# On the asymptotic integration of a class of sublinear fractional differential equations 

Dumitru Băleanu ${ }^{1, a)}$ and Octavian G. Mustafa ${ }^{2, b)}$<br>${ }^{1}$ Department of Mathematics and Computer Science, Çankaya University, Ögretmenler Cad. 14, Balgat, Ankara 06530, Turkey<br>${ }^{2}$ Department of Mathematics, DAL, University of Craiova, Craiova 200534, Romania

(Received 31 March 2009; accepted 4 November 2009; published online 22 December 2009)

We estimate the growth in time of the solutions to a class of nonlinear fractional differential equations $D_{0+}^{\alpha}\left(x-x_{0}\right)=f(t, x)$ which includes $D_{0+}^{\alpha}\left(x-x_{0}\right)=H(t) x^{\lambda}$ with $\lambda \in(0,1)$ for the case of slowly decaying coefficients $H$. The proof is based on the triple interpolation inequality on the real line and the growth estimate reads as $x(t)=o\left(t^{a \alpha}\right)$ when $t \rightarrow+\infty$ for $1>\alpha>1-a>\lambda>0$. Our result can be thought of as a noninteger counterpart of the classical Bihari asymptotic integration result for nonlinear ordinary differential equations. By a carefully designed example we show that in some circumstances such an estimate is optimal. © 2009 American Institute of Physics. [doi:10.1063/1.3271111]

## I. INTRODUCTION

The fractional calculus deals with the generalization of differentiation and integration to noninteger orders. Fractional calculus is as old as the classical differential calculus and it gained more and more importance during the last three decades in many fields of science ad engineering. ${ }^{10-12,31,26,34,32,19,37,23,20,36,35,22,29}$

The fractional order differential equations play a significant role in modeling the anomalous dynamics of various processes related to complex systems in most areas of science and engineering. However, only a few steps have been taken toward what may be called a coherent theory of these equations in the applied sciences in a manner analogous to the ordinary case.

We would like to recall that the fractional differential equations have been found effective to describing various physical phenomena such as damping laws, rheology, and diffusion processes. ${ }^{18,17}$ Various applications of fractional differential equations in the reaction kinetics of proteins, the anomalous electron transport in amorphous materials, the dielectrical or mechanical relaxation of polymers, the modeling of glass-forming liquids, and others are presented in Refs. $25,16,33$, and 24.

On the theoretical side, several fractional variational principles were formulated and the corresponding Euler-Lagrange equations consisting of left and the right derivatives were analyzed during the last few years. ${ }^{2,4,5,21}$ The fractional derivatives are also the infinitesimal generators of a class of translation invariant convolution semigroups which appear universally as attractors. ${ }^{37,22}$

New methods for solving the nonlinear fractional partial differential equations were developed recently. ${ }^{27,28}$

Let us consider the initial value problem for a fractional differential equation (FIVP) below,

[^0]\[

\left\{$$
\begin{array}{l}
D_{0+}^{\alpha}\left(x-x_{0}\right)(t)=f(t, x(t)), \quad t>0  \tag{1}\\
x(0)=x_{0}
\end{array}
$$\right.
\]

where the nonlinearity $f: D=[0,+\infty) \times \mathbb{R} \rightarrow \mathrm{R}$ is assumed continuous and $x_{0}$ is some real number.
The differential operator $D_{0+}^{\alpha}$ in problem (1) is the Riemann-Liouville differential operator of order $0<\alpha<1$, namely,

$$
D_{0+}^{\alpha}(u)(t)=\frac{1}{\Gamma(1-\alpha)} \cdot \frac{d}{d t}\left[\int_{0}^{t} \frac{u(s)}{(t-s)^{1-\alpha}} d s\right]
$$

where $\Gamma(1-\alpha)=\int_{0}^{+\infty} e^{-t} t^{-\alpha} d t$ is the Gamma function. See Ref. 20, p. 70.
The above fractional differential equation represents a wide class of dissipative nonlinear nonlocal equations, therefore the asymptotic behavior of its solutions is an interesting topic of investigation.

The motivation for inserting the initial datum $x_{0}$ into the differential operator comes from the physical origin of such mathematical models and the reader can find comprehensive details in this respect in Refs. 32, p. 80, and Ref. 13, p. 230. The Riemann-Liouville operator with an inserted datum is called a Caputo differential operator-see Refs. 9 and 20, p. 91.

Assuming that the FIVP has a solution $x(t)$, the formulas $\Gamma(\alpha) \Gamma(1-\alpha)=\pi / \sin \pi \alpha$ and

$$
\int_{0}^{t} f(s, x(s)) d s=\frac{\sin \pi \alpha}{\pi} \int_{0}^{t} \frac{1}{(t-s)^{\alpha}} \int_{0}^{s} \frac{f(\tau, x(\tau))}{(s-\tau)^{1-\alpha}} d \tau d s
$$

see Ref. 3, p. 196, allow us to rewrite (1) via an integration as

$$
\begin{equation*}
\int_{0}^{t} \frac{1}{\Gamma(1-\alpha)(t-s)^{\alpha}}\left[x(s)-x_{0}-\frac{1}{\Gamma(\alpha)} \int_{0}^{s} \frac{f(\tau, x(\tau))}{(s-\tau)^{1-\alpha}} d \tau\right] d s=0 \tag{2}
\end{equation*}
$$

where $t>0$.
We can regard from now on, by means of (2), any solution of FIVP (1) as a (continuous) solution of the singular integral equation,

$$
\begin{equation*}
x(t)=x_{0}+\frac{1}{\Gamma(\alpha)} \int_{0}^{t} \frac{f(s, x(s))}{(t-s)^{1-\alpha}} d s, \quad t>0 \tag{3}
\end{equation*}
$$

and vice versa. This fact is established in a rigorous manner in Ref. 13. See also the monograph (Ref. 20, Theorem 3.24, pp. 199-200) (for $\gamma=r=0$ in the original notation).

Given the generality of $f$, we can replace it with $\Gamma(\alpha)^{-1} f$, which means that we are interested in analyzing the next integral equation

$$
\begin{equation*}
x(t)=x_{0}+\int_{0}^{t} \frac{f(s, x(s))}{(t-s)^{1-\alpha}} d s, \quad t \geq 0 . \tag{4}
\end{equation*}
$$

Our aim here is to investigate the long-time behavior of the solutions to a large class of Eqs. (4). Noticing that, for a constant $f$ in (4), the solution reads as

$$
x(t)=x_{0}+\frac{f}{\alpha} \cdot t^{\alpha}=O\left(t^{\alpha}\right) \quad \text { as } t \rightarrow+\infty
$$

we shall compare the growth of solutions to the nonlinear integral equation (4) with the natural test function $t^{\alpha}$.

