



Article

# Certain Results Comprising the Weighted Chebyshev Function Using Pathway Fractional Integrals

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**Abstract:** An analogous version of Chebyshev inequality, associated with the weighted function, has been established using the pathway fractional integral operators. The result is a generalization of the Chebyshev inequality in fractional integral operators. We deduce the left sided Riemann Liouville version and the Laplace version of the same identity. Our main deduction will provide noted results for an appropriate change to the Pathway fractional integral parameter and the degree of the fractional operator.

**Keywords:** Riemann Liouville fractional integral operator; pathway fractional order integral operator; Chebyshev functional

### 1. Introduction and Preliminaries

Fractional calculus has been broadly used in various disciplines of engineering and science as a mathematical model. Nowadays, fractional order (arbitrary order) calculus is more realistic than classical models of integer order for the memory description and hereditary rules of various physical problems and actions. Some important generalised fractional integral operators include the Hadamard operator, Erdélyi-Kober operators, the Saigo operator, the Gaussian hypergeometric operator, the Marichev–Saigo–Maeda fractional integral operator, etc.; out of those, the Riemann Liouville fractional integral operator has been widely used by researchers in theory as well as applications. For additional details about fractional calculus operators and their utility, one may refer to the treatises by Miller and Ross [1], Samko et al. [2], Kiryakova [3] and Baleanu et al. [4]. The pathway fractional integral operator is a generalised Riemann Liouville fractional integral operator in higher dimensions. Indeed, the pathway fractional integral operator has been used to define certain probability density functions and has interesting applications in statistics (also see [5–7]).

Let  $\lambda$  be an arbitrary real number preferably less than 1. Let a class of conformal functions:

$$C_{\lambda}^{(k)} = \{ f(t) = t^p \tilde{f}(t); p > \lambda, \tilde{f} \in C^{(k)}([\alpha, \beta]) \}$$

be a Banach space for our consideration for different non-negative values of k. The space we have chosen for k = 0 [3] is induced by the norm:

$$||f||_r = \left[\int_{\alpha}^{\beta} |f(t)|^r dt\right]^{\frac{1}{r}} < \infty.$$

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Suppose  $\phi(t) \in C_{\lambda}^{(k)}$  for k = 0 and  $\nu \in \mathbb{C}$ , which has a positive real part, and a is positive. The pathway fractional integral operator is defined as ([5]):

$$P_{0+}^{(\nu,\lambda)}(\phi(t)) = \int_0^{\frac{t}{a(1-\lambda)}} t^{\nu} \left[ 1 - \frac{a(1-\lambda)u}{t} \right]^{\nu/(1-\lambda)} \phi(u) du \tag{1}$$

We call  $\lambda$  the pathway parameter and  $\nu$  the order of the integral operator. While  $\lambda$  converges to 1 from the left side, the operator gets the form:

$$P_{0+}^{(\nu,1)}(\phi(t)) = \int_0^{t/a(1-\lambda)} t^{\nu} e^{\frac{-a\nu u}{t}} \phi(u) du = t^{\nu} \mathfrak{L}\left[\frac{a\nu}{t}\right]. \tag{2}$$

Meaning thereby that the fractional integral operator of pathway type reduces to the Laplace transform with the factor (av/t) for a particular value of  $\lambda$ .

The pathway fractional order integral operator transformed into a left-sided fractional order integral operator of Riemann–Liouville type, i.e.,  $I_{0+}^{\nu}(\phi(t))$  for a particular value of  $\lambda=0$  and a=1:

$$\int_0^t (t-u)^{\nu-1} \phi(u) du = \Gamma(\nu) I_{0+}^{\nu}(\phi(t)), \tag{3}$$

by replacing  $\nu$  to  $\nu-1$ . The reader may refer to the papers of Mathai and Haubold [6,7], Nair [8] and Nisar et al. [9] for more details on pathway operators.

