{-1} \right], \quad (10)$$

$$T_{slip} = P \frac{\partial T}{\partial y}, \quad C_{slip} = Q_1 \frac{\partial C}{\partial y}, \quad N_{slip} = R \frac{\partial N}{\partial y}, \quad (11)$$



**Fig. 2.** Significance of  $Ha$ ,  $\gamma$ ,  $S$ ,  $Rb$ ,  $Rc$  on  $f'$ .

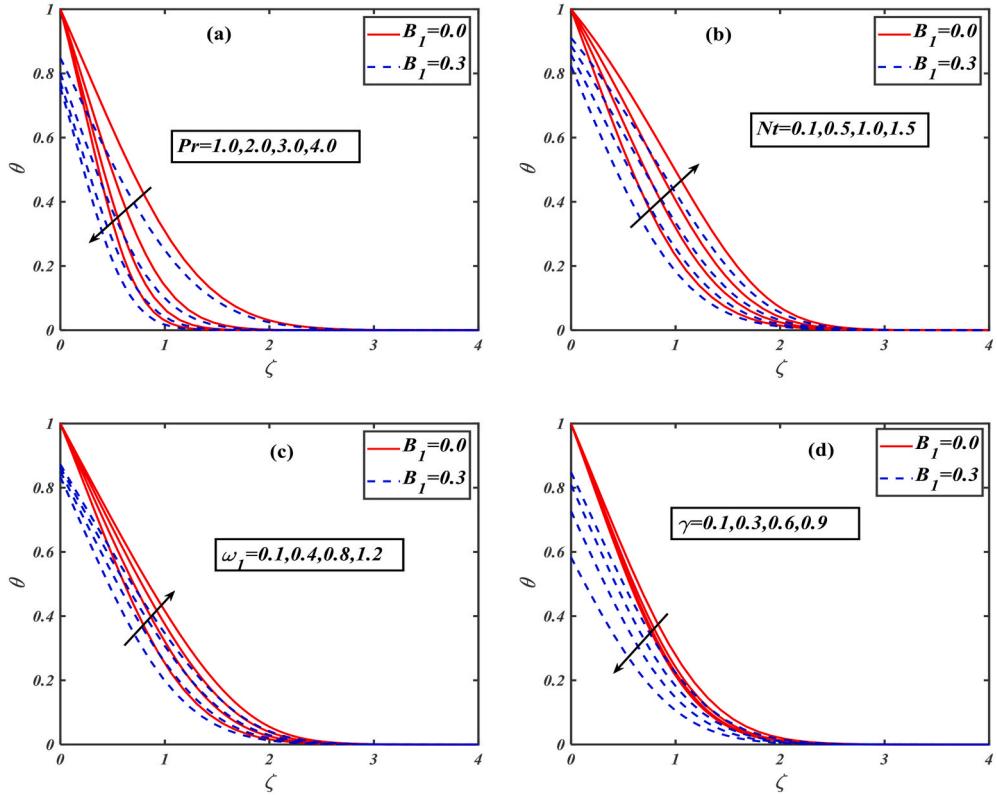
where  $L, P, Q$  and  $R$  are Navier, thermal, volumetric concentration and rescaled density of motile microorganisms slip coefficients correspondingly.

## 2.7. Similarity analysis

The dimensionless variables are taken as [51,61]:

$$\left. \begin{aligned} \psi(x, y) &= \sqrt{\frac{2\nu x U_e}{m+1}} \chi^{\frac{m+1}{2}} f(\zeta), \quad \zeta = y \sqrt{\frac{(m+1)U_e}{2\nu}} x^{\frac{m-1}{2}} \\ \theta(\zeta) &= \frac{T - T_\infty}{T_w - T_\infty}, \quad \varphi(\zeta) = \frac{C - C_\infty}{C_w - C_\infty}, \quad \chi(\zeta) = \frac{N - N_\infty}{N_w - N_\infty} \end{aligned} \right\}, \quad (12)$$

where  $\psi$  shows stream function,  $f(\zeta)$  dimensionless velocity,  $\theta(\zeta)$  dimensionless temperature and  $\varphi(\zeta), \chi(\zeta)$  volumetric concentration of nanoparticles and rescaled density of microorganisms respectively,  $T$  temperature of the fluid,  $C$  fluid concentration and  $N$  mortality of fluid. Equation (1) is satisfied while Eqs. (2)–(9) take the following forms:

Fig. 3. Significance of  $\text{Pr}$ ,  $Nt$ ,  $\omega_1$ ,  $\gamma$  on  $\theta$ .

$$[1 + (1 - n)(We f'')^n]f''' + [ff'' + Ha^2(1 - f')] [1 + (We f'')^n]^2 + \gamma [1 - (f')^2] [1 + (We f'')^n]^2 + S(\theta - Rb\varphi - Rc\chi) = 0, \quad (13)$$

