

FRACTIONAL HEAT EQUATION OPTIMIZED BY A CHAOTIC FUNCTION

by

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In this effort, we propose a new fractional differential operator in the open unit disk. The operator is an extension of the Atangana-Baleanu differential operator without singular kernel. We suggest it for a normalized class of analytic functions in the open unit disk. By employing the extended operator, we study the time-2-D space heat equation and optimizing its solution by a chaotic function.

Key words: fractional calculus, thermal, heat equation, subordination, chaotic, univalent function, analytic function

Introduction

The class of fractional heat equations is investigated by many researchers. They modeled different physical environments, involving time-space, random walks, non-local transport theory and delayed flux-force associations [1-4]. Moreover, some investigators introduced a general physical introduction to fractional diffusion equations, motivated by Atangana-Baleanu differential operator [5] to simulate heat transfer processes. Optimization by using chaotic functions is used in financial studies. Chaotic functions play a significant role in improving diffusion, symmetry ergodicity and stochasticity of chaos.

In this work, we shall optimize the fractional heat equation type time-2-D space in terms of a special class of chaotic functions, which is used to define a chaotic map [6]. Our method is based on the majorization and subordination theory in the open unit disk [7, 8]. For two analytic functions φ and ψ in the open unit disk $\mathbb{U} = \{z \in \mathbb{C} : |z| < 1\}$, we say that φ is majorized by ψ ($\varphi \ll \psi$) if there is an analytic function $\varpi, |\varpi| < 1$ such that $\varphi(z) = \varpi(z)\psi(z)$. Moreover, φ is subordinated to ψ if $\varphi(z) = \psi[\varpi(z)]$, [9].

Preparation

A fractional differential operator for the complex Atangana and Baleanu is defined [10]:

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$${}^C \Delta^\nu g(z) = \frac{\beta(\nu)}{2\pi i(1-\nu)} \int_{\mathbb{D}} f'(\zeta) \Xi_\nu[-\mu_\nu(z-\zeta)^\nu] d\zeta \quad (1)$$

where $\beta(\nu)$ is normalized by $\beta(0) = \beta(1) = 1$ and $\Xi_\nu(\omega)$ is the Mittag-Leffler function. Moreover, they introduced the following fractional differential operator:

$${}^R \Delta^\nu g(z) = \frac{\beta(\nu)}{2\pi i(1-\nu)} \frac{d}{dz} \int_{\mathbb{D}} f(\zeta) \Xi_\nu[-\mu_\nu(z-\zeta)^\nu] d\zeta \quad (2)$$

$$\mu_\nu = \frac{\nu}{1-\nu}, \quad \nu \in (0,1), \quad \mathbb{D} = [z + re^{i\pi}(z-\ell) : 0 < r < 1]$$

To modify the previous operators, we define a class of analytic functions by $f(z) = z + \sum_{n=2}^{\infty} a_n z^n$, $z \in \mathbb{U}$. This class is denoted by Λ calling the class of univalent functions and normalized by $f(0) = f'(0) - 1 = 0$.

Definition 1 Let $f \in \Lambda$. Then the modified operators of (1) and (2) are given by the integrals, respectively:

$${}^C \Delta_z^\nu f(z) = \frac{\beta(\nu)}{1-\nu} \int_0^z f'(\zeta) \Xi_{\nu,\nu}(-\mu_\nu \zeta^\nu) \Xi_\nu[-\mu_\nu(z-\zeta)^\nu] d\zeta \quad (3)$$

and

$${}^R \Delta_z^\nu f(z) = \frac{\beta(\nu)}{1-\nu} \frac{d}{dz} \int_0^z f(\zeta) \Xi_{\nu,\nu}(-\mu_\nu \zeta^\nu) \Xi_\nu[-\mu_\nu(z-\zeta)^\nu] d\zeta \quad (4)$$

where ν indicates the power of z .