To the best of our knowledge, the literature of fractional differential equations is scarce with respect to such enterprises. Most of the attempts regard exact solutions to (very) restricted types of fractional differential equations. For example, in Ref. 16 the linear fractional differential equation $\tau_{0}^{-\alpha} D_{0+}^{\alpha}[\Phi(t)]+\Phi(t)-\Phi(0)=0$ has been used to model the complex protein dynamics. The solu-
tions were expressed using a Mittag-Leffler transcendental function of order $\alpha$, namely, $\Phi(t)$ $=\Phi(0) E_{\alpha}\left(-\left(t / \tau_{0}\right)^{\alpha}\right)$, see Ref. 16, p. 49. So, $\Phi(t) \sim t^{-\alpha}$ as the time $t$ increases indefinitely.

The paper is organized as follows. In Sec. II we sketch an asymptotic integration theory for nonlinear ordinary differential equations of integer order that motivates our investigation. In Sec. III we introduce the hypotheses concerning the functional quantity $f(t, x)$ from 3 . The main result is established in Sec. IV. Section V contains a detailed presentation of the fundamental class of nonlinear differential equations given by $f(t, x)=H(t) x^{\lambda}$, with $\alpha>\lambda>0$. Section VI is devoted to an analysis of the sharpness of our asymptotic estimate for the continuable solutions to (3).

Let us conclude this section by stating that, even though we have chosen to use the simple Riemann-Liouville differential operator inside the nonlinear fractional differential equation (1) for reasons of clarity in the exposition, the technique presented here works equally well for most of the (far more complicated) differential operators of noninteger order that can be encountered in both the mathematics and physics literature. Another point that has to be emphasized is that, due to the interpolation-type approach described in this paper, we allow the coefficient $H$ to be more flexible (both in the vicinity of the initial moment 0 and for the large values of $t$ ) than it is usually allowed in the recent investigations based on the (weighted) Hölder-type inequalities.

## II. THE BIHARI ASYMPTOTIC INTEGRATION THEORY

In two seminal papers ${ }^{7,8}$ on the asymptotic integration of nonlinear ordinary differential equations of integer order, Bihari established that, given the equation

$$
x^{\prime \prime}+f(t, x)=0, \quad t \geq t_{0}>0
$$

where

$$
|f(t, x)| \leq a(t) \cdot g\left(\frac{|x|}{t}\right), \quad x \in \mathbb{R}, \quad t \geq t_{0}
$$

and $f, a, g$ are continuous functions, $g$ is monotone nondecreasing and

$$
\int_{t_{0}}^{+\infty} a(t) d t<+\infty
$$

all the continuable, real-valued, solutions $x$ satisfy the conditions

$$
\lim _{t \rightarrow+\infty} x^{\prime}(t)=\lim _{t \rightarrow+\infty} \frac{x(t)}{t}=l(x) \in \mathbb{R}
$$

that is, they resemble with a certain degree of accuracy to the lines $l(x) \cdot t$.
For the case $l(x)=0$, of interest for applications in mathematical physics, see also the recent contribution. ${ }^{30}$

The core of Bihari's approach relies on using his famous integral inequality, see the presentation from Ref. 1, to establishing that, if a solution $x$ is globally defined in the future, the quantity

$$
\frac{x(t)}{t}, \quad t \geq t_{0}
$$

is bounded in $\left[t_{0},+\infty\right)$.
This implies further that the improper integral,

$$
\int_{t_{0}}^{+\infty} f(s, x(s)) d s=\lim _{t \rightarrow+\infty} \int_{t_{0}}^{t} f(s, x(s)) d s=x^{\prime}\left(t_{0}\right)-\lim _{t \rightarrow+\infty} x^{\prime}(t)
$$

is convergent, its value being $x^{\prime}\left(t_{0}\right)-l(x)$.

To search for possible answers to the following question-are there any noninteger counterparts of the classical Bihari asymptotic integration result?-we have thus (1) to produce a quantity like $x(t) / t$ that will ultimately be bounded and further (2) to establish that $\lim x^{\prime}(t)$ is finite.

In the following, the quantity from (1) will be

$$
\mathcal{B}(x, t)=\frac{x(t)}{(t+1)^{\alpha}}
$$

The part (2), however, cannot be translated into the noninteger case: this is a consequence of the nonlocal nature of the differential operator $D_{0+}^{\alpha}$.

To answer affirmatively to the preceding question, we have replaced the finiteness of $\lim x^{\prime}(t)$, which is equivalent with the finiteness of $\lim x(t) / t$, with a condition in the same spirit, $t \rightarrow+\infty$ namely,

$$
\lim _{t \rightarrow+\infty} t^{(1-a) \alpha} \cdot \frac{x(t)}{(t+1)^{\alpha}} \in \mathbb{R}
$$

for some $a \in(0,1)$.

## III. PRELIMINARIES

Several technical requirements for the nonlinearity $f$ will be presented next.
Given $\alpha \in(0,1)$, set $a \in(0,1)$ and $p_{2}>1$, such that $\alpha>1-a$ and

$$
\begin{equation*}
(1-a) \alpha=\frac{(1-a)\left(\alpha+\alpha^{2}\right)}{1+\alpha}<\frac{1}{p_{2}}<\frac{(1-a) \alpha+\alpha^{2}}{1+\alpha}<\alpha \tag{5}
\end{equation*}
$$

Introduce further $p_{1}, p_{3}>1$ via the formulas

$$
\frac{1}{p_{1}}+\frac{1}{p_{2}}=\frac{1+(1-a) \alpha}{1+\alpha} \quad \text { and } \quad p_{3}=\frac{1+\alpha}{a \alpha}
$$

Notice that

$$
\begin{equation*}
\frac{1}{p_{1}}+\frac{1}{p_{2}}+\frac{1}{p_{3}}=1 \tag{6}
\end{equation*}
$$

which makes $p_{1}, p_{2}, p_{3}$ perfect candidates for the triple interpolation inequality on the real line, see, e.g., Ref. 15, p. 146. For further use, remark that the double inequality regarding $p_{2}$ from (5) is equivalent with

$$
\begin{equation*}
1-\alpha<\frac{1}{p_{1}}<\frac{1-(1-a) \alpha^{2}}{1+\alpha} \tag{7}
\end{equation*}
$$

Consider now the continuous, nondecreasing function $g:[0,+\infty) \rightarrow[0,+\infty)$ with the following features:

$$
g(0)=0, \quad g(\xi)>0 \quad \text { when } \quad \xi>0
$$

Moreover-take $x_{0} \neq 0$ in (4)—

$$
\begin{equation*}
\lim _{u \rightarrow+\infty} \frac{W(u)}{u^{a}}=+\infty, \quad W(u)=\int_{0}^{u} \frac{d \xi}{\left[g\left(\left|x_{0}\right|+\xi^{1 / p_{3}}\right)\right]^{p_{3}}} . \tag{8}
\end{equation*}
$$

