

# Analytical optical soliton solutions of the Schrödinger-Poisson dynamical system

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## ABSTRACT

The article studies the exact traveling wave solutions to the Schrödinger-Poisson system which has applications in gravity's role of quantum state and approximate the coupling between quantum mechanics with gravitation. Diverse exact solutions in hyperbolic, trigonometric and plane wave forms are obtained using two norms of integration. For this sake modified extended direct algebraic (MEDA) and ( $G'/G$ )-expansion techniques are used. The 3D plots and their corresponding contour graphs are also depicted. The constraints conditions for the exact of solutions are also emerged during the derivation of solution.

## Introduction

Nonlinear partial differential equations (NPDEs) have been comprehensively studied in recent years. These equations play fundamental role in the modeling of various fields of science and engineering such as solid state physics, thermodynamics, civil engineering, soil mechanics, economics and quantum physics [1–7]. Many nonlinear systems are also emerged in coupled form, and these systems have great important impact in different fields of sciences [1].

Thus in this article, the Schrödinger-Poisson system is under investigation. This system is fundamentally the Schrödinger equation together with the gravitational potential, which has application of gravity's role in quantum state [2]. Also this nonlinear system can approximate the coupling between quantum mechanics with gravitation [3]. Further this system can also be viewed as the junction of Einstein-Klein Gordon and Einstein-Dirac system [4]. MEDA and ( $G'/G$ )-expansion methods [8–11] are applied to find the exact solutions to this system to find the exact solutions, has impressive impact in the theory of solitons [12–15]. These types of solutions play multiple significant role in the appropriate understanding of qualitative features of numerous phenomena and processes in different areas of natural science. The

Schrödinger-Poisson system is read as

$$\begin{cases} -i \frac{\partial \psi}{\partial t} = -\Delta \psi + \phi(x)\psi \\ -\Delta \phi = |\psi|^2 \end{cases} \quad (1)$$

The wave profile function is defined by  $\psi(x, t)$ . Here, function  $\phi$  is additionally has an elliptic behavior *i.e.*, a solution of Schrödinger-Poisson system. In this case where  $\phi(x)$  is determine by charge of the wave function itself. Using the definition of  $\Delta$ , the last system can be converted into single equation which has the following form,

$$-i \frac{\partial^3 \psi}{\partial x^2 \partial t} + \frac{\partial^4 \psi}{\partial x^4} + 2\phi'(x) \frac{\partial \psi}{\partial x} - \phi(x) \frac{\partial^2 \psi}{\partial x^2} + |\psi|^2 \psi = 0. \quad (2)$$

Much concentration has been paid to the models of NLPDEs for finding exact solutions in last few years. There are many powerful methods have been constructed and developed to analytically solve NPDEs with the aid of computational software such as Maple, Mathematica and Matlab [5]. The efficiency and reliability of such software is much better. In recent decades, exact solutions [6], analytical solutions [7] and numerical

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solutions [8] of many NPDEs have been successfully obtained [9]. The methods for obtaining exact explicit solutions of NPDEs are the tanh-function method [10], the exp-function method [11], the F-expansion method [16], Hirota method [17], Kudryashov method [9], the MEDA method [18], the extended auxiliary equation method [19] and modified method of simplest equation, the  $(G'/G^2)$ -expansion method [20], modified mapping method [21] and extended Fan-sub equation method [22]. Examples of the methods for solving NPDEs numerically are the finite element method [23], finite volume method [24], generalized finite difference method [25], spectral collocation method [26] and Galerkin finite element method [27];  $(G'/G), 1/G$ -expansion method [28]; the unified method [29]; simplified Hirota's method [30]; the exp  $(-\Phi(\xi))$  method [31]. Hence, there are results containing the extended and modified direct algebraic method, extended mapping method, and Seadawy technique of valuable researches in which may well complement the existing literature such as [32–39] In the following section, the model is a investigated analytically.

**Analysis of Schrödinger-Poisson system**

In this section, we adopt two integration norms namely MEDA [40] and  $(G'/G)$ -expansion method [41].