For example, let $f(z) = z$, then by [11], Theorem 2.4 or [12], Theorem 11.2, we have:

$$\begin{aligned} {}^C \Delta_z^\nu(z) &= [\beta(\nu)/1-\nu] \int_0^z \Xi_\nu(-\mu_\nu \zeta^\nu) \Xi_\nu[-\mu_\nu(z-\zeta)^\nu] d\zeta = \\ &= [\beta(\nu)/1-\nu] z \Xi_{\nu,2}^2[-\mu_\nu(z)^\nu] = \\ &= [\beta(\nu)/1-\nu] z \sum_{k=0}^{\infty} \frac{(2)_k z^k}{k! \Gamma(k\nu+2)}, \quad (\wp)_0 = 1, \quad (\wp)_n = \wp(\wp+1)\dots(\wp+n-1) \end{aligned}$$

And in view of [11], Theorem 2.2, we have:

$$\begin{aligned} {}^R \Delta_z^\nu(z) &= [\beta(\nu)/1-\nu] \frac{d}{dz} \int_0^z \Xi_\nu(-\mu_\nu \zeta^\nu) \Xi_\nu[-\mu_\nu(z-\zeta)^\nu] \zeta d\zeta = \\ &= [\beta(\nu)/1-\nu] \{ z^2 \Xi_{\nu,3}^2[-\mu_\nu(z)^\nu] \} = \\ &= [\beta(\nu)/1-\nu] \{ z \Xi_{\nu,2}^2[-\mu_\nu(z)^\nu] \} \end{aligned}$$

It is clear that ${}^C \Delta_z^\nu(z) = {}^R \Delta_z^\nu(z)$. In general, we have:

$${}^C \Delta_z^\nu(z^n) = [\beta(\nu) / 1 - \nu](n-1)z^n \{\Xi_{\nu,1+n}^2[-\mu_\nu(z)^\nu]\}, \quad n > 1$$

$${}^R \Delta_z^\nu(z^n) = [\beta(\nu) / 1 - \nu]z^n \{\Xi_{\nu,1+n}^2[-\mu_\nu(z)^\nu]\}$$

We have the following property.

Proposition 1 Consider the operators (3) and (4) for $f \in \Lambda$. Then by letting $b(\nu) := [\beta(\nu) / 1 - \nu]$:

- ${}^e \Delta_z^\nu f(z) := \frac{{}^C \Delta_z^\nu f(z)}{b(\nu)\Xi_{\nu,2}^2[-\mu_\nu(z)^\nu]} \in \Lambda$ and ${}^R \Delta_z^\nu f(z) := \frac{{}^R \Delta_z^\nu f(z)}{b(\nu)\Xi_{\nu,2}^2[-\mu_\nu(z)^\nu]} \in \Lambda$
- ${}^R \Delta_z^\nu f(z) \ll {}^e \Delta_z^\nu f(z)$
- ${}^R \Delta_z^\nu f(z) \prec {}^e \Delta_z^\nu f(z)$, provided that ${}^e \Delta_z^\nu f(z)$, is locally univalent of the first order (like convex function [13]) when $|z| \in (0.28, \sqrt{2}-1)$ or locally univalent of the second order (like the class of univalent functions [13]) when $|z| \in (0.21, 0.3)$.

Proof 1 Let $f \in \Lambda$. Then a direct computation yields:

$$\begin{aligned} \frac{{}^C \Delta_z^\nu f(z)}{b(\nu)\Xi_{\nu,2}^2[-\mu_\nu(z)^\nu]} &= \frac{b(\nu)\Xi_{\nu,2}^2[-\mu_\nu(z)^\nu]z + \sum_{n=2}^{\infty} a_n b(\nu)(n-1)\{\Xi_{\nu,1+n}^2[-\mu_\nu(z)^\nu]\}z^n}{b(\nu)\Xi_{\nu,2}^2[-\mu_\nu(z)^\nu]} = \\ &= z + \sum_{n=2}^{\infty} a_n(n-1) \left\{ \frac{\Xi_{\nu,1+n}^2[-\mu_\nu(z)^\nu]}{\Xi_{\nu,2}^2[-\mu_\nu(z)^\nu]} \right\} z^n \Rightarrow {}^e \Delta_z^\nu f(z) \in \Lambda \end{aligned}$$