An example of $g$ is that of $g(u)=u^{1-(1+\zeta) a}$ for some $\zeta \in(0,1)$ and $a \in(0,1 /(1+\zeta))$. It is easy to see that

$$
\begin{aligned}
W(u) & =\int_{0}^{u} \frac{d \xi}{\left[\left(\left|x_{0}\right|+\xi^{1 / p_{3}}\right)^{1-(1+\zeta) a}\right]^{p_{3}}} \geq \int_{\left|x_{0}\right|^{p_{3}}}^{u} \frac{d \xi}{\left[\left(2 \cdot \xi^{1 / p_{3}}\right)^{1-(1+\zeta) a}\right]^{p_{3}}} \\
& =2^{-[1-(1+\zeta) a] p_{3}} \cdot \frac{u^{(1+\zeta) a}-\left|x_{0}\right|^{(1+\zeta) a p_{3}}}{(1+\zeta) a} \sim c \cdot u^{(1+\zeta) a} \quad \text { when } u \rightarrow+\infty
\end{aligned}
$$

where $c=2^{-[1-(1+\zeta) a] p_{3}} \cdot 1 /(1+\zeta) a$.
Remark 1: A crucial fact should be emphasized at this point. We shall use in Secs. V and VI the function $g$ given by $g(u)=u^{\lambda}$ for $\alpha>1-a>\lambda$. This restriction upon $\lambda$, which can be recast as $1-\lambda>a$, is natural in the above computation since we can take $(1+\zeta) a=1-\lambda$ for a $\zeta$ small enough.

We introduce next the class of Bihari-like nonlinearities $f$ by means of the inequality

$$
\begin{equation*}
|f(t, x)| \leq h(t) g\left(\frac{|x|}{(t+1)^{\alpha}}\right) \tag{9}
\end{equation*}
$$

where the function $h:[0,+\infty) \rightarrow[0,+\infty)$ is continuous and such that

$$
\begin{equation*}
t^{\left(p_{3} / p_{1}\right)\left[1-p_{1}(1-\alpha)\right]} \cdot\|h\|_{L^{p_{2}(0, t)}}^{p_{3}}=O\left(t^{\alpha}\right) \quad \text { when } t \rightarrow+\infty \tag{10}
\end{equation*}
$$

Remark 2: The ratio inside $g$ from (9) is designed, following the tradition of Ref. 8, pp. 277-278, according to our aim of studying the so-called long-time behavior (that is, for $t \rightarrow+\infty$ ) of solutions to (1). This is why the choice for the argument of $g$ looks sloppy when $t$ is close to 0 . We emphasize that it is only the quest for clarity and simplicity that has lead us to introducing such a peculiar element $(t+1)^{\alpha}$. A careful inspection of the computations in the following will allow the interested reader to adapt the description of $f$ for the vicinity of 0 . This principle is also valid for the important issue of (any) solution $x$ of (3) being or not being allowed to take the values $\pm \infty$ for $t=0$. Not to complicate the analysis, we opted for asking $x$ to be continuous as well in 0 . This limitation yields $x(0)=x_{0}$ for all the solutions of (1) and (3).

Remark 3: Our main interest in this paper is about fractional differential equations where the nonlinearity $f(t, x)$ is powerlike. This is motivated by many applied examples: logistic equations used in population dynamics and biochemistry (e.g., polymer growth) and Emden-Fowler and Thomas-Fermi type of equations (from stellar dynamics where the chaotic expansions of the outer space gases might be modeled with fractional differential equations). The architecture of the integral $W$ from Eq. (8) has been chosen, however, as to give a certain freedom to the nonlinearity. One can have even some nondeterministic/nonsmooth alterations of the powerlike form.

The restriction (10) can be recast as

$$
\begin{equation*}
t^{\left(p_{3} / p_{1}\right)\left[1-p_{1}(1-\alpha)\right]} \cdot\left\{\int_{0}^{t}[h(s)]^{p_{2}} d s\right\}^{p_{3} / p_{2}} \leq M(t+1)^{\alpha}, \quad t \geq 0 \tag{11}
\end{equation*}
$$

for some sufficiently large constant $M$.
Remark 4: Notice the inequality of the exponents of $t$, namely,

$$
\frac{p_{3}}{p_{1}}\left[1-p_{1}(1-\alpha)\right]<\alpha
$$

which is equivalent with the second part of the double inequality (7). This means that the class of coefficients $h$ is quite substantial, including, for instance, the space $L^{p_{2}}(0,+\infty)$. It is a common practice in the asymptotic analysis of ordinary differential equations of integer order to study carefully the case of slowly decaying coefficients $h$, e.g., $h(t) \sim t^{-\varepsilon}$ for some $\varepsilon \in(0,1)$ when $t \rightarrow$ $+\infty$. This issue seems not to be addressed yet in the literature devoted to asymptotic integration of fractional differential equations. For this reason, we hope that our interpolation-type approach will lead to further investigations in this area of research.

## IV. A BIHARI-TYPE RESULT FOR FRACTIONAL DIFFERENTIAL EQUATIONS

Our main contribution in this work is given in the sequel.
Theorem 1: Assume that the function from (4) verifies the conditions (9) and (11). Then, all the continuable solutions (that is, continuous solutions defined in $[0,+\infty)$ ) of the FIVP (1) and (3) have the asymptotic behavior

$$
\begin{equation*}
x(t)=o\left(t^{a \alpha}\right) \quad \text { as } t \rightarrow+\infty \tag{12}
\end{equation*}
$$

Remark 5: We have refrained here from discussing the delicate matter of the global existence of solutions to (3). This requires the tools of functional analysis and its insertion would have lead us too far afield. However, a long but straightforward adaptation of the procedures from Ref. 1 or Ref. 14 will conclude that such globally defined solutions of the FIVP (1) and (3) do exist. See also Lemma 3 of the present paper and Ref. 6 for a significant particular case.

Proof: We split the demonstration into two steps.
In step 1, we shall establish the raw asymptotic description of solutions to (1) and (3) given by the formula

$$
\begin{equation*}
x(t)=o\left(t^{\alpha}\right) \quad \text { when } t \rightarrow+\infty . \tag{13}
\end{equation*}
$$

To this end, consider $x(t)$ a continuable solution of the integral equation (4) and set $y(t)$ $=x(t)-x_{0}$ for every $t \geq 0$.