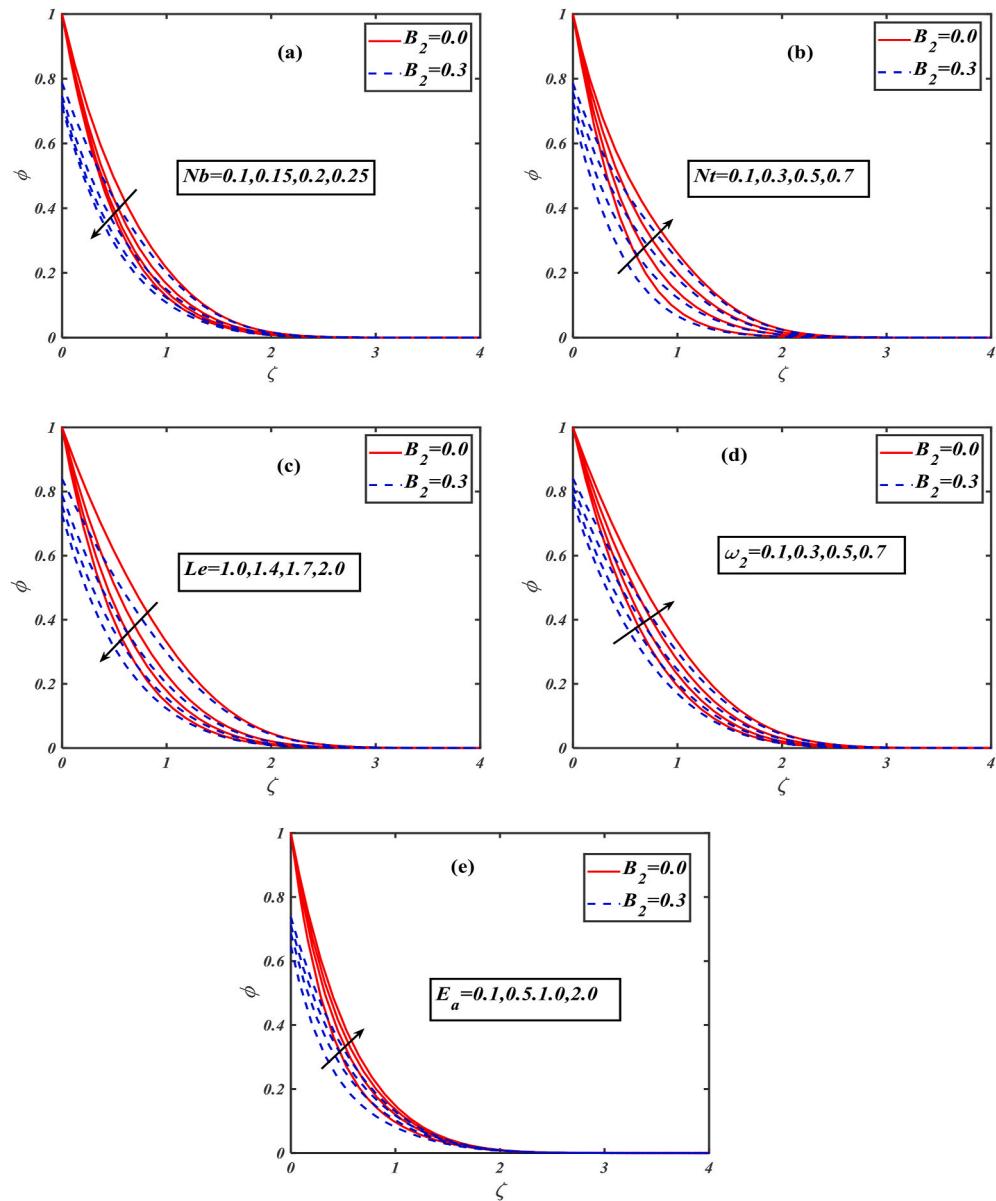
$$(1 + \omega_1 \theta)\theta'' + \omega_1 (\theta')^2 + Pr f\theta' + Pr(Nb\theta'\varphi' + Nt\theta'2) + \delta_0 Pr\theta = 0, \quad (14)$$

$$(1 + \omega_2 \varphi)\varphi'' + \omega_2 (\varphi')^2 + \left(\frac{Nt}{Nb}\right)\theta'' + LePr[f\varphi'] - PrLe\sigma(1 + \omega\theta)^m \exp\left(\frac{-E_a}{1 + \omega\theta}\right)\varphi = 0, \quad (15)$$

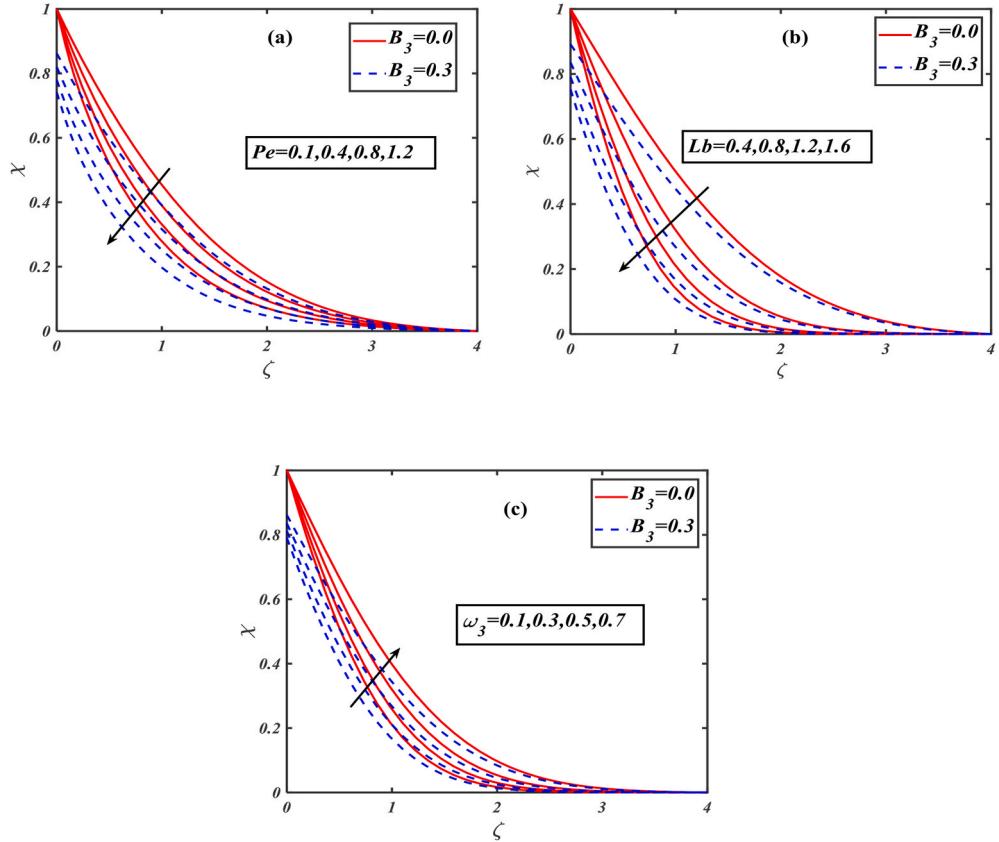
$$(1 + \omega_3 \chi)\chi'' + \omega_3 \chi'^2 + Lbf\chi' - Pe(\varphi''(\chi + \delta_1) + \chi'\varphi') = 0, \quad (16)$$