Further, an operator T defined from X and  $\phi$ ,  $\psi \in X$ , called the weighted version of the Chebyshev functional, is:

$$T(\phi, \psi, p) = \int_{\alpha}^{\beta} p(t)dt \int_{\alpha}^{\beta} \phi(t)\psi(t)p(t)dt - \int_{\alpha}^{\beta} \phi(t)p(t)dt \int_{\alpha}^{\beta} \psi(t)p(t)dt, \tag{4}$$

where p(t) is the weighted function with positive values. A result has been established by Dragomir [10] related to Equation (4), given by:

$$2|T(\phi,\psi,p)| \le ||\phi'||_r ||\psi'||_s \left[ \int_{\alpha}^{\beta} \int_{\alpha}^{\beta} |u-v|p(u)p(v)dudv \right]. \tag{5}$$

Here,  $\phi$ ,  $\psi$  are derivable functions and  $\phi' \in L_r([\alpha, \beta])$ ,  $\psi' \in L_s([\alpha, \beta])$  with the condition that  $r^{-1} + s^{-1} = 1$ . Recently, by using fractional integral operators, several extensions of the classical inequalities, including Equation (5), have been studied by many authors, see [11–24] and the references therein. We attempt to generate the inequality of Equation (5) by making use of the fractional integral operator of pathway type. So, in this article our prime intention is to provide an analogous version of the Chebyshev inequality by means of pathway integral operators of fractional orders. The result is a generalization of the Chebyshev inequality in fractional integral operators. We also deduce the left-sided Riemann–Liouville version and Laplace version of the same inequality.

The article has been organised as follows: First we prove a generalised version of the vital result and then state a specific but useful theorem. In Section 3, we derive certain specific results using these two theorems in the form of corollaries.

# 2. Main Theorems

We state the subsequent results:

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**Theorem 1.** Let  $\phi(t), \psi(t) \in C_{\lambda}^{(k)}$  for k = 0 and  $\nu \in \mathbb{C}$ , such that the real part of  $\nu$  is positive and  $\phi' \in L_r([0,\infty))$  and  $\psi' \in L_s([0,\infty))$  have the condition  $r^{-1} + s^{-1} = 1$  (when r > 1). p is a positive function and we choose a > 0 and  $\lambda < 1$ . Therefore, for every  $t \in [0,\infty)$ :

$$\begin{split} |P_{0+}^{(\nu,\lambda)}(\phi(t)\psi(t)p(t))P_{0+}^{(\nu',\lambda')}(p(t)) - P_{0+}^{(\nu,\lambda)}(\psi(t)p(t))P_{0+}^{(\nu',\lambda')}(\phi(t)p(t)) \\ - P_{0+}^{(\nu,\lambda)}(\phi(t)p(t))P_{0+}^{(\nu',\lambda')}(\psi(t)p(t)) + P_{0+}^{(\nu,\lambda)}(p(t))P_{0+}^{(\nu',\lambda')}(\phi(t)\psi(t)p(t))| \\ &\leq ||\phi'||_r ||\psi'||_s t \int_0^{\frac{t}{a(1-\lambda)}} \int_0^{\frac{t}{a(1-\lambda')}} F(t,u)G(t,v)p(u)p(v)dvdu \\ &\leq ||\phi'||_r ||\psi'||_s t P_{0+}^{(\nu,\lambda)}(p(t))P_{0+}^{(\nu',\lambda')}(p(t)), \end{split} \tag{6}$$

where  $P_{0^+}^{(\nu,\lambda)}(\phi(t))$  and  $P_{0^+}^{(\nu',\lambda')}(\psi(t))$  are pathway operators in  $C_{\lambda}^{(k)}$  space and H(t,v), F(t,u) and G(t,v) are functions in two-dimensional space defined as follows:

$$H(u,v) = (\phi(u) - \phi(v))(\psi(u) - \psi(v)) \tag{7}$$

$$F(t,u) = t^{\nu} \left[ 1 - \frac{a(1-\lambda)u}{t} \right]^{\nu/(1-\lambda)}$$
(8)

where  $u \in [0, t]$  when  $t \in [0, \infty)$  and:

$$G(t,v) = t^{\nu'} \left[ 1 - \frac{a(1-\lambda')v}{t} \right]^{\nu'/(1-\lambda')}.$$
 (9)

**Proof.** Multiplying H(u,v), defined in Equation (7), by F(t,u)p(u) and integrating u over 0 to  $t/a(1-\lambda)$ :

$$\int_{0}^{\frac{t}{a(1-\lambda)}} H(t,v)F(t,u)p(u)du = \int_{0}^{t/a(1-\lambda)} t^{\nu} \left[ 1 - \frac{a(1-\lambda)u}{t} \right]^{\nu/(1-\lambda)} [\phi(u)\psi(u)p(u)du - \phi(v)\psi(u)p(u)du - \phi(u)\psi(v)p(u)du + \phi(v)\psi(v)p(u)du].$$
(10)