We have the estimates

$$
\begin{aligned}
|y(t)| & \leq \int_{0}^{t} \frac{1}{(t-s)^{1-\alpha}} \cdot h(s) \cdot g\left(\frac{|y(s)|+\left|x_{0}\right|}{(s+1)^{\alpha}}\right) d s \leq\left\{\int_{0}^{t}\left[\frac{1}{(t-s)^{1-\alpha}}\right]^{p_{1}}\right\}^{1 / p_{1}} \cdot\|h\|_{L^{p_{2}(0, t)}} \cdot[z(t)]^{1 / p_{3}}, \quad t \\
& >0
\end{aligned}
$$

where

$$
z(t)=\int_{0}^{t}\left[g\left(\frac{|y(s)|+\left|x_{0}\right|}{(s+1)^{\alpha}}\right)\right]^{p_{3}} d s
$$

Further-recall (11)—,

$$
\begin{equation*}
|y(t)|^{p_{3}} \leq c \cdot t^{\left(p_{3} / p_{1}\right)\left[1-p_{1}(1-\alpha)\right]}\|h\|_{L^{p_{2}(0, t)}}^{p_{3}} \cdot z(t) \leq M_{1}(t+1)^{\alpha} \cdot z(t), \quad t \geq 0 \tag{14}
\end{equation*}
$$

with $c=\left[1-p_{1}(1-\alpha)\right]^{-p_{3} / p_{1}}$ and $M_{1}=c \cdot M$.
We deduce now that

$$
z^{\prime}(t) \leq\left[g\left(|y(t)|+\left|x_{0}\right|\right)\right]^{p_{3}} \leq g\left(\left|x_{0}\right|+\left[M_{1}(t+1)^{\alpha} \cdot z(t)\right]^{1 / p_{3}}\right)^{p_{3}}
$$

Fix $t_{0}>0$. Notice that, when $t \in\left[0, t_{0}\right]$, the preceding inequality reads as

$$
z^{\prime}(t) \leq g\left(\left|x_{0}\right|+\left[c_{1} \cdot z(t)\right]^{1 / p_{3}}\right)^{p_{3}}, \quad c_{1}=M_{1}\left(1+t_{0}\right)^{\alpha},
$$

and

$$
\frac{\left[c_{1} z(t)\right]^{\prime}}{g\left(\left|x_{0}\right|+\left[c_{1} z(t)\right]^{1 / p_{3}}\right)^{p_{3}}} \leq c_{1}
$$

An integration with respect to $t$ leads to

$$
\begin{equation*}
W\left(c_{1} z(t)\right) \leq c_{1} t \leq c_{1} t_{0}<M_{1}\left(1+t_{0}\right)^{\alpha+1}, \quad t \in\left[0, t_{0}\right] . \tag{15}
\end{equation*}
$$

Replacing $t$ with $t_{0}$ in (15), we get-recall (14)—

$$
\begin{equation*}
W\left(|y(t)|^{p_{3}}\right) \leq W\left(M_{1}(t+1)^{\alpha} z(t)\right) \leq M_{1}(t+1)^{\alpha+1}, \quad t \geq 0 . \tag{16}
\end{equation*}
$$

The inequality (16) can be rewritten as

$$
\begin{equation*}
\frac{W\left(|y(t)|^{p_{3}}\right)}{\left[|y(t)|^{p_{3}}\right]^{a}} \cdot\left[\frac{|y(t)|}{(t+1)^{\alpha}}\right]^{a p_{3}} \cdot \frac{(t+1)^{a \alpha p_{3}}}{(t+1)^{\alpha+1}} \leq M_{1} \tag{17}
\end{equation*}
$$

Notice that the third factor from the left-hand side of (17) is 1 , that is,

$$
\begin{equation*}
\frac{W\left(|y(t)|^{p_{3}}\right)}{\left[|y(t)|^{p_{3}}\right]^{a}} \cdot\left[\frac{|y(t)|}{(t+1)^{\alpha}}\right]^{a p_{3}} \leq M_{1}, \quad t \geq 0 \tag{18}
\end{equation*}
$$

Suppose now, for the sake of contradiction, that $y(t)$ is not $o\left(t^{\alpha}\right)$ when $t \rightarrow+\infty$. This means that there exist an increasing, unbounded from above, sequence of positive numbers $\left(t_{n}\right)_{n \geq 1}$ and $\varepsilon$ $>0$ such that

$$
\begin{equation*}
\frac{\left|y\left(t_{n}\right)\right|}{\left(t_{n}\right)^{\alpha}} \geq \varepsilon, \quad n \geq 1 \tag{19}
\end{equation*}
$$

Since (19) implies that $\lim \left|y\left(t_{n}\right)\right|=+\infty$, we have-recall (8)—

$$
\begin{aligned}
& n \rightarrow+\infty \\
& \lim _{n \rightarrow+\infty} \frac{W\left(\left|y\left(t_{n}\right)\right|^{p_{3}}\right)}{\left[\left|y\left(t_{n}\right)\right|^{p_{3}}\right]^{a}}=\lim _{u \rightarrow+\infty} \frac{W(u)}{u^{a}}=+\infty .
\end{aligned}
$$

According to the inequality (18), we obtain that $\lim \left|y\left(t_{n}\right)\right| /\left(t_{n}\right)^{\alpha}=0$ and this contradicts (19). $n \rightarrow+\infty$
Step 2: According to the first step, $y(t)=o\left(t^{\alpha}\right)$ for all the large values of $t$. This yields, via the L'Hôpital rule, that

$$
\lim _{t \rightarrow+\infty} \frac{z(t)}{t}=\lim _{t \rightarrow+\infty} \frac{\int_{0}^{t}\left[g\left(\frac{|y(s)|+\left|x_{0}\right|}{(s+1)^{\alpha}}\right)\right]^{p_{3}} d s}{t}=\lim _{t \rightarrow+\infty} g\left(\frac{|y(t)|+\left|x_{0}\right|}{(t+1)^{\alpha}}\right)=g(0)=0 .
$$

Further, by means of the estimate (14), we conclude that

$$
|y(t)|=o\left(t^{(1+\alpha) / p_{3}}\right)=o\left(t^{a \alpha}\right) \quad \text { when } t \rightarrow+\infty .
$$

The proof is complete.