Similarly, we obtain ${}^R \Delta_z^\nu f(z) \in \Lambda$. This completes the first part. For the second part, it is sufficient to prove that, [14], $|{}^R \Delta_z^\nu f(z)| \leq |{}^e \Delta_z^\nu f(z)|$. A computation yields:

$$\begin{aligned} |{}^R \Delta_z^\nu f(z)| &= \left| z + \sum_{n=2}^{\infty} a_n \left\{ \frac{\Xi_{\nu,1+n}^2[-\mu_\nu(z)^\nu]}{\Xi_{\nu,2}^2[-\mu_\nu(z)^\nu]} \right\} z^n \right| \leq \\ &\leq \left| z + \sum_{n=2}^{\infty} a_n(n-1) \left\{ \frac{\Xi_{\nu,1+n}^2[-\mu_\nu(z)^\nu]}{\Xi_{\nu,2}^2[-\mu_\nu(z)^\nu]} \right\} z^n \right| = \\ &= |{}^e \Delta_z^\nu f(z)| \end{aligned}$$

The last part immediately comes from [14] Corollary 1 and 2, respectively. Based on *Proposition 1*, we shall focus on ${}^e \Delta_z^\nu f(z)$.

Heat equation associated with ${}^e \Delta_z^\nu$

The Koebe function is an extreme function in the field of geometric function theory. To determine the heat equation associated with ${}^e \Delta_z^\nu$, we deal with the parametric Koebe function of the form:

$$f_{\sigma}(t, z) = z/(1-tz)^{\sigma} = z + \sum_{n=2}^{\infty} \frac{(\sigma)_{n-1}}{(n-1)!} t^{n-1} z^n, \quad t < |z| < 1$$

Then the generalized heat equation is given by:

$$\Psi(t, z) = [{}^c \Delta_z^{\nu} f_{\sigma}(t, z)]_t - [{}^c \Delta_z^{\nu} f_{\sigma}(t, z)]_{zz}, \quad z \in \mathbb{U} \tag{5}$$

Our aim is to optimize the solution of (5) by the chaotic function, fig. 1:

$$\begin{aligned} \sin[z/(1-tz)^{\sigma}] &= z + \sigma tz^2 + z^3 [1/2\sigma(\sigma+1)t^2 - 1/6] + 1/6\sigma tz^4 [(\sigma+1)(\sigma+2)t^2 - 3] + \\ &+ 1/120z^5 [-60\sigma^2 t^2 + 5\sigma(\sigma+1)(\sigma+2)(\sigma+3)t^4 - 30\sigma(\sigma+1)t^2 + 1] + \\ &+ 1/120\sigma tz^6 [-10(9\sigma^2 + 9\sigma + 2)t^2 + (\sigma^4 + 10\sigma^3 + 35\sigma^2 + 50\sigma + 24)t^4 + 5] + O(z^7) \\ &:= z + \sum_{n=2}^{\infty} \zeta_n(\sigma, t) z^n, \quad t \leq |z| < 1 \end{aligned} \tag{6}$$

Note that $\sin(\omega)$ is univalent in the disk $|z| < \pi/2$, see [15].