$$\left. \begin{array}{l} f(0) = 0, f'(0) = \frac{\alpha_1}{\sqrt{2 - \gamma}} f''(0) \left[ \frac{1}{1 + (We f'(0))^n} \right], \\ f'(\infty) \rightarrow 1, \\ \theta(0) = 1 + \frac{B_1}{\sqrt{2 - \gamma}} \theta'(0), \theta(\infty) \rightarrow 0, \\ \varphi(0) = 1 + \frac{B_2}{\sqrt{2 - \gamma}} \varphi'(0), \varphi(\infty) \rightarrow 0, \\ \chi(0) = 1 + \frac{B_3}{\sqrt{2 - \gamma}} \chi'(0), \chi(\infty) \rightarrow 0, \end{array} \right\}, \quad (17)$$

The variables appearing in Eqs. (13)-(17) are local Weissenberg number  $We$ , Prandtl number  $Pr$ , Lewis number  $Le$ , Hartmann number  $Ha$ , temperature difference parameter  $\omega$ , mixed convection parameter  $S$ , buoyancy ratio parameter  $Rb$ , bioconvection Rayleigh number  $Rc$ , Brownian motion parameter  $Nb$ , thermophoresis parameter  $Nt$ , wedge angle parameter  $\gamma$ , heat source-sink parameter  $\delta_0$ , motile microorganisms difference parameter  $\delta_1$ , Peclet number  $Pe$ , bioconvection Lewis number  $Lb$ , velocity slip parameter  $\alpha_1$ , thermal slip parameter  $B_1$ , solutal slip parameter  $B_2$  and density of motile microorganisms slip parameter  $B_3$  which are expressed by



**Fig. 4.** Significance of  $Nb$ ,  $Nt$ ,  $Le$ ,  $\omega_2$ ,  $E_a$  on  $\phi$ .



**Fig. 5.** Significance of  $Pe$ ,  $Lb$ ,  $\omega_3$  on  $\chi$ .

$$\left. \begin{aligned} We &= \sqrt{\frac{F^2(m+1)U_0^3x^{3m-1}}{2\nu}}, \quad Pr = \frac{\mu_0 cp}{k_\infty}, \quad Le = \frac{\alpha}{D_B}, \quad Ha = \frac{2\sigma B_0^2}{\rho U_0(m+1)x^{m-1}}, \quad \omega = \frac{T_w - T_\infty}{T_\infty}, \\ S &= \frac{\beta^* g(1-C_\infty)(T_w - T_\infty)}{(m+1)u_e^2}, \quad Rb = \frac{(\rho_p - \rho_f)(C_w - C_\infty)}{(1-C_\infty)(T_w - T_\infty)\beta^*}, \quad Rc = \frac{\gamma^{**}(\rho_m - \rho_f)(N_w - N_\infty)}{(1-C_\infty)(T_w - T_\infty)\beta}, \\ Nb &= \frac{\tau D_B(C_w - C_\infty)}{\alpha}, \quad Nt = \frac{\tau D_T(T_w - T_\infty)}{T_\infty \alpha}, \quad \gamma = \frac{2m}{m+1}, \quad \delta_0 = \frac{Q}{u_e \rho c_p}, \quad \delta_1 = \frac{N_\infty}{N_w - N_\infty}, \\ Pe &= \frac{bW_c}{D_m}, \quad Lb = \frac{\nu}{D_m}, \quad \alpha_1 = \frac{L}{x} \sqrt{Re_x}, \quad B_1 = \frac{P}{x} \sqrt{Re_x}, \quad B_2 = \frac{Q_1}{x} \sqrt{Re_x}, \quad B_3 = \frac{R}{x} \sqrt{Re_x}, \end{aligned} \right\} \quad (18)$$

## 2.8. Engineering physical quantities of interest

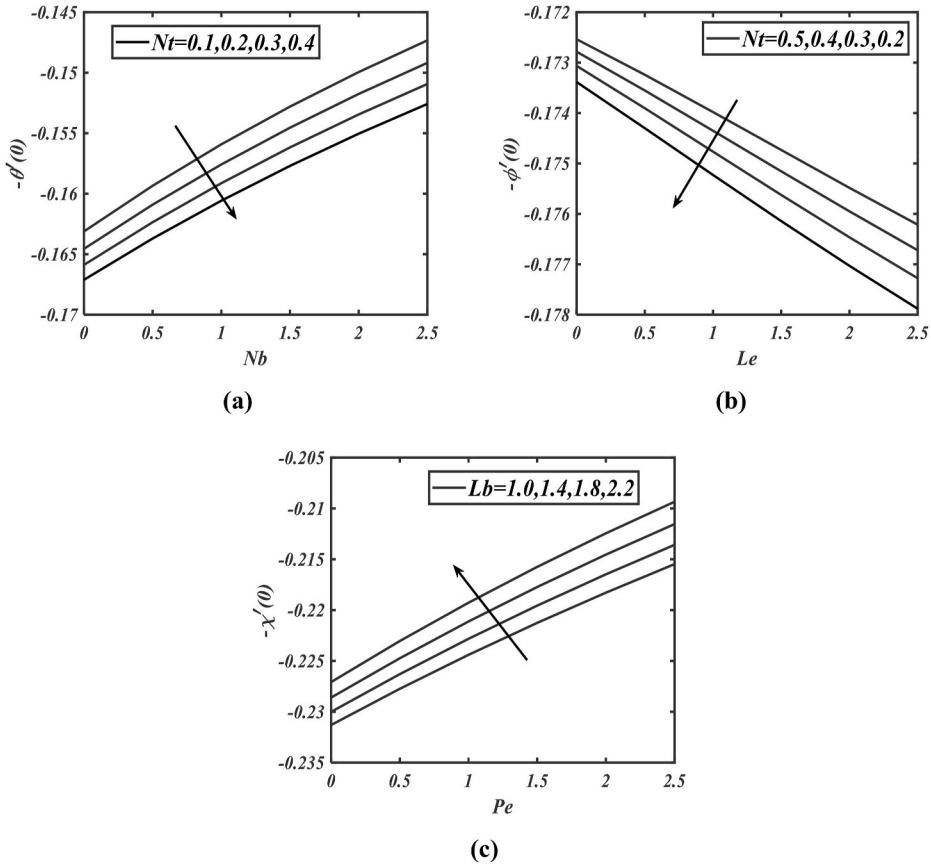
The engineering expressions for drag force, Nusselt, Sherwood and motile numbers are calculated from following definitions:

$$Cf_x = \frac{\tau_w}{\mu u_w^2 / 2}, \quad Nu_x = \frac{xq_w}{k(T_w - T_\infty)}, \quad Sh_x = \frac{xq_m}{D_B(C_w - C_\infty)}, \quad Sh_n = \frac{xq_n}{D_m(N_w - N_\infty)}, \quad (19)$$

In which  $\tau_w, q_w, q_m$  and  $q_n$  are defined below:

$$\begin{aligned} \tau_w &= \mu_0 \frac{\partial u}{\partial y} \left( \left( 1 + \left( \Lambda \frac{\partial u}{\partial y} \right) \right)^{-1} \right) \Big|_{y=0}, \quad q_w = -k \left( \frac{\partial T}{\partial y} \right) \Big|_{y=0}, \quad q_m = -D \left( \frac{\partial C}{\partial y} \right) \Big|_{y=0}, \\ q_n &= -D_m \left( \frac{\partial N}{\partial y} \right) \Big|_{y=0}, \end{aligned} \quad (20)$$

In dimensionless forms one has



**Fig. 6.** Significance of  $Nb$ ,  $Nt$ ,  $Le$ ,  $Pe$ ,  $Lb$  on  $-\theta'(0)$ ,  $-\varphi'(0)$  and  $-\chi'(0)$ .

**Table 1**  
Comparison of numerical computation  $-f''(0)$  with results of Shahzad et al. [51].

Parameters	$n = 0.9$			Present values			
	$We$	$Ha$	$\gamma$	$\alpha_1 = 0.0$	$\alpha_1 = 1.0$	$\alpha_1 = 0.0$	$\alpha_1 = 1.0$
0.0	0.3	1.0		2.536528	1.666042	2.536527	1.666040
0.1				2.400638	1.626165	2.400634	1.625863
0.4				2.092247	1.523816	2.092245	1.523618
0.8	0.5			4.043851	1.426457	4.043853	1.426458
	0.9			3.125334	1.519509	3.125343	1.519510
	1.4			2.433076	1.661463	2.433067	1.661465
0.8	0.3	1.2		2.088515	1.522247	2.088517	1.522246
		1.6		3.879954	1.651639	3.879952	1.651637
		1.8		4.684523	1.458006	4.68525	1.458004

$$\left. \begin{aligned} Re_x^{1/2} Cf_x &= \frac{2}{\sqrt{2-\gamma}} f''(0) \left[ \frac{1}{1 + (We f''(0))^n} \right], \\ Re_x^{-1/2} Nu_x &= -\frac{2}{\sqrt{2-\gamma}} \theta'(0), \\ Re_x^{-1/2} Sh_x &= -\frac{2}{\sqrt{2-\gamma}} \varphi'(0), \\ Re_x^{-1/2} Sh_n &= -\frac{2}{\sqrt{2-\gamma}} \chi'(0), \end{aligned} \right\}, \quad (21)$$

**Table 2**Variations of  $-f''(0)$  for  $Ha$ ,  $n$ ,  $We$ ,  $S$ ,  $Rb$ ,  $Rc$ ,  $\alpha_1$ ,  $\gamma$ .

$Ha$	$n$	$We$	$S$	$Rb$	$Rc$	$\alpha_1$	$\gamma$	$-f''(0)$
0.1	0.5	1.0	0.1	0.1	0.1	0.3	0.2	0.9548
0.2	0.5	1.0	0.1	0.1	0.1	0.3	0.2	1.0417
0.6 1.0	0.1 0.4 0.7	1.0	0.1	0.1	0.1	0.3	0.2	1.1162
0.1	0.5	0.5 1.5 2.5	0.1	0.1	0.1	0.3	0.2	0.9111
0.1	0.5	1.0	0.2 0.5 0.8	0.1	0.1	0.3	0.2	0.9504
0.1	0.5	1.0	0.1	0.2 0.5 0.7	0.1	0.3	0.2	0.9361
0.1	0.5	1.0	0.1	0.1	0.2 0.5 0.7	0.3	0.2	0.9250
0.1	0.5	1.0	0.1	0.1	0.1	0.3	0.2	0.9111
0.1	0.5	1.0	0.1	0.1	0.1	0.3	0.2	0.9406
0.1	0.5	1.0	0.1	0.1	0.1	0.3	0.2	0.9866
0.1	0.5	1.0	0.1	0.1	0.1	0.3	0.2	1.0305
0.1	0.5	1.0	0.1	0.2 0.5 0.7	0.1	0.3	0.2	0.9218
0.1	0.5	1.0	0.1	0.1	0.1	0.3	0.2	0.9126
0.1	0.5	1.0	0.1	0.1	0.2 0.5 0.7	0.3	0.2	0.9064
0.1	0.5	1.0	0.1	0.1	0.1	0.3	0.2	0.9228
0.1	0.5	1.0	0.1	0.1	0.2 0.5 0.7	0.3	0.2	0.9167
0.1	0.5	1.0	0.1	0.1	0.1	0.1	0.2	0.9126
0.1	0.5	1.0	0.1	0.1	0.1	0.1	0.2	1.2606
0.1	0.5	1.0	0.1	0.1	0.1	0.1	0.2	0.8235
0.1	0.5	1.0	0.1	0.1	0.1	0.1	0.2	0.6294
0.1	0.5	1.0	0.1	0.1	0.1	0.3	0.1	0.8160
						0.3	0.4	0.9597
						0.3	0.7	1.0106