## V. APPLICATIONS

To illustrate this work, let us fix $\alpha, \lambda \in(0,1)$ with $\alpha>\lambda$. There exists $\zeta \in(0,1)$ with the property that

$$
\begin{equation*}
\alpha>\frac{\lambda+\zeta}{1+\zeta} \tag{20}
\end{equation*}
$$

Set also $a=a(\zeta)=(1-\lambda) /(1+\zeta)$.
The number $p_{2}>1$ satisfying (5) can now be introduced. Noticing that the function $\zeta \mapsto(\lambda$ $+\zeta) /(1+\zeta)=1-a(\zeta)$ is increasing, the conditions (5) can be simplified using the stronger restriction,

$$
\begin{equation*}
\lambda \alpha<\frac{\lambda+\zeta}{1+\zeta} \cdot \alpha<\frac{1}{p_{2}}<\lambda \alpha+(1-\lambda) \frac{\alpha^{2}}{1+\alpha}=\frac{\lambda \alpha+\alpha^{2}}{1+\alpha}<\frac{\frac{\lambda+\zeta}{1+\zeta} \cdot \alpha+\alpha^{2}}{1+\alpha} \tag{21}
\end{equation*}
$$

In fact, let us fix $p_{2}$ such that

$$
\begin{equation*}
\lambda \alpha<\frac{1}{p_{2}}<\lambda \alpha+(1-\lambda) \frac{\alpha^{2}}{1+\alpha} . \tag{22}
\end{equation*}
$$

Then, we can replace (20) and (21) with the more practical requirements for $\zeta$ given by

$$
\alpha>\frac{\lambda+\zeta}{1+\zeta}, \quad \frac{\lambda+\zeta}{1+\zeta} \cdot \alpha<\frac{1}{p_{2}} .
$$

Now, the data regarding the numbers $a,\left(p_{i}\right)_{1 \leq i \leq 3}$ can be summarized in the following algorithmlike structure.

Step 1: Take $\alpha, \lambda$ such that

$$
1>\alpha>\lambda>0 .
$$

Step 2: Take $p_{2}$ such that

$$
\left[\lambda \alpha+(1-\lambda) \frac{\alpha^{2}}{1+\alpha}\right]^{-1}<p_{2}<(\lambda \alpha)^{-1} .
$$

Step 3: Take $\zeta$ such that

$$
0<\zeta<\min \left\{\frac{\alpha-\lambda}{1-\alpha}, \frac{1-\lambda \alpha p_{2}}{\alpha p_{2}-1}\right\} .
$$

Step 4: Set

$$
a=\frac{1-\lambda}{1+\zeta} .
$$

Step 5: Set

$$
p_{1}=\left[\frac{1+(1-a) \alpha}{1+\alpha}-\frac{1}{p_{2}}\right]^{-1} .
$$

Step 6: Set

$$
p_{3}=\frac{1+\alpha}{a \alpha} .
$$

Consider the initial value problem for a fractional (general) logistic equation with sublinear exponents,

$$
\left\{\begin{array}{l}
D_{0+}^{\alpha}\left(x-x_{0}\right)=H(t) x^{\lambda}\left[1+P(t) x^{-\mu}\right], \quad t>0 \\
x(0)=x_{0},
\end{array}\right.
$$

where $\mu \in(0, \lambda)$, the function $P:[0,+\infty) \rightarrow \mathrm{R}$ is continuous and bounded, and

$$
H(t)=(t+1)^{-\alpha \lambda} h(t), \quad h \in L^{p_{2}}(0,+\infty) .
$$

Then, all the positive-valued, continuable solutions of the initial value problem have the asymptotic behavior $x(t)=o\left(t^{\alpha}\right)$ when $t \rightarrow+\infty$, according to the first step of Theorem 1. This lack of sharpness in the asymptotic estimate is explained in the next paragraph.

Notice that

$$
\begin{aligned}
\left|H(t) x^{\lambda}\left[1+P(t) x^{-\mu}\right]\right| & \leq(t+1)^{-\alpha \lambda} h(t) \cdot\left(1+\|P\|_{\infty}\right)\left(x^{\lambda}+x^{\eta}\right) \leq 2\left(1+\|P\|_{\infty}\right) h(t)\left[\frac{1+x}{(t+1)^{\alpha}}\right]^{\lambda} \\
& \leq h(t) g\left(\frac{x}{(t+1)^{\alpha}}\right), \quad t \geq 0, \quad x>0
\end{aligned}
$$

where $\eta=\lambda-\mu \in(0,1)$ and

$$
g(u)=2\left(1+\|P\|_{\infty}\right)(1+u)^{\lambda} .
$$

Since this function $g$ does not verify the hypothesis $g(0)=0$, we can extract from the inferences of Theorem 1 only the conclusion of the first step.

Let us observe also that for $a, \lambda$, and $p_{2}$ subjected to the preceding restrictions, that is, 0 $<\lambda<\alpha<1$ and $p_{2}$ confined by (22), all the continuable solutions of the sublinear fractional differential equation,

$$
\left\{\begin{array}{l}
D_{0+}^{\alpha}\left(x-x_{0}\right)=H(t) x^{\lambda}, \quad t>0  \tag{23}\\
x(0)=x_{0}
\end{array}\right.
$$

where

$$
H(t)=(t+1)^{-\alpha \lambda} h(t), \quad h \in L^{p_{2}}(0,+\infty),
$$

have the asymptotic behavior $x(t)=o\left(t^{a \alpha}\right)$ when $t \rightarrow+\infty$, according to Theorem 1 .
As detailed in Sec. III, the asymptotic integration results for these initial value problems are manageable illustrations of our technique. Recently, their superlinear variant, namely, the case when $P \equiv 0, \alpha \in\left(\frac{1}{2}, 1\right)$, and $\lambda>1$ has been discussed in the beautiful contribution, ${ }^{14}$ however, under a quite demanding restriction on the coefficient $h$-for instance, $h \in L^{p}(0,+\infty)$ for all $p$ $\in(0,1)$. Such a restriction excludes all hope for an analysis of the long-time behavior of solutions in the case of slowly decaying coefficients. We emphasize the fact that, even though most real-life applications of (fractional) differential equations involve nonlinear equations, in the research literature such enterprises are scarce. In this respect, the contribution ${ }^{14}$ is quite singular since its methods are mathematically rigorous. On the formal level, a solution of type $O\left(t^{\alpha}\right)$ as $t \rightarrow+\infty$ for $\alpha=\frac{1}{2}$ and $\lambda=4$ has been presented in Ref. 32, p. 236, for a fractional model of radiative cooling of a semi-infinite body.

## VI. A DIRECT ANALYSIS

Section VI deals with the crucial issue of sharpness for our asymptotic estimate in Theorem 1. We claim that the asymptotic estimate given in this paper by means of a Bihari-type approach in conjunction with an interpolation inequality is almost optimal. To establish the validity of our claim, we use a method frequently encountered in the asymptotic integration theory of differential equations of integer order, namely, we build an "extreme" example and show that, with respect to it, our theorem does the best job ever possible.