This implies:

$$\int_{0}^{\frac{t}{a(1-\lambda)}} H(t,v)F(t,u)p(u)du = P_{0+}^{(\nu,\lambda)}(\phi(t)\psi(t)p(t))$$

$$-\phi(v)P_{0+}^{(\nu,\lambda)}(\psi(t)p(t)) - \psi(v)P_{0+}^{(\nu,\lambda)}(\phi(t)p(t)) + \phi(v)\psi(v)P_{0+}^{(\nu,\lambda)}(p(t)).$$
(11)

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Now by multiplying Equation (11) by G(t,v)p(v) and integrating with respect to v between 0 and  $v'/a(1-\lambda')$ , we get:

$$\int_{0}^{\frac{t}{a(1-\lambda)}} \int_{0}^{\frac{t}{a(1-\lambda')}} H(u,v)F(t,u)G(t,v)p(u)p(v)dvdu$$

$$= \left(\int_{0}^{\frac{t}{a(1-\lambda)}} t^{\nu} \left[1 - \frac{a(1-\lambda)u}{t}\right]^{\nu/(1-\lambda)} \phi(u)\psi(u)p(u)du\right)$$

$$\times \left(\int_{0}^{\frac{t}{a(1-\lambda')}} t^{\nu'} \left[1 - \frac{a(1-\lambda')v}{t}\right]^{\nu'/(1-\lambda')} p(v)dv\right)$$

$$- \left(\int_{0}^{\frac{t}{a(1-\lambda)}} t^{\nu} \left[1 - \frac{a(1-\lambda)u}{t}\right]^{\nu/(1-\lambda)} \psi(u)p(u)du\right)$$

$$\times \left(\int_{0}^{\frac{t}{a(1-\lambda)}} t^{\nu} \left[1 - \frac{a(1-\lambda')v}{t}\right]^{\nu'/(1-\lambda')} \phi(v)p(v)dv\right)$$

$$- \left(\int_{0}^{\frac{t}{a(1-\lambda)}} t^{\nu} \left[1 - \frac{a(1-\lambda)u}{t}\right]^{\nu/(1-\lambda)} \phi(u)p(u)du\right)$$

$$\times \left(\int_{0}^{\frac{t}{a(1-\lambda)}} t^{\nu} \left[1 - \frac{a(1-\lambda')v}{t}\right]^{\nu'/(1-\lambda')} \psi(v)p(v)dv\right)$$

$$+ \left(\int_{0}^{\frac{t}{a(1-\lambda)}} t^{\nu} \left[1 - \frac{a(1-\lambda)u}{t}\right]^{\nu/(1-\lambda')} \phi(v)\psi(v)p(v)dv\right)$$

$$\times \left(\int_{0}^{\frac{t}{a(1-\lambda)}} t^{\nu} \left[1 - \frac{a(1-\lambda)u}{t}\right]^{\nu/(1-\lambda')} \phi(v)\psi(v)p(v)dv\right)$$

This implies:

$$\int_{0}^{\frac{t}{a(1-\lambda)}} \int_{0}^{\frac{t}{a(1-\lambda')}} H(u,v) F(t,u) G(t,v) p(u) p(v) dv du = P_{0+}^{(\nu,\lambda)} (\phi(t)\psi(t)p(t)) 
\times P_{0+}^{(\nu',\lambda')} (p(t)) - P_{0+}^{(\nu,\lambda)} (\psi(t)p(t)) P_{0+}^{(\nu',\lambda')} (\phi(t)p(t)) 
- P_{0+}^{(\nu,\lambda)} (\phi(t)p(t)) P_{0+}^{(\nu',\lambda')} (\psi(t)p(t)) + P_{0+}^{(\nu,\lambda)} (p(t)) P_{0+}^{(\nu',\lambda')} (\phi(t)\psi(t)p(t)).$$
(13)