Using MEDA-method

Let us find the new exact traveling wave solutions to the system Eq. (1). For this, we convert Eq. (2) into nonlinear ordinary differential equation using the following complex transformation.

$$\psi(x, t) = U(z) \times e^{ip}, \tag{3}$$

where  $\psi(x, t)$  is the wave amplitude of wave profile,  $z = x - vt$ , and  $x$  is independent spatial variable and  $t$  represents the temporal variable,  $v$  is the velocity of the wave. Also note that  $p = p_1x + p_2t$ , where  $p_1$  and  $p_2$  are arbitrary constants. Hence substituting the results of Eq. (3) into Eq. (1) can get the equation, then separating the real and imaginary parts yields a pair of relations. The imaginary part gives the following constraint condition.

$$\rho_1 v = (2\phi(x) + 2p_2 + p_1^2) \tag{4}$$

and real parts take the form,

$$U^{iv}(z) + (2vp_1 - 6p_1^2 - p_2 - \phi(x))U''(z) - 2(\phi'(x)U'(z) + (p_1^4 + p_2p_1^2 + \phi(x)p_1^2)U(z) + U^3(z)) = 0. \tag{5}$$

The value of the positive integer  $N$  can be determined using the homogeneous balance principle, i.e, by balancing between the highest order derivatives and the nonlinear terms occurring in Eq. (5).

$$N + 4 = 3N \Rightarrow N = 2.$$

$$16p_1\gamma^2v = 48p_1^2\gamma^2 + (\phi'(x))^2 + 8\phi(x)\gamma^2 - 64\gamma^3 + 8p_2\gamma^2 - \sqrt{512(\phi'(x))^2\gamma^3 + (\phi'(x))^4}$$

Thus, Eq. (5) the formal solutions are obtained using the following result.

$$U(z) = \beta_0 + \beta_1\eta(z) + \beta_2\eta^2(z) + \gamma_1\eta^{-1}(z) + \gamma_2\eta^{-2}(z), \tag{6}$$

where  $\beta_0, \beta_1, \beta_2, \gamma_1, \gamma_2$  are unknown constants, and  $\eta'(z) = \gamma + \eta^2(z)$  in which  $\gamma$  is a free-parameter. Substituting Eq. (6) into Eq. (5) with the aid

of  $\eta'(z) = \gamma + \eta^2(z)$ , by collecting all the terms with the same power of  $\eta(z)$  [42]. After equating each coefficient equal to zero, this yields a set of following algebraic equations. Solving this system of algebraic equations with the aid of maple, we obtained following set of solutions.

For Case-1: The  $\beta_2, \gamma_2, v, \gamma$  are free parameters while along with the value of  $\beta_0 = 0, \beta_1 = 0, \gamma_1 = 0, \beta_2 = \pm 2\sqrt{-30}, \gamma_2 = \pm 2\gamma^2\sqrt{-30}, v = \frac{-40\gamma + 6p_1^2 + \phi(x) + p_2}{2p_1}$ . The constraint condition hold for the following results, when  $v$  is putting in it,  $\gamma = -\frac{3}{40}\phi(x) - \frac{3}{40}p_2 + \frac{1}{10}p_1^2$ . For  $\gamma < 0$ , thus following families of hyperbolic solutions are obtained.

$$\psi_1 = \left( \pm 2, \sqrt{-30} (-\sqrt{-\gamma} \tanh(\sqrt{-\gamma}z))^2 \pm 2\gamma^2\sqrt{-30} (-\sqrt{-\gamma} \tanh(\sqrt{-\gamma}z))^{-2} \right) e^{ip},$$

or

$$\psi_2 = \left( \pm 2\sqrt{-30} (-\sqrt{-\gamma} \coth(\sqrt{-\gamma}z))^2 \pm 2\gamma^2\sqrt{-30} (-\sqrt{-\gamma} \coth(\sqrt{-\gamma}z))^{-2} \right) e^{ip}.$$

For  $\gamma > 0$ , thus following periodic wave solutions are obtained.

$$\psi_3 = \left( \pm 2\sqrt{-30} (\sqrt{\gamma} \tan(\sqrt{\gamma}z))^2 \pm 2\gamma^2\sqrt{-30} (\sqrt{\gamma} \tan(\sqrt{\gamma}z))^{-2} \right) e^{ip},$$

or

$$\psi_4 = \left( \pm 2\sqrt{-30} (-\sqrt{\gamma} \cot(\sqrt{\gamma}z))^2 \pm 2\gamma^2\sqrt{-30} (-\sqrt{\gamma} \cot(\sqrt{\gamma}z))^{-2} \right) e^{ip}.$$

For  $\gamma = 0$ , the following plan solutions are obtained.

$$\psi_5 = \left( \pm 2\sqrt{-30} \pm 2\gamma^2\sqrt{-30}z^{-2} \right) e^{ip}$$

The graph and its corresponding contour plot of the solutions  $\psi_1(x, t)$  are shown in Fig. 1, for the values of parameters  $\rho_1 = 0.5, \rho_2 = 0.008, \gamma = 0.0005, \phi(x) = \sin(x)$  and  $v = 0.0009$ .