Before proceeding with the computations, let us mention that, similarly to the case of the classical Bihari investigation, ${ }^{1}$ we obtain here the following type of error: for an exact solution having the numerical values described by $O\left(t^{\varepsilon \cdot q}\right)$ when $t \rightarrow+\infty$, where $\varepsilon \in\left(\varepsilon_{0}, 1\right)$ for a prescribed $\varepsilon_{0}>0$ and $q>0$, the asymptotic behavior obtained using inequalities and averaging reads as $o\left(t^{q}\right)$ for $t \rightarrow+\infty$. Since we can set $\varepsilon_{0}$ as close to 1 as desired, it is obvious that the estimate provided by such a Bihari-type analysis is optimal.

Consider now the integral equation below,

$$
\begin{equation*}
x(t)=x_{0}+\int_{0}^{t} \frac{H(s)}{(t-s)^{1-\alpha}}[x(s)]^{\lambda} d s, \quad t>0, \tag{24}
\end{equation*}
$$

where $x_{0}>0$ and $\alpha, \lambda \in(0,1)$ are fixed numbers such that

$$
\begin{equation*}
1>\frac{3 \alpha}{2} \quad \text { and } \quad \alpha>2 \lambda \tag{25}
\end{equation*}
$$

Here, $H:[0,+\infty) \rightarrow[0,+\infty)$ is a continuous function.
The next lemmas address three fundamental issues regarding the investigation of (24): the sign of $x$, its uniqueness, and its global existence in the future.

Lemma 1: If $x:[0, T) \rightarrow \mathrm{R}$, where $T \leq+\infty$, is a continuous solution of Eq. (24) such that $x(0)=x_{0}$ then

$$
\begin{equation*}
x(t) \geq x_{0}, \quad t \in[0, T) . \tag{26}
\end{equation*}
$$

Proof: Notice that it is enough to prove that $x$ takes positive values everywhere in $[0, T)$. In fact, if we assume that $x$ has a zero in $[0, T)$ then, given its continuity, there exists the number $t_{0}>0$, such that-recall that $x(0)=x_{0}>0$ -

$$
x(t)>0 \quad \text { for all } t \in\left[0, t_{0}\right), \quad x\left(t_{0}\right)=0 .
$$

By taking into account (24), we get

$$
x\left(t_{0}\right)=x_{0}+\int_{0}^{t_{0}} \frac{H(s)}{\left(t_{0}-s\right)^{1-\alpha}}[x(s)]^{\lambda} d s \geq x_{0}
$$

We have reached a contradiction.
Finally, since $x(t)>0$ throughout $[0, T)$, we obtain (26).
Lemma 2: Given the number $x_{0}>0$, there exists at most one continuous solution $x:[0, T)$ $\rightarrow \mathrm{R}$, where $T \leq+\infty$, of Eq. (24) such that $x(0)=x_{0}$.

Proof: Assume that there exist two continuous solutions $x, y:[0, T) \rightarrow \mathbb{R}$ of Eq. (24) with $x(0)=y(0)=x_{0}$. Then, via Lemma 1, we have

$$
x(t), y(t) \geq x_{0} \quad \text { for all } t \in[0, T)
$$

The mean value theorem, applied to the function $k:\left(x_{0} / 2,+\infty\right) \rightarrow(0,+\infty)$ with the formula $k(\xi)=\xi^{\lambda}$, yields

$$
\left|x^{\lambda}-y^{\lambda}\right|=\left|\frac{\lambda}{\zeta^{1-\lambda}} \cdot(\max \{x, y\}-\min \{x, y\})\right| \leq \lambda\left(\frac{2}{x_{0}}\right)^{1-\lambda} \cdot|x-y|,
$$

where $\xi$ lies between $x$ and $y$.
The latter estimate implies that

$$
\begin{equation*}
|x(t)-y(t)| \leq \int_{0}^{t} \frac{H(s)}{(t-s)^{1-\alpha}} \cdot\left|[x(s)]^{\lambda}-[y(s)]^{\lambda}\right| d s \leq c \cdot \int_{0}^{t} \frac{H(s)}{(t-s)^{1-\alpha}} \cdot|x(s)-y(s)| d s, \quad t \in[0, T) \tag{27}
\end{equation*}
$$

where $c=\lambda\left(2 x_{0}^{-1}\right)^{1-\lambda}$.
The Lipschitz-type formula (27) allows one to apply the iterative procedure presented, e.g., in Ref. 32, pp. 127-131, to conclude that the solutions $x, y$ coincide.

Lemma 3: Given the number $x_{0}>0$, there exists at least one continuous solution $x:[0,+\infty)$ $\rightarrow \mathbb{R}$ of Eq. (24) such that $x(0)=x_{0}$.

Proof: The argument is based on a technique developed by the authors in Ref. 6, Theorem 2.
First, set the numbers $p, q, L>0$, such that

$$
\frac{1}{p}+\frac{1}{q}=1, \quad 1<p<\min \left\{\frac{1}{\alpha}, \frac{1}{1-\alpha}\right\}
$$

and

$$
L>C(\alpha, p)^{q} \quad \text { for } C(\alpha, p)=p^{1 / q-\alpha} \cdot \Gamma(1-p(1-\alpha)) .
$$

Second, introduce the functions-here, $f(t, x)=H(t) x^{\lambda}$ -

$$
k(t)=1+x_{0}+\int_{0}^{t} \frac{\left|f\left(s, x_{0}\right)\right|}{(t-s)^{1-\alpha}} d s, \quad F(t)=c \cdot H(t)
$$

and

$$
H_{L}(t)=k(t) \cdot \exp \left(t+\frac{L}{q} \int_{0}^{t}[k(s) F(s)]^{q} d s\right), \quad t \geq 0
$$

where $c$ has been defined in the preceding proof.
Now, consider the complete metric space $\mathcal{X}=\left(X, d_{L}\right)$, where $X$ is the set of all elements of $C\left([0,+\infty),\left[x_{0},+\infty\right)\right)$ that behave as $O\left(H_{L}(t)\right)$ when $t$ goes to $+\infty$ and

$$
d_{L}(x, y)=\sup _{t \geq 0}\left\{\frac{|x(t)-y(t)|}{H_{L}(t)}\right\}, \quad x, y \in X
$$

Notice that, as in Lemma 2, the following inequality holds:

$$
|f(t, x(t))-f(t, y(t))| \leq F(t)|x(t)-y(t)|, \quad t \geq 0, \quad x, y \in X
$$

Introduce the operator $T: \mathcal{X} \rightarrow \mathcal{X}$ with the formula

$$
\mathcal{T}(x)(t)=x_{0}+\int_{0}^{t} \frac{f(s, x(s))}{(t-s)^{1-\alpha}} d s, \quad x \in X, t \geq 0
$$

Following verbatim the computations from Ref. 6, we conclude that $\mathcal{T}$ is a contraction of coefficient $C(\alpha, p) \cdot L^{-1 / q}$ which implies, via the Banach contraction principle, that our problem has a solution globally defined in the future.