Now, by the Hölder inequality:

$$|H(u,v)| \le |u-v| \left| \int_{u}^{v} |\phi'|^{r} dx \right|^{r-1} \left| \int_{u}^{v} |\psi'|^{s} dt \right|^{s-1}, \tag{14}$$

with  $r^{-1} + s^{-1} = 1$ . We apply the Hölder inequality on the left side of Equation (13). The left side of Equation (13) converts to:

$$\int_{0}^{\frac{t}{a(1-\lambda)}} \int_{0}^{\frac{t}{a(1-\lambda')}} H(u,v)F(t,u)G(t,v)p(u)p(v)dvdu \leq \int_{0}^{\frac{t}{a(1-\lambda)}} \int_{0}^{\frac{t}{a(1-\lambda')}} F(t,u)G(t,v)p(u)p(v) \\
\times |u-v| \left| \int_{u}^{v} |\phi'|^{r} dx \right|^{r-1} \left| \int_{u}^{v} |\psi'|^{s} dt \right|^{s-1} dvdu. \tag{15}$$

Now, we will use the Hölder inequality, but in the form of:

$$\left| \int_{u}^{v} \int_{u}^{v} f(x)g(t)dxdt \right| \le \left| \int_{u}^{v} \int_{u}^{v} |f(x)|^{r} dxdt \right|^{r-1} \left| \int_{u}^{v} \int_{u}^{v} |g(t)|^{s} dxdt \right|^{s-1}$$
(16)

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to get:

$$\int_{0}^{\frac{t}{a(1-\lambda)}} \int_{0}^{\frac{t}{a(1-\lambda')}} H(u,v)F(t,u)G(t,v)p(u)p(v)dvdu$$

$$\leq \left[\int_{0}^{\frac{t}{a(1-\lambda)}} \int_{0}^{\frac{t}{a(1-\lambda')}} F(t,u)G(t,v)p(u)p(v)|u-v| \left|\int_{u}^{v} |\phi'(x)|^{r} dx \right| dudv\right]^{r-1}$$

$$\times \left[\int_{0}^{\frac{t}{a(1-\lambda)}} \int_{0}^{\frac{t}{a(1-\lambda')}} F(t,u)G(t,v)p(u)p(v)|u-v| \left|\int_{u}^{v} |\psi'(t)|^{s} dt \right| dudv\right]^{s-1}$$

$$= \left[||\phi'||_{r}^{r} \int_{0}^{\frac{t}{a(1-\lambda)}} \int_{0}^{\frac{t}{a(1-\lambda')}} F(t,u)G(t,v)p(u)p(v)|u-v|dudv\right]^{r-1}$$

$$\times \left[||\psi'||_{s}^{s} \int_{0}^{\frac{t}{a(1-\lambda)}} \int_{0}^{\frac{t}{a(1-\lambda')}} F(t,u)G(t,v)p(u)p(v)|u-v|dudv\right]^{s-1}$$

$$\leq ||\phi'||_{r}||\psi'||_{s} \left|\int_{0}^{\frac{t}{a(1-\lambda)}} \int_{0}^{\frac{t}{a(1-\lambda')}} F(t,u)G(t,v)p(u)p(v)|u-v|dvdu\right|$$

We look at this inequality:

$$|P_{0+}^{(\nu,\lambda)}(\phi(t)\psi(t)p(t))P_{0+}^{(\nu',\lambda')}(p(t)) - P_{0+}^{(\nu,\lambda)}(\psi(t)p(t))P_{0+}^{(\nu',\lambda')}(\phi(t)p(t)) - P_{0+}^{(\nu',\lambda')}(\phi(t)p(t))P_{0+}^{(\nu',\lambda')}(\phi(t)p(t)) + P_{0+}^{(\nu,\lambda)}(p(t))P_{0+}^{(\nu',\lambda')}(\phi(t)\psi(t)p(t))|$$

$$\leq ||\phi'||_r ||\psi'||_s \int_0^{\frac{t}{a(1-\lambda)}} \int_0^{\frac{t}{a(1-\lambda')}} F(t,u)G(t,v)p(u)p(v)|u - v|dvdu$$

$$(18)$$

since |u - v| < t:

$$|P_{0+}^{(\nu,\lambda)}(\phi(t)\psi(t)p(t))P_{0+}^{(\nu',\lambda')}(p(t)) - P_{0+}^{(\nu,\lambda)}(\psi(t)p(t))P_{0+}^{(\nu',\lambda')}(\phi(t)p(t)) - P_{0+}^{(\nu,\lambda)}(\phi(t)p(t))P_{0+}^{(\nu',\lambda')}(\phi(t)p(t)) + P_{0+}^{(\nu,\lambda)}(p(t))P_{0+}^{(\nu',\lambda')}(\phi(t)\psi(t)p(t))|$$

$$\leq ||\phi'||_r ||\psi'||_s t \int_0^{\frac{t}{a(1-\lambda)}} \int_0^{\frac{t}{a(1-\lambda')}} F(t,u)G(t,v)p(u)p(v)dvdu$$

$$\leq ||\phi'||_r ||\psi'||_s t P_{0+}^{(\nu,\lambda)}(p(t))P_{0+}^{(\nu',\lambda')}(p(t))$$

$$(19)$$

This proves the result of Equation (6).  $\Box$ 

A specific case of Theorem 1 can be deduced by assuming both pathway fractional integral operators with the same parameters.

**Theorem 2.** Suppose  $\phi(t)$ ,  $\psi(t) \in C_{\lambda}^{(k)}$  for k = 0 and  $v \in \mathbb{C}$ , such that the real part of v is positive and  $\phi' \in L_r([0,\infty))$  and  $\psi' \in L_s([0,\infty))$  with the condition  $r^{-1} + s^{-1} = 1$  (when r > 1). p is a positive function and we are choosing a > 0 and a < 1. Thereupon for the entire  $a \in [0,\infty)$ :

$$\begin{split} |P_{0^{+}}^{(\nu,\lambda)}(\phi(t)\psi(t)p(t))P_{0^{+}}^{(\nu,\lambda)}(p(t)) - P_{0^{+}}^{(\nu,\lambda)}(\psi(t)p(t)).P_{0^{+}}^{(\nu,\lambda)}(\phi(t)p(t))| \\ &\leq \frac{t}{2}||\phi'||_{r}||\psi'||_{s} \int_{0}^{\frac{t}{a(1-\lambda)}} \int_{0}^{\frac{t}{a(1-\lambda)}} F(t,u)G(t,v)p(u)p(v)dvdu \\ &\leq \frac{t}{2}||\phi'||_{r}||\psi'||_{s} \left(P_{0^{+}}^{(\nu,\lambda)}(p(t))\right)^{2}, \end{split} \tag{20}$$

where  $P_{0^+}^{(\nu,\lambda)}(\phi(t))$  is a pathway operator in  $C_{\lambda}^{(k)}$  space and H(t,v), F(t,u) and G(t,v) are functions in two dimensional space defined in Theorem 1.

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In the next section, we deduce some interesting and established (or new) results from Theorem 1.

#### 3. Derived Results

# 3.1. Choose $\lambda = 0$ and a = 1

Theorem 2 will become the subsequent known inequality due to Dahmani et al. [19]:

$$|I_{0+}^{(\nu)}(\phi(t)\psi(t)p(t))I_{0+}^{(\nu)}(p(t)) - I_{0+}^{(\nu)}(\psi(t)p(t))I_{0+}^{(\nu)}(\phi(t)p(t))| \\ \leq \frac{t}{2(\Gamma(\nu))^2}||\phi'||_r||\psi'||_s \int_0^t \int_0^t F(t,u)G(t,v)p(u)p(v)dvdu$$
(21)

or:

$$|I_{0+}^{(\nu)}(\phi(t)\psi(t)p(t))I_{0+}^{(\nu)}(p(t)) - I_{0+}^{(\nu)}(\psi(t)p(t))I_{0+}^{(\nu)}(\phi(t)p(t))| \\ \leq \frac{t}{2}||\phi'||_r||\psi'||_s \left(I_{0+}^{(\nu)}(p(t))\right)^2.$$
(22)

Here  $I_{0^+}^{(\nu)}(\phi(t))$  is a left-sided fractional integral operator of Riemann–Liouville type.