The graph and its corresponding contour plot of the solutions  $\psi_2(x, t)$  are shown in Fig. 2, for the values of parameters  $\rho_1 = 0.5, \rho_2 = 0.008, \gamma = 0.0005, \phi(x) = \sin(x)$  and  $v = 0.0009$ .

The graph and its corresponding contour plot of the solutions  $\psi_4(x, t)$  are shown in Fig. 3, for the values of parameters  $\rho_1 = 0.5, \rho_2 = 0.008, \gamma = 0.0005, \phi(x) = \sin(x)$  and  $v = 0.0009$ .

The graph and its corresponding contour plot of the solutions  $\psi_5(x, t)$  are shown in Fig. 4, for the values of parameters  $\rho_1 = 0.5, \rho_2 = 0.008, \gamma = 0.0005, \phi(x) = \sin(x)$  and  $v = 0.0009$ .

For Case-2: The  $\gamma_1, \gamma_2, \gamma, v$  are free parameters while along with  $\beta_0 = 0, \beta_1 = 0, \beta_2 = 0, \gamma_1 = \pm \frac{1}{4}\sqrt{-2} \left( \phi'(x) - \sqrt{512\gamma^3 + (\phi'(x))^2} \right), \gamma_2 = \pm 2\gamma^2\sqrt{-30}$ .

The constraint condition holds for the following results, when  $v$  is putting in it,  $24\gamma^2\phi(x) = -24p_2\gamma^2 + 32p_1^2\gamma^2 + ((\phi'(x))^2 - 64\gamma^3 - \sqrt{512((\phi'(x))^2\gamma^3 + ((\phi'(x))^4)})$ , For  $\gamma < 0$ , thus following families of hyperbolic solutions are obtained.

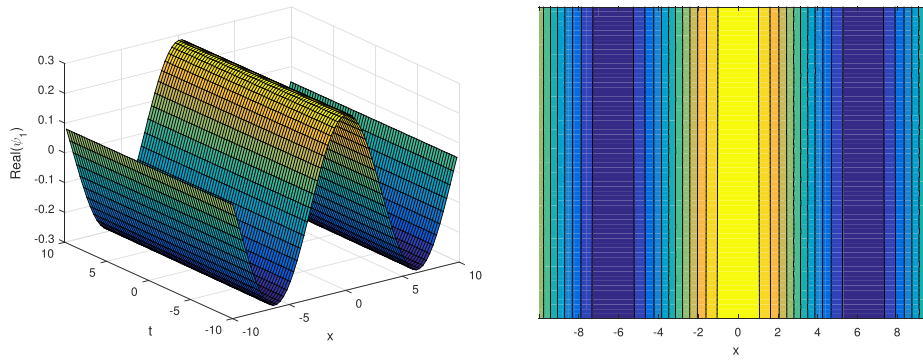


Fig. 1. The 3D plot and the corresponding contour representation of the traveling wave solutions  $\psi_1(x, t)$  for different values of parameters.

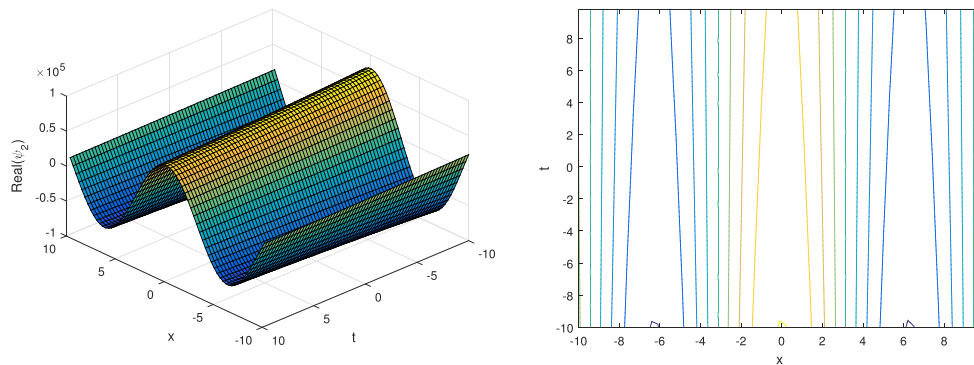


Fig. 2. The 3D plot and the corresponding contour representation of the traveling wave solutions  $\psi_2(x, t)$  for different values of parameters.