### 3. Numerical approach

The local similar solution of Eqs. (13)-(16) supported by the restricting conditions (17) is accomplished through the bvp4c technique [53]. The consequences of effective parameters that are involved in current scrutinization are exhibited through graphical and tabular data for the velocity field, temperature distribution, concentration of nanoparticles and rescaled density of motile microorganisms. Dimensionless ODEs of current flow problems are highly nonlinear in nature. The exact solution of the non-linear system of flow problems is the main challenge of research. The built-in numerical scheme known as bvp4c [54] under the commercial software MATLAB is more applicable to get the numerical solution of the current model. The bvp4c scheme is a different code that implements the Lobbato-IIla formula. Initially higher order differential equations in velocity, temperature, volumetric concentration and motile microorganisms respectively are reduced into the first-order flow problem before starting the process. The detail of these simulations is given. Let

$$\left. \begin{array}{l} f = q_1, f' = q_2, f'' = q_3, f''' = q_3', \\ \theta = q_4, \theta' = q_5, \theta'' = q_5', \\ \varphi = q_6, \varphi' = q_7, \varphi'' = q_7', \\ \chi = q_8, \chi' = q_9, \chi'' = q_9', \end{array} \right\}, \quad (22)$$

$$q_3' = \frac{1}{\beta_1} \left[ [-q_1 q_3 - Ha^2 (1 - q_2) - \gamma (1 - q_2^2)] [1 + (We q_3)^n]^2 - S[q_4 - Rb q_6 - Rc q_8] \right], \quad (23)$$

$$q_5' = \frac{1}{\beta_2} \left[ -\omega_1 q_5^2 - Pr q_1 q_5 - Pr(Nb q_5 q_7 + Nt q_5) - \delta_0 Pr q_4 \right], \quad (24)$$

$$q_7' = \frac{1}{\beta_3} \left[ -\omega_2 q_6^2 - Le Pr [q_1 q_7] - \frac{Nt}{Nb} q_5' + Le Pr \sigma (1 + \omega q_4)^m \exp \left( \frac{-E}{1 + \omega q_4} \right) q_6 \right], \quad (25)$$

$$q_9' = \frac{1}{\beta_4} \left[ -\omega q_9^2 - Lb q_1 q_9 + Pe [q_7'(q_8 + \delta_1) + q_9 q_7] \right], \quad (26)$$

Here

$$\left. \begin{array}{l} \beta_1 = [1 + (1 - n)(We q_3)^n], \\ \beta_2 = (1 + \omega_1 q_4), \\ \beta_3 = (1 + \omega_2 q_6), \\ \beta_4 = (1 + \omega_3 q_9), \end{array} \right\}, \quad (27)$$

**Table 3**Variations of  $-\theta'(0)$  for  $Ha$ ,  $n$ ,  $We$ ,  $S$ ,  $Rb$ ,  $Rc$ ,  $\gamma$ ,  $Pr$ ,  $Nb$ ,  $Nt$ ,  $D$ ,  $B_1$ .

$Ha$	$n$	$We$	$S$	$Rb$	$Rc$	$\gamma$	$Pr$	$Nb$	$Nt$	$\omega_1$	$B_1$	$-\theta'(0)$
		0.5	1.0	0.1	0.1	0.2	1.2	0.2	0.3	0.3	0.3	
0.2	0.5	1.0	0.1	0.1	0.1	0.2	1.2	0.2	0.3	0.3	0.3	0.7400
0.6												0.7520
1.0												0.7602
0.1	0.1	1.0	0.1	0.1	0.1	0.2	1.2	0.2	0.3	0.3	0.3	0.7302
0.4												0.7345
0.7												0.7390
0.1	0.5	0.5	0.1	0.1	0.1	0.2	1.2	0.2	0.3	0.3	0.3	0.7261
1.5												0.7442
2.5												0.7574
0.1	0.5	1.0	0.2	0.1	0.1	0.2	1.2	0.2	0.3	0.3	0.3	0.7383
0.5												0.7451
0.8												0.7515
0.1	0.5	1.0	0.1	0.2	0.1	0.2	1.2	0.2	0.3	0.3	0.3	0.7355
0.5												0.7339
0.7												0.7329
0.1	0.5	1.0	0.1	0.1	0.2	0.2	1.2	0.2	0.3	0.3	0.3	0.7356
1.5												0.7347
2.5												0.7341
0.1	0.5	1.0	0.1	0.1	0.1	0.1	1.2	0.2	0.3	0.3	0.3	0.7201
0.5												0.7408
0.7												0.7473
0.1	0.5	1.0	0.1	0.1	0.1	0.2	2.0	0.2	0.3	0.3	0.3	0.9057
1.5							3.0					1.0416
2.5							4.0					1.1293
0.1	0.5	1.0	0.1	0.1	0.1	0.2	1.2	0.1	0.3	0.3	0.3	0.7707
0.5								0.4				0.6685
0.7								0.7				0.5735
0.1	0.5	1.0	0.1	0.1	0.1	0.2	1.2	0.2	0.1	0.3	0.3	0.7878
0.5								0.4				0.7112
0.7								0.7				0.6418
0.1	0.5	1.0	0.1	0.1	0.1	0.2	1.2	0.2	0.2	0.1	0.3	0.7483
0.5								0.4				0.7301
0.7								0.7				0.7137
0.1	0.5	1.0	0.1	0.1	0.1	0.2	1.2	0.2	0.2	0.3	0.1	0.7939
0.5											0.3	0.7077
0.7											0.7	0.6289