Here we note that by taking a particular value of  $\nu$  as  $\nu=1$ , we can get the Chebyshev inequality:

$$|I_{0+}^{(1)}(\phi(t)\psi(t)p(t))I_{0+}^{(1)}(p(t)) - I_{0+}^{(1)}(\psi(t)p(t))I_{0+}^{(1)}(\phi(t)p(t))| \\ \leq \frac{t}{2}||\phi'||_r||\psi'||_s \left(I_{0+}^{(1)}(p(t))\right)^2,$$
(23)

that is:

$$|\int_{0}^{t} \phi(u)\psi(u)p(u)du \int_{0}^{t} p(u)du - \int_{0}^{t} \psi(u)p(u)du \int_{0}^{t} \phi(u)p(u)du| \\ \leq \frac{t}{2}||\phi'||_{r}||\psi'||_{s} \left(\int_{0}^{t} p(u)du\right)^{2}, \tag{24}$$

for unweighted space:

$$|\int_{0}^{t} \phi(u)\psi(u)du - \int_{0}^{t} \psi(u)du \int_{0}^{t} \phi(u)du| \le \frac{t}{2}||\phi'||_{r}||\psi'||_{s}.$$
 (25)

# 3.2. $\lambda$ Tends to 1 from the Left Side

From Theorem 2:

$$|\mathcal{L}_{fgp}(\xi)\mathcal{L}_{p}(\xi) - \mathcal{L}_{gp}(\xi)\mathcal{L}_{fp}(\xi)| \leq \frac{1}{2t^{2\nu-1}}||\phi'||_{r}||\psi'||_{s}$$

$$\times \int_{0}^{\infty} \int_{0}^{\infty} F(t,u)G(t,v)p(u)p(v)dvdu$$
(26)

or:

$$|\mathcal{L}_{fgp}(\xi)\mathcal{L}_{p}(\xi) - \mathcal{L}_{gp}(\xi)\mathcal{L}_{fp}(\xi)| \le \frac{t}{2}||\phi'||_{r}||\psi'||_{s} \left(\mathcal{L}_{fgp}(\xi)\right)^{2},\tag{27}$$

where  $\xi = av/t$  and  $\mathcal{L}_f(\xi)$  are Laplace of the function f.

Since p(t) is a weighted function, we can reduce Theorem 2 in unweighted form, which is given as follows:

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**Corollary 1.** Suppose  $\phi(t)$ ,  $\psi(t) \in C_{\lambda}^{(k)}$  for k = 0 and  $\nu \in \mathbb{C}$ , such that the real part of  $\nu$  is positive and  $\phi' \in L_r([0,\infty))$  and  $\psi' \in L_s([0,\infty))$  with the condition  $r^{-1} + s^{-1} = 1$  (when r > 1). Choose a > 0 and  $\lambda < 1$ . Then for whole  $t \in [0,\infty)$ :

$$|\mathcal{A}t^{\nu+1}P_{0+}^{(\nu,\lambda)}(\phi(t)\psi(t)) - P_{0+}^{(\nu,\lambda)}(\psi(t))P_{0+}^{(\nu,\lambda)}(\phi(t))|$$

$$\leq \frac{t}{2}||\phi'||_{r}||\psi'||_{s} \int_{0}^{\frac{t}{a(1-\lambda)}} \int_{0}^{\frac{t}{a(1-\lambda)}} F(t,u)G(t,v)dvdu$$

$$\leq \frac{\mathcal{A}^{2}t^{2\nu+3}}{2}||\phi'||_{r}||\psi'||_{s},$$
(28)

where:

$$\mathcal{A}(\nu,\lambda) = \frac{\Gamma\left(1 + \frac{\nu}{1-\lambda}\right)}{a(1-\lambda)\Gamma\left(2 + \frac{\nu}{1-\lambda}\right)},$$

which depends upon the parameter of the pathway operator.

In similar fashion, the inequality in Theorem 1 becomes:

$$|\mathcal{B}t^{\nu'+1}P_{0+}^{(\nu,\lambda)}(\phi(t)\psi(t)) - P_{0+}^{(\nu,\lambda)}(\psi(t))P_{0+}^{(\nu',\lambda')}(\phi(t)) - P_{0+}^{(\nu,\lambda)}(\phi(t))P_{0+}^{(\nu',\lambda')}(\psi(t)) + \mathcal{A}t^{(\nu+1)}P_{0+}^{(\nu',\lambda')}(\phi(t)\psi(t))| \leq ||\phi'||_r ||\psi'||_s t \int_0^{\frac{t}{a(1-\lambda)}} \int_0^{\frac{t}{a(1-\lambda')}} F(t,u)G(t,v)dvdu$$

$$\leq \mathcal{A}\mathcal{B}t^{\nu+\nu'+3}||\phi'||_r ||\psi'||_s,$$
(29)

where:

$$\mathcal{B}(\nu,\lambda) = \frac{\Gamma\left(1 + \frac{\nu'}{1 - \lambda'}\right)}{a(1 - \lambda')\Gamma\left(2 + \frac{\nu'}{1 - \lambda'}\right)}.$$