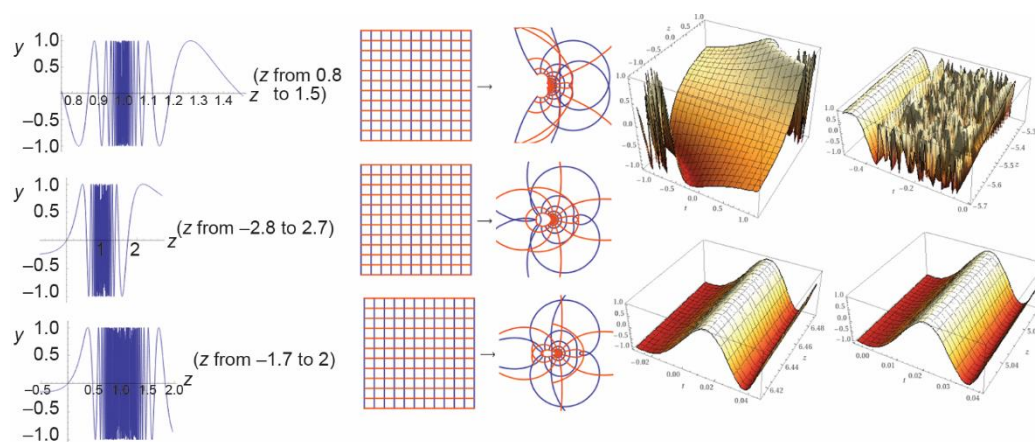


Figure 1. The plot of $\sin[z/(1-tz)^{\sigma}]$, when $t = 1, \sigma = 1, 2, 3$; the last two columns are 2-D plot for $\sigma = 1, 2, 3, 4$ (for color image see journal web site)

Theorem 1 Consider the heat eq. (5). For a small value of $\nu \in [0,1]$, the solution of (5) is optimized by the chaotic function $\sin[z/(1-tz)^{\sigma}]$.

Proof 2 By Proposition 1, we indicate that $\Psi(t,0) = 0$. Also, $\nu \rightarrow 0$, implies:

$$\left(\frac{\Xi_{\nu,1+n}^2[-\mu_{\nu}(z)^{\nu}]}{\Xi_{\nu,2}^2[-\mu_{\nu}(z)^{\nu}]} \right) \approx 1, \text{ then}$$

$$\begin{aligned}
 {}^c \Delta_z^\nu f(t, z) &= z + \sum_{n=2}^{\infty} \frac{(\sigma)_{n-1}}{(n-1)!} (n-1) \left\{ \frac{\Xi_{\nu,1+n}^2[-\mu_\nu(z)^\nu]}{\Xi_{\nu,2}^2[-\mu_\nu(z)^\nu]} \right\} t^{n-1} z^n \approx \\
 &\approx z + \sum_{n=2}^{\infty} \frac{(\sigma)_{n-1}}{(n-1)!} (n-1) t^{n-1} z^n = \\
 &= z + \sum_{n=2}^{\infty} \frac{(\sigma)_{n-1}}{(n-2)!} t^{n-1} z^n = \\
 &:= z + \sum_{n=2}^{\infty} \kappa_n(\sigma, t) z^n
 \end{aligned} \tag{7}$$

To optimize the solution of (5), it is sufficient to show that $|\kappa_n(\sigma, t)| \leq |\zeta_n(\sigma, t)|$. This means that we must find the value of σ whenever $t < 1$. A comparison between the coefficients $|\kappa_n(\sigma, t)|$ and $|\zeta_n(\sigma, t)|$, we obtain the value $0 < \sigma \leq 1/\sqrt{3} \approx 0.57735\dots$. This completes the proof.

Corollary 1 Consider the heat equation (5). Then for $\nu, t \rightarrow 1$:

$${}^c \Delta_z^\nu f(t, z) \prec \sin[z/(1-tz)^\sigma], \quad 0.21 < |z| < 0.3$$

Proof 3 In view of *Theorem 1*, we have ${}^c \Delta_z^\nu f(t, z) \ll \sin[z/(1-tz)^\sigma]$. Since $\sin(\omega)$ is univalent and $[{}^c \Delta_z^\nu f(t, 0)]_z = 1 > 0$, then in view of [14] *Corollary 2*, we conclude that ${}^c \Delta_z^\nu f(t, z) \prec \sin[z/(1-tz)^\sigma]$.