with boundaries

$$\left. \begin{aligned} q_1(0) = 0, \quad q_2(0) = \frac{\alpha_1}{\sqrt{2-\gamma}} q_3 \left[ \frac{1}{1 + (Weq_3)^n} \right], \\ q_4(0) = 1 + \frac{B_1}{\sqrt{2-\chi}} q_5, \quad q_6(0) = 1 + \frac{B_2}{\sqrt{2-\gamma}} q_7, \\ q_8(0) = 1 + \frac{B_3}{\sqrt{2-\gamma}} q_9 \end{aligned} \right\}, \quad (28)$$

$$q_2 \rightarrow 1, \quad q_4 \rightarrow 0, \quad q_6 \rightarrow 0, \quad q_8 \rightarrow 0 \text{ as } \zeta \rightarrow \infty \quad (29)$$

#### 4. Results and discussion

This section is reserved to discuss the significance of different parameters that appear in this article. Impacts of local Weissenberg number, Prandtl number, Lewis number, Hartmann number, bioconvection Rayleigh number, thermal conductivity parameter, microorganism conductivity parameter, wedge angle parameter, mixed convection parameter, bioconvection Lewis number, Peclet number, activation energy and Brownian motion on the velocity of the fluid, temperature profile, concentration profile and motile microorganisms profile are depicted in the graphical illustration.

Fig. 2(a–e) are captured to depicts the behavior of physical parameters namely, Hartmann number  $Ha$ , wedge angle parameter  $\gamma$ , mixed convection parameter  $S$ , bioconvection Rayleigh number  $Rc$  and buoyancy ratio parameter  $Rb$  against velocity field  $f(\zeta)$ . Fig. 2 (a) illustrates the nature of the velocity distribution  $f'(\zeta)$  for different values of the Hartmann number  $Ha$ . In the present sketch, solutions are analyzed for  $Ha = 0.1, 0.4, 0.8$  and  $1.2$ . For the swelling amount of Hartmann number, the velocity profile  $f'(\zeta)$  is declined for both cases  $\alpha_1 = 0.0$  and  $\alpha_1 = 0.3$ . Physically, an increment in Hartmann number produces stronger Lorentz force which cause a

**Table 4**Variations of  $-\varphi'(0)$  for  $Ha$ ,  $n$ ,  $We$ ,  $S$ ,  $Rb$ ,  $Rc$ ,  $\gamma$ ,  $Pr$ ,  $Nb$ ,  $Nt$ ,  $Le$ ,  $E_a$ ,  $B_2$ .

$Ha$	$n$	$We$	$S$	$Rb$	$Rc$	$\gamma$	$Pr$	$Nb$	$Nt$	$Le$	$E_a$	$B_2$	$-\varphi'(0)$
		0.5	1.0	0.1	0.1	0.2	1.2	0.2	0.3	0.1	0.1	0.3	
0.2	0.5	1.0	0.1	0.1	0.1	0.2	1.2	0.2	0.3	0.1	0.1	0.3	0.5029
0.6													0.5082
1.0													0.5117
0.1	0.1	1.0	0.1	0.1	0.1	0.2	1.2	0.2	0.3	0.1	0.1	0.3	0.4983
	0.4												0.5004
	0.7												0.5024
0.1	0.5	0.5	0.1	0.1	0.1	0.2	1.2	0.2	0.3	0.1	0.1	0.3	0.4977
	1.5												0.5040
	2.5												0.5087
0.1	0.5	1.0	0.2	0.1	0.1	0.2	1.2	0.2	0.3	0.1	0.1	0.3	0.5023
	0.5												0.5059
	0.8												0.5093
0.1	0.5	1.0	0.1	0.2	0.1	0.2	1.2	0.2	0.3	0.1	0.1	0.3	0.5007
	0.5												0.4998
	0.7												0.4991
0.1	0.5	1.0	0.1	0.1	0.2	0.2	1.2	0.2	0.3	0.1	0.1	0.3	0.5009
	0.5												0.5004
	0.7												0.5000
0.1	0.5	1.0	0.1	0.1	0.1	0.1	1.2	0.2	0.3	0.1	0.1	0.3	0.4962
	0.4												0.5017
	0.7												0.5106
0.1	0.5	1.0	0.1	0.1	0.1	0.2	2.0	0.2	0.3	0.1	0.1	0.3	0.7151
	0.5						3.0						0.9650
	0.7						4.0						1.1934
0.1	0.5	1.0	0.1	0.1	0.1	0.2	1.2	0.1	0.3	0.1	0.1	0.3	0.0700
	0.4						0.4						0.7842
	0.7						0.7						0.9032
0.1	0.5	1.0	0.1	0.1	0.1	0.2	1.2	0.2	0.1	0.1	0.1	0.3	0.8060
	0.5						0.4						0.3754
	0.7						0.7						0.0935
0.1	0.5	1.0	0.1	0.1	0.1	0.2	1.2	0.2	0.2	0.2	0.1	0.3	0.3230
	0.5						0.4						0.2475
	0.6						0.6						0.1711
0.1	0.5	1.0	0.1	0.1	0.1	0.2	1.2	0.2	0.2	0.1	0.2	0.3	0.4966
	0.5						1.0						0.4723
	2.0						2.0						0.4590
0.1	0.5	1.0	0.1	0.1	0.1	0.2	1.2	0.2	0.2	0.1	0.1	0.1	0.5986
	0.4						0.4						0.4633
	0.7						0.7						0.3779

reduction in flow velocity of the fluid. Therefore, the velocity field reduces.