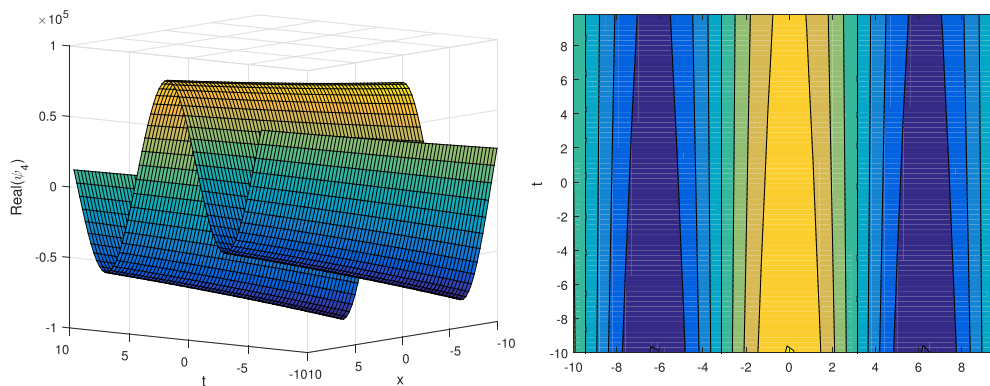


Fig. 3. The 3D plot and the corresponding contour representation of the traveling wave solutions  $\psi_4(x, t)$  for different values of parameters.

$$\psi_6 = \left( \left( \pm \frac{\sqrt{-2}\phi'(x)}{4} - \sqrt{512\gamma^3 + (\phi'(x))^2} \right) (-\sqrt{-\gamma} \tanh(-\sqrt{-\gamma}z))^{-1} \pm 2\gamma^2\sqrt{-30} (-\sqrt{-\gamma} \tanh(-\sqrt{-\gamma}z))^{-2} \right) e^{ip},$$

or

$$\psi_7 = \left( \left( \pm \frac{\sqrt{-2}\phi'(x)}{4} - \sqrt{512\gamma^3 + (\phi'(x))^2} \right) (-\sqrt{-\gamma} \coth(-\sqrt{-\gamma}z))^{-1} \pm 2\gamma^2\sqrt{-30} (-\sqrt{-\gamma} \coth(-\sqrt{-\gamma}z))^{-2} \right) e^{ip}.$$

For  $\gamma > 0$ , thus following periodic wave solutions are obtained.

$$\psi_8 = \left( \left( \pm \frac{\sqrt{-2}\phi'(x)}{4} - \sqrt{512\gamma^3 + (\phi'(x))^2} \right) (\sqrt{\gamma} \tan(\sqrt{\gamma}z))^{-1} \pm 2\gamma^2\sqrt{-30} (\sqrt{\gamma} \tan(\sqrt{\gamma}z))^{-2} \right) e^{ip},$$

or

$$\psi_9 = \left( \left( \pm \frac{\sqrt{-2}\phi'(x)}{4} - \sqrt{512\gamma^3 + (\phi'(x))^2} \right) (-\sqrt{\gamma} \cot(\sqrt{\gamma}z))^{-1} \pm 2\gamma^2\sqrt{-30} (-\sqrt{\gamma} \cot(\sqrt{\gamma}z))^{-2} \right) e^{ip}.$$

For  $\gamma = 0$ , thus following plane solutions are obtained.

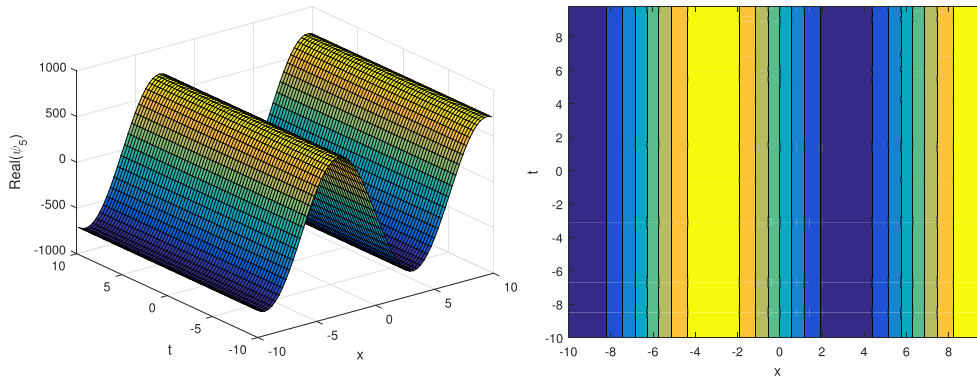


Fig. 4. The 3D plot and the corresponding contour representation of the traveling wave solutions  $\psi_5(x, t)$  for different values of parameters.