Corollary 2 Consider the heat eq. (5). Then for $t \rightarrow 1$:

$$[{}^c \Delta_z^\nu f(t, z)]_z \ll \{\sin[z/(1-tz)^\sigma]\}_z, \quad |z| \leq 0.26794$$

Proof 4 In view of *Theorem 1*, we obtain ${}^c \Delta_z^\nu f(t, z) \ll \sin[z/(1-tz)^\sigma]$. According to [14] *Theorem 1*, where $\sin(\omega)$ is of the second kind of locally univalent function, we get the require assertion.

Remark 1 In view of *Proposition 1 (C)* and *Corollary 1*, we confirm that:

$$\Re {}^c \Delta_z^\nu f(t, z) \prec {}^c \Delta_z^\nu f(t, z) \prec \sin[z/(1-tz)^\sigma], \quad |z| \in (0.21, 0.3)$$

Conclusion

We formulated a modified Atangana-Baleanu differential operator of a class of normalized analytic functions in the open unit disk. We presented a new generalization of time-2-D heat equations based on the suggested operator. Analytic solution is indicated by using the chaotic function $\sin[z/(1-tz)^\sigma]$. The optimal solution is appeared when $\sigma = 0.57735$ (see fig. 2). For future works, one may suggest another class of analytic function in the open unit disk such as meromorphic, multivalent and harmonic functions.

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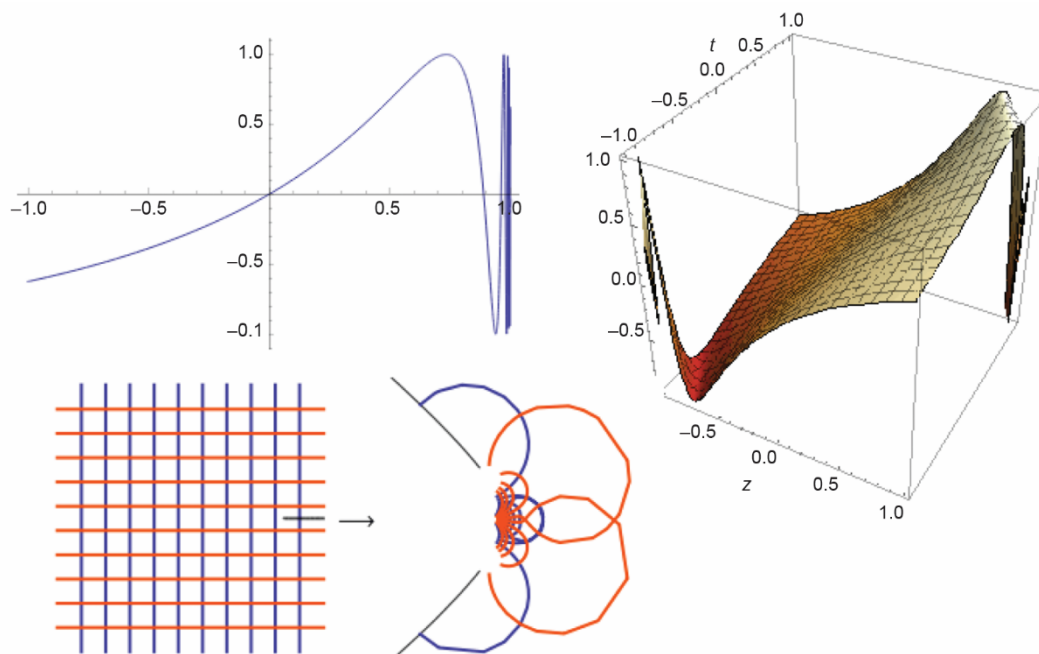


Figure 2. The plot of $\sin[z/(1 - tz)^{0.577}]$, which is the optimal solution of heat eq. (5)

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