The consequences of the wedge angle parameter  $\gamma$  against the velocity field  $f'(\zeta)$  are revealed in Fig. 2(b). The growth of the wedge angle parameter  $\gamma$  leads to decay in velocity field  $f'(\zeta)$  for both cases  $\alpha_1 = 0.0$  and  $\alpha_1 = 0.3$ . Fig. 2(c) accounts for the characteristics of mixed convection parameter  $S$  on the velocity field  $f'(\zeta)$ . The velocity field  $f'(\zeta)$  increases for growing estimations of mixed convection parameter  $S$  for both cases  $\alpha_1 = 0.0$  and  $\alpha_1 = 0.3$ . The effects of the buoyancy ratio parameter  $Rb$  on the velocity profile  $f'(\zeta)$  with different values  $Rb = 0.1, 0.4, 0.8$  and  $1.2$  are exhibited in Fig. 2(d). It is detected that the velocity profile  $f'(\zeta)$  declines by varying the values of the buoyancy ratio parameter  $Rb$  on both conditions  $\alpha_1 = 0.0$  and  $\alpha_1 = 0.3$ . Fig. 2(e) is captured to elucidate the features of bioconvection Rayleigh number  $Rc$  against velocity field  $f'(\zeta)$  with various estimations of  $Rc = 0.1, 0.4, 0.8$  and  $1.2$ . The velocity profile  $f'(\zeta)$  deteriorates for progressive variation of bioconvection Rayleigh number  $Rc$ .

Fig. 3(a-d) are displayed to visualize the impacts of involved parameters such as Prandtl number  $Pr$ , wedge angle parameter  $\gamma$ , thermophoresis parameter  $Nt$  and thermal conductivity parameter  $\omega_1$  on the temperature field  $\theta(\zeta)$  for both cases  $B_1 = 0.0$  and  $B_1 = 0.3$ . Fig. 3(a) describes the impact of the Prandtl number  $Pr$  over temperature profile  $\theta(\zeta)$ . The enhancing magnitude of the Prandtl number  $Pr$  diminishes the temperature distribution  $\theta(\zeta)$ . The sketch lines of the thermophoresis parameter  $Nt$  versus temperature distribution  $\theta(\zeta)$  are delineated in Fig. 3(b). This sketch shows the enhancing trend for temperature field  $\theta(\zeta)$  towards  $Nt = 0.1, 0.5, 1.0$  and  $1.5$  for both situations  $B_1 = 0.0$  and  $B_1 = 0.3$ . The features of the thermal conductivity parameter  $\omega_1$  on the temperature field  $\theta(\zeta)$  are demonstrated in Fig. 3(c). From this analysis, it is revealed that temperature distribution  $\theta(\zeta)$  is enhanced for both cases  $B_1 = 0.0$  and  $B_1 = 0.3$  by enhancing the values of the thermal conductivity parameter  $\omega_1$ . Fig. 3(d) manifests the nature of the wedge angle parameter  $\gamma$  on temperature profile  $\theta(\zeta)$  with various amount of  $\gamma = 0.1, 0.3, 0.6$  and  $0.9$ . The temperature distribution  $\theta(\zeta)$  is knockdown by enhancing the variations of wedge angle parameter  $\gamma$  in both cases  $B_1 = 0.0$  and  $B_1 = 0.3$ .

Fig. 4(a-e) demonstrate the characteristics of prominent parameters like Brownian motion parameter  $Nb$ , thermophoresis parameter  $Nt$ , activation energy  $E_a$ , Lewis number  $Le$  and thermal conductivity parameter  $\omega_2$  on the volumetric concentration of

**Table 5**Variations of  $-\chi'(0)$  for  $Ha$ ,  $n$ ,  $We$ ,  $S$ ,  $Rb$ ,  $Rc$ ,  $\gamma$ ,  $Pe$ ,  $Lb$ ,  $B_3$ .

$Ha$	$n$	$We$	$S$	$Rb$	$Rc$	$\gamma$	$Pe$	$Lb$	$B_3$	$-\chi'(0)$
0.1	0.5	1.0	0.1	0.1	0.1	0.2	0.3	1.0	0.3	0.8505
0.2	0.5	1.0	0.1	0.1	0.1	0.2	0.3	1.0	0.3	0.8630
0.6										0.8714
1.0										
0.1	0.1	1.0	0.1	0.1	0.1	0.2	0.3	1.0	0.3	0.8403
0.4										0.8448
0.7										0.8495
0.1	0.5	0.5	0.1	0.1	0.1	0.2	0.3	1.0	0.3	0.8362
		1.5								0.8547
		2.5								0.8682
0.1	0.5	1.0	0.2	0.1	0.1	0.2	0.3	1.0	0.3	0.8588
		0.5								0.8559
		0.8								0.8627
0.1	0.5	1.0	0.1	0.2	0.1	0.2	0.3	1.0	0.3	0.8458
		0.5								0.8442
		0.7								0.8431
0.1	0.5	1.0	0.1	0.1	0.2	0.2	0.3	1.0	0.3	0.8460
		0.5								0.8450
		0.7								0.8424
0.1	0.5	1.0	0.1	0.1	0.1	0.1	0.3	1.0	0.3	0.8340
		0.4								0.8489
		0.7								0.8502
0.1	0.5	1.0	0.1	0.1	0.1	0.2	0.2	1.0	0.3	0.9054
		0.6								0.7070
		1.0								0.5811
0.1	0.5	1.0	0.1	0.1	0.1	0.2	0.3	1.2	0.3	0.9371
		1.6								1.0991
		2.0								1.2913
0.1	0.5	1.0	0.1	0.1	0.1	0.2	0.3	1.0	0.1	0.8470
		0.4								0.8472
		0.7								0.8495