Some more results can also be established by choosing  $p(t) = t^{\gamma - 1}$ .

**Corollary 2.**  $\phi(t), \psi(t)$  and  $\nu$  are the same as mentioned in Theorem 2; then for all  $t \in [0, \infty)$ :

$$|\mathcal{C}t^{\nu+\gamma}P_{0+}^{(\nu,\lambda)}(t^{\gamma-1}\phi(t)\psi(t)) - P_{0+}^{(\nu,\lambda)}(t^{\gamma-1}\psi(t))P_{0+}^{(\nu,\lambda)}(t^{\gamma-1}\phi(t))|$$

$$\leq \frac{t}{2}||\phi'||_{r}||\psi'||_{s} \int_{0}^{\frac{t}{a(1-\lambda)}} \int_{0}^{\frac{t}{a(1-\lambda)}} u^{\gamma-1}v^{\gamma-1}F(t,u)G(t,v)dvdu$$

$$\leq \frac{\mathcal{C}^{2}t^{2\gamma-1}}{2}||\phi'||_{r}||\psi'||_{s},$$
(30)

where:

$$C(\nu,\lambda,\gamma) = \frac{(\Gamma(\gamma))\Gamma\left(1+\frac{\nu}{1-\lambda}\right)}{[a(1-\lambda)]^{\gamma}\Gamma\left(1+\gamma+\frac{\nu}{1-\lambda}\right)},$$

which depends upon the parameter of the pathway operator as well as  $\gamma$ .

Similarly, the inequality of Theorem 1 reduces to:

$$|\mathcal{D}t^{\nu+\gamma}P_{0+}^{(\nu,\lambda)}(t^{\gamma-1}\phi(t)\psi(t)) - P_{0+}^{(\nu,\lambda)}(t^{\gamma-1}\psi(t))P_{0+}^{(\nu',\lambda')}(t^{\gamma-1}\phi(t)) - P_{0+}^{(\nu,\lambda)}(t^{\gamma-1}\phi(t))P_{0+}^{(\nu',\lambda')}(t^{\gamma-1}\phi(t))P_{0+}^{(\nu',\lambda')}(t^{\gamma-1}\phi(t)\psi(t))| \\ \leq ||\phi'||_r ||\psi'||_s t \int_0^{\frac{t}{a(1-\lambda)}} \int_0^{\frac{t}{a(1-\lambda')}} u^{\gamma-1}v^{\gamma-1}F(t,u)G(t,v)dvdu \\ \leq \mathcal{C}\mathcal{D}t^{\nu+\nu'+2\gamma+1}||\phi'||_r ||\psi'||_s,$$
(31)

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where:

$$\mathcal{D}(\nu',\lambda',\gamma) = \frac{(\Gamma(\gamma))\Gamma\left(1 + \frac{\nu'}{1-\lambda'}\right)}{[a(1-\lambda')]^{\gamma}\Gamma\left(1 + \gamma + \frac{\nu'}{1-\lambda'}\right)},$$

and  $t \in [0, \infty)$ .

### 4. Discussion and Conclusions

We have proved two theorems in this paper, one involving pathway fractional integral operators with different parameters and the other pathway fractional integral operators with the same parameters. These two theorems are analogous versions of the Chebyshev inequality. Further, we have taken some particular cases of these theorems, choosing specific values of the parameters. Since the pathway operator is a unification of the distinct nature of operators, we can discover a number of inequalities by selecting the values applicable to the restrictions and the weighted function p(t). Two of them have been shown here, viz., the Laplace operator and the left-sided Riemann–Liouville operator. Following similar methodology, we can also generalize other inequalities of the literature using pathway fractional integral operators.

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