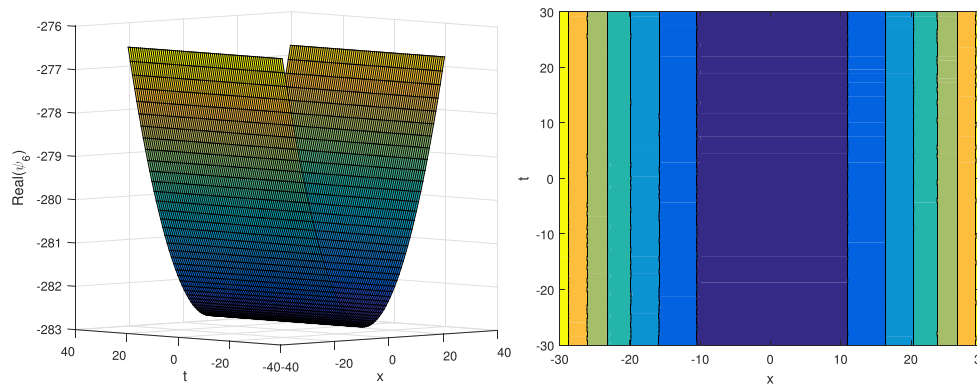


Fig. 5. The 3D plot and the corresponding contour representation of the traveling wave solutions  $\psi_6(x, t)$  for different values of parameters.

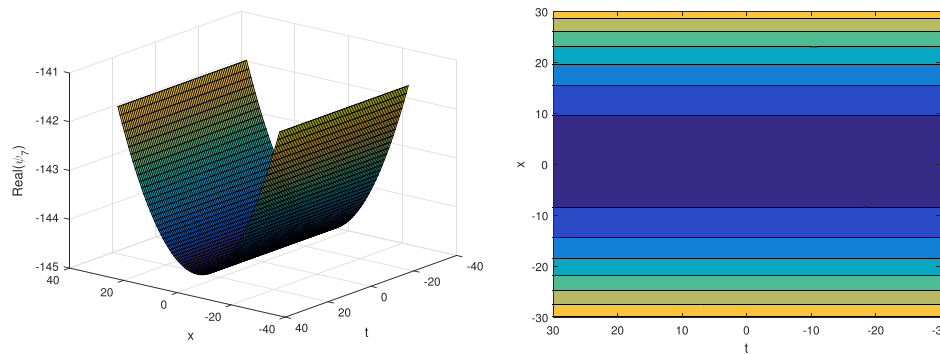


Fig. 6. The 3D plot and the corresponding contour representation of the traveling wave solutions  $\psi_7(x, t)$  for different values of parameters.

$$\psi_{10} = \left( \left( \pm \frac{\sqrt{-2}\phi'(x)}{4} - \sqrt{512\gamma^3 + \phi'(x)^2}(-z^{-1}) \right) \pm 2\gamma^2\sqrt{-30}z^{-2} \right) e^{ip}$$

The graphs and its corresponding contours plot of the solutions  $\psi_6(x, t)$  are shown in Fig. 5, for the values of parameters  $\rho_1 = 0.007, \rho_2 = 0.08, \gamma = 0.005, \phi(x) = \sin(x)$  and  $\nu = 0.0019$ .

The graphs and its corresponding contour plot of the solutions  $\psi_7(x, t)$  are shown in Fig. 6, for the values of parameters  $\rho_1 = 0.007, \rho_2 = 0.08, \gamma = 0.005, \phi(x) = \sin(x)$  and  $\nu = 0.0019$ .

The graphs and its corresponding contour plot of the solutions  $\psi_8(x, t)$  are shown in Fig. 7, for the values of parameters  $\rho_1 = 0.007, \rho_2 = 0.08, \gamma = 0.005, \phi(x) = \sin(x)$  and  $\nu = 0.0019$ .

For Case-3: The  $\beta_2, \gamma_1, \nu, \gamma$  are free parameters while along with  $\gamma_2 = 0$ ,

$$\begin{aligned} \beta_0 &= 0, & \beta_1 &= 0, & \beta_2 &= \pm 2\sqrt{-30}, & \gamma_1 &= \pm \\ & \sqrt{-40\gamma - 4\nu p_1 + 12p_1^2 + 2\phi(x) + 2p_2}, & 2p_1\nu &= -40\gamma + 6p_1^2 + \phi(x) + \\ & p_2, & \gamma &= \frac{1}{8}\sqrt{p_1^2 + p_2 + \phi(x)p_1}, \end{aligned}$$

The constraint condition holds for the following results, when vis putting in it,  $\gamma = -\frac{3}{40}\phi(x) - \frac{3}{40}p_2 + \frac{1}{10}p_1^2$ .

For  $\gamma < 0$ , thus following families of hyperbolic solutions are obtained.