nanoparticles  $\varphi(\zeta)$  for both cases  $B_2 = 0.0$  and  $B_2 = 0.3$ . The behavior of the Brownian motion parameter  $Nb$  on the volumetric concentration of nanoparticles  $\varphi(\zeta)$  is depicted in Fig. 4(a). As expected concentration field of nanoparticles  $\varphi(\zeta)$  is reducing function of the Brownian motion parameter  $Nb$  for both cases  $B_2 = 0.0$  and  $B_2 = 0.3$ . Fig. 4(b) is the representation of concentration profile  $\varphi(\zeta)$  due to thermophoresis parameter  $Nt$  at different values of  $Nt = 0.1, 0.3, 0.5$  and  $0.7$ . The concentration field of nano-shape particles  $\varphi(\zeta)$  upsurges by growing thermophoresis parameter  $Nt$ . The physical significance of the Lewis number  $Le$  versus nanoparticle concentration  $\varphi(\zeta)$  is interpreted in Fig. 4(c). Here for both cases  $B_2 = 0.0$  and  $B_2 = 0.3$ , the concentration profile  $\varphi(\zeta)$  is retarded when the Lewis number  $Le$  enhances. Fig. 4(d) is devoted to scrutinize the estimation in the concentration field  $\varphi(\zeta)$  for different values of the thermal conductivity parameter  $\omega_2$ . It is analyzed that the concentration field  $\varphi(\zeta)$  increases for higher peaches of thermal conductivity parameter  $\omega_2$ . Fig. 4(e) is engrossed to depict the impact of activation energy  $E_a$  on the volumetric concentration field of nanoparticles  $\varphi(\zeta)$  for both situations  $B_2 = 0.0$  and  $B_2 = 0.3$ . It can be perceived that volumetric concentration  $\varphi(\zeta)$  is enhanced by an increment in the activation energy  $E_a$ .

Fig. 5(a–c) are elucidated to interpret the significance of bioconvection Lewis number  $Lb$ , Peclet number  $Pe$  and microorganism conductivity parameter  $\omega_3$  on the motile microorganisms field  $\chi(\zeta)$  for both cases  $B_3 = 0.0$  and  $B_3 = 0.3$ . The curves of motile microorganism's field  $\chi(\zeta)$  versus Peclet number  $Pe$  are elaborated in Fig. 5(a). It is found that the microorganisms field  $\chi(\zeta)$  declines by an increment in the variation of Peclet number  $Pe$  for  $B_3 = 0.0$  and  $B_3 = 0.3$ . Fig. 5(b) prevail the interesting scenario about the bioconvection Lewis number  $Lb$  on motile microorganism field  $\chi(\zeta)$  for  $B_3 = 0.0$  and  $B_3 = 0.3$ . It is noticed that the rescaled density of motile microorganisms  $\chi(\zeta)$  reduces by enlarging the bioconvection Lewis number  $Lb$ . Fig. 5(c) is examined the effect of the microorganism conductivity parameter  $\omega_3$  on the swimming microorganisms profile  $\chi(\zeta)$ . It is concluded that the motility of fluid  $\chi(\zeta)$  decreases at higher peaches of microorganism conductivity parameter  $\omega_3$  for both cases  $B_3 = 0.0$  and  $B_3 = 0.3$ .

Fig. 6(a–c) show the effects of  $Nb$ ,  $Nt$ ,  $Le$ ,  $Pe$ ,  $Lb$  on  $-\theta'(0)$ ,  $-\varphi'(0)$  and  $-\chi'(0)$  respectively. Here  $-\theta'(0)$  and  $-\varphi'(0)$  are enhanced for larger values of thermophoresis parameter. Furthermore, the  $-\chi'(0)$  is an increasing function of  $Lb$ .

In Table 1, to attained the numerical solution of the skin friction coefficient for a chosen variation of Weissenberg number  $We$ , Hartmann number  $Ha$  and wedge angle parameter  $\gamma$  is compared with those studied by Shahzad et al. [51] and found to be in good agreement.

Table 2 is drawn to observe the behavior of involving parameters such as  $Ha$ ,  $n$ ,  $We$ ,  $S$ ,  $Rb$ ,  $Rc$ ,  $\alpha_1$ ,  $\gamma$  on skin friction coefficient. It is noticed that the skin friction coefficient enhances for the variation of  $Ha$ ,  $Rb$  and  $\gamma$ .

The features of Nusselt number against prominent parameters are presented in Table 3. The Nusselt number rises for a different variation of  $\omega_1$ ,  $S$  while decreases for  $Pr$ .

Table 4 is constructed to examine the behavior of  $Ha$ ,  $n$ ,  $We$ ,  $S$ ,  $Rb$ ,  $Rc$ ,  $\gamma$ ,  $Pr$ ,  $Nb$ ,  $Nt$ ,  $D$  and  $B_1$  versus Sherwood number. With growing amount of

$Ha$ , the Sherwood number reduces.

Table 5 explicates the significance of rescaled density number of motile microorganisms versus  $Ha$ ,  $n$ ,  $We$ ,  $S$ ,  $Rb$ ,  $Rc$ ,  $\gamma$ ,  $Pe$ ,  $Lb$ ,  $B_3$ . The rescaled density number of motile microorganisms is reduced for higher variations of  $Pe$ ,  $Lb$ .

## 5. Conclusions

This article discusses the heat and mass transfer features in non-Newtonian (cross fluid) nano-material flow with multiple slip conditions. The salient features of motile microorganism, nonlinear heat source/sink and activation energy are accounted. The key points of current article are.

- Velocity profile exaggerates for higher values of mixed convection parameter while opposite behavior is observed for buoyancy ratio parameter and bioconvection Lewis number.
- The temperature distribution is a reducing function of the Prandtl number [55].
- The temperature field is enhanced for thermophoresis parameter and thermal conductivity parameter.
- The volumetric concentration of nanoparticles increases for growing thermal conductivity parameter.
- The density of motile microorganisms depicts diminishing trend for bioconvection Lewis number and Peclet number [56–65].
- The current mathematical flow model is more useful in the field of nanotechnology, electrical and mechanical engineering, biotechnology, biofuel, microbiology etc.
- In future, one can extend this work by considering magnetic dipole, convective boundary conditions, homogeneous-heterogeneous reactions, nonlinear thermal radiation etc.

## Declaration of competing interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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