Consider now the particular case of Eq. (24) given by

$$
H(t)=(t+1)^{-\alpha \lambda} h(t)=(t+1)^{-\alpha \lambda} \cdot t^{-(1+\varepsilon) / p_{2}}, \quad t \geq 1
$$

where $\varepsilon \in(0,1)$ is small enough to have $(1-\lambda) \alpha>(1+\varepsilon) / p_{2}$. The existence of $\varepsilon$ will be explained in the next paragraph. As regards the definition of $H$, this function is continued downward to $t$ $=0$ in such a way as to keep its sign and continuity in $[0,+\infty)$.

It is supposed that the number $p_{2}$ verifies the double inequality (22). To see how to choose $\varepsilon$, notice that the inequality below,

$$
(1-\lambda) \alpha>\lambda \alpha+(1-\lambda) \frac{\alpha^{2}}{1+\alpha}\left(>\frac{1}{p_{2}}\right)
$$

is equivalent with

$$
1>\lambda(2+\alpha)
$$

The hypotheses (25) lead to

$$
1>\frac{3 \alpha}{2}>\alpha+\frac{\alpha^{2}}{2}=\frac{\alpha}{2} \cdot(2+\alpha)>\lambda \cdot(2+\alpha)
$$

so such a number $\varepsilon$ exists truly.
Let now $x$ be the solution of Eq. (24) given by Lemmas $1-3$ for $T=+\infty$-such a solution exists always. By taking into account the formula of $H$, we obtain that

$$
\begin{align*}
x(t) & \geq x_{0}+\int_{1}^{t} \frac{h(s)}{(t-s)^{1-\alpha}} \cdot\left[\frac{x(s)}{(s+1)^{\alpha}}\right]^{\lambda} d s \geq x_{0}+\int_{1}^{t} \frac{s^{-(1+\varepsilon) / p_{3}}}{(t-s)^{1-\alpha}} \cdot\left[\frac{x_{0}}{(s+1)^{\alpha}}\right]^{\lambda} d s \\
& \geq \int_{1}^{t} \frac{d s}{(t-s)^{1-\alpha}} \cdot \frac{1}{t^{(1+\varepsilon) / p_{3}}} \cdot \frac{x_{0}^{\lambda}}{(t+1)^{\alpha \lambda}} \sim c \cdot t^{(1-\lambda) \alpha-(1+\varepsilon) / p_{3}} \quad \text { when } t \rightarrow+\infty \tag{28}
\end{align*}
$$

where $c=x_{0}^{\lambda} / \alpha$.
According to our theorem, $x(t)=o\left(t^{\alpha a}\right)$ for $t \rightarrow+\infty$, where $\lambda<1-a<\alpha$. Since we can fix $a$ in this range freely, assume that $a$ is taken such that

$$
\begin{equation*}
\alpha>2 \lambda>1-a>\lambda . \tag{29}
\end{equation*}
$$

This choice is in perfect agreement with the hypotheses (25).
The restriction (29) leads to-recall (22)-

$$
\frac{1}{p_{2}}>\lambda \alpha>(1-\lambda-a) \alpha
$$

Further, we have

$$
a \alpha>(1-\lambda) \alpha-\frac{1}{p_{2}}>(1-\lambda) \alpha-\frac{1+\varepsilon}{p_{2}}>0
$$

We would like to evaluate the difference between the exponent $a \alpha$ of $t$ in the asymptotic formula of the solution $x$ given by our theorem and the exponent $(1-\lambda) \alpha-(1+\varepsilon) / p_{2}$ provided by (28). This reads as

$$
\begin{aligned}
a \alpha-\left[(1-\lambda) \alpha-\frac{1+\varepsilon}{p_{2}}\right]= & \frac{1+\varepsilon}{p_{2}}+(a+\lambda-1) \alpha<(1+\varepsilon)\left[\lambda \alpha+(1-\lambda) \frac{\alpha^{2}}{1+\alpha}\right]+(a+\lambda-1) \alpha \\
= & \alpha\left\{[a+(2+\varepsilon) \lambda-1]+\frac{1+\varepsilon}{1+\alpha}[\alpha(1-\lambda)]\right\}<\alpha\{[a+(2+\varepsilon) \lambda-1]+(1 \\
& +\varepsilon) \alpha\}=\Omega \alpha
\end{aligned}
$$

Fix now $\eta_{0} \in(0,1)$ and replace the first of hypotheses (25) with a stronger condition, namely,

$$
\eta_{0}>\frac{7 \alpha}{2}
$$

Thus, we have

$$
\Omega<\eta_{0} a
$$

In fact, this inequality is equivalent with

$$
\left(1-\eta_{0}\right) a+(2+\varepsilon) \lambda+(1+\varepsilon) \alpha<1
$$

The latter inequality is valid as a consequence of the estimate below-recall the second of hypotheses (25)-

$$
\left(1-\eta_{0}\right) a+(2+\varepsilon) \lambda+(1+\varepsilon) \alpha<\left(1-\eta_{0}\right)+3 \lambda+2 \alpha<1-\eta_{0}+\frac{3 \alpha}{2}+2 \alpha<1
$$

The conclusion of this analysis of Eq. (24) is that the solution $x$, if computed numerically, will be at least as big as $O\left(t^{(1-\eta) \cdot a \alpha}\right)$ for $t \rightarrow+\infty$, where $\eta \in\left(0, \eta_{0}\right)$ and $\eta_{0}$ is a prescribed quantity, while the description of $x$ given by our theorem reads as $o\left(t^{a \alpha}\right)$ when $t \rightarrow+\infty$. In other words, the theorem
provides a sharp asymptotic estimate for the solutions of (1) and (3) and, in some particular cases, this estimate is optimal.

## ACKNOWLEDGMENTS

The work on this paper has been done during the visit of O.G.M. to Çankaya University in the fall of 2008. O.G.M. is grateful to the people from the Department of Mathematics and Computer Science for the friendly and enthusiastic working atmosphere. Both authors are grateful to the referees for their helpful comments.
${ }^{1}$ Agarwal, R. P., Djebali, S., Moussaoui, T., and Mustafa, O. G., "On the asymptotic integration of nonlinear differential equations," J. Comput. Appl. Math. 202, 352 (2007).
${ }^{2}$ Agrawal, O. P., "Formulation of Euler-Lagrange equations for fractional variational problems," J. Math. Anal. Appl. 272, 368 (2002).
${ }^{3}$ Ahlfors, L. V., Complex Analysis. An Introduction to the Theory of Analytic Functions of One Complex Variable (McGraw-Hill, New York, 1979).
${ }^{4}$ Atanackovic, T. M., Konjik, S., and Pipiloviç, S., "Variational problems with fractional derivatives: Euler-Lagrange equations," J. Phys. A 41, 095201 (2008).
${ }^{5}$ Băleanu, D. and Trujillo, J. J., "On exact solutions of a class of fractional Euler-Lagrange equations," Nonlinear Dyn. 52, 331 (2008).
${ }^{6}$ Băleanu, D. and Mustafa, O. G., "A note on global existence of solutions to fractional differential equations," Comput. Math. Appl. (in press).
${ }^{7}$ Bihari, I., "A generalization of a lemma of Bellman and its application to uniqueness problems of differential equations," Acta Math. Acad. Sci. Hung. 7, 81 (1956).
${ }^{8}$ Bihari, I., "Researches of the boundedness and stability of the solutions of nonlinear differential equations," Acta Math. Acad. Sci. Hung. 8, 261 (1957).
${ }^{9}$ Caputo, M., "Linear models of dissipation whose $Q$ is almost frequency indepedent II," Geophys. J. R. Astron. Soc. 13, 529 (1967).
${ }^{10}$ Chen, W., "An iterative learning observer for fault detection and accommodation in nonlinear time-delay systems," Int. J. Robust Nonlinear Control 16, 1 (2006).
${ }^{11}$ Delbosco, D. and Rodino, L., "Existence and uniqueness for a nonlinear fractional differential equation," J. Math. Anal. Appl. 204, 609 (1996).
${ }^{12}$ Deng, W., Li, C., and Lü, L., "Stability analysis of linear fractional differential system with multiple time delays," Nonlinear Dyn. 48, 409 (2007).
${ }^{13}$ Diethelm, K. and Ford, N. J., "Analysis of fractional differential equations," J. Math. Anal. Appl. 265, 229 (2002).
${ }^{14}$ Furati, K. M. and Tatar, N.-E., "Power-type estimates for a nonlinear fractional differential equation," Nonlinear Anal. Theory, Methods Appl. 62, 1025 (2005).
${ }^{15}$ Gilbarg, D. and Trudinger, N. S., Elliptic Partial Differential Equations of Second Order (Springer-Verlag, Berlin, 1998).
${ }^{16}$ Glöckle, W. G. and Nonnenmacher, T. F., "A fractional calculus approach to self-similar protein dynamics," Biophys. J. 68, 46 (1995).
${ }^{17}$ Gorenflo, R. and Mainardi, F., "Some recent advances in theory and simulation of fractional diffusion processes," J. Comput. Appl. Math. 229, 400 (2009).
${ }^{18}$ Heymans, N. and Podlubny, I., "Physical interpretation of initial conditions for fractional differential equations with Rieman-Liouville fractional derivatives," Rheol. Acta 45, 765 (2006).
${ }^{19}$ Hilfer, R., Applications of Fractional Calculus in Physics (World Scientific, Singapore, 2000).
${ }^{20}$ Kilbas, A. A., Srivastava, H. M., and Trujillo, J. J., Theory and Applications of Fractional Differential Equations, North-Holland Mathematical Studies Vol 204 (Elsevier, Amsterdam, 2006).
${ }^{21}$ Klimek, M., "Lagrangian and Hamiltonian fractional sequential mechanics," Czech. J. Phys. 52, 1247 (2002).
${ }^{22}$ Lakshmikantham, V., Leela, S., and Vasundhara Devi, J., Theory of Fractional Dynamic Systems (Cambridge Scientific, Cambridge, 2009).
${ }^{23}$ Magin, R. L., Fractional Calculus in Bioengineering (Begell House, Connecticut, 2006).
${ }^{24}$ Magin, R. L., Abdullah, O., Băleanu, D., and Xiaohong, J. Z., "Anomalous diffusion expressed through fractional order differential operators in the Bloch-Torrey equation," J. Magn. Reson. 190, 255 (2008).
${ }^{25}$ Metzler, R., Schick, W., Kilian, H. G., and Nonennmacher, T. F., "Relaxation in filled polymers: A fractional calculus approach," J. Chem. Phys. 103, 7180 (1995).
${ }^{26}$ Miller, K. S. and Ross, B., An Introduction to the Fractional Integrals and Derivatives. Theory and Applications (Wiley, New York, 1993).
${ }^{27}$ Momani, S. and Odibat, Z., "A novel method for nonlinear fractional partial differential equations: combination of DTM and generalized Taylor's formula," J. Comput. Appl. Math. 220, 85 (2008).
${ }^{28}$ Momani, S. and Odibat, Z., "Homotopy perturbation method for nonlinear partial differential equations of fractional order," Phys. Lett. A 365, 345 (2007).
${ }^{29}$ Muslih, S. I. and Băleanu, D., "Hamiltonian formulation of systems with linear velocities within Riemann-Liouville fractional derivatives," J. Math. Anal. Appl. 304, 599 (2005).
${ }^{30}$ Mustafa, O. G., "Existence of positive evanescent solutions to some quasilinear elliptic equations," Bull. Austral. Math. Soc. 78, 157 (2008).
${ }^{31}$ Oldham, K. B. and Spanier, J., The Fractional Calculus (Academic, New York, 1974).
${ }^{32}$ Podlubny, I., Fractional Differential Equations (Academic, San Diego, 1999).
${ }^{33}$ Sabatier, J., Agrawal, O. P., and Tenreiro Machado, J. A., Advances in Fractional Calculus. Theoretical Developments and Applications in Physics and Engineering (Springer-Verlag, Dordrecht, 2007).
${ }^{34}$ Samko, S. G., Kilbas, A. A., and Marichev, O. I., Fractional Integrals and Derivatives. Theory and Applications (Gordon and Breach, London, 1993).
${ }^{35}$ Uchaikin, V. V., The Method of the Fractional Derivatives (Art-i-Shock, Ulianovsk, 2008) (in Russian).
${ }^{36}$ West, B. J., Bologna, M., and Grigolini, P., Physics of Fractal Operators (Springer-Verlag, New York, 2003).
${ }^{37}$ Zaslavsky, G., Hamiltonian Chaos and Fractional Dynamics (Oxford University Press, Oxford, 2005).

Journal of Mathematical Physics is copyrighted by the American Institute of Physics (AIP). Redistribution of journal material is subject to the AIP online journal license and/or AIP copyright. For more information, see http://ojps.aip.org/jmp/jmpcr.jsp


[^0]:    ${ }^{\text {a) }}$ On leave from The National Institute for Laser, Plasma and Radiation, Physics, Institute of Space Sciences, MăgureleBucureşti, P.O. Box MG-23, R 76911, Romania. Electronic mail: dumitru@cankaya.edu.tr.
    ${ }^{\text {b) }}$ Electronic mail: octawian @yahoo.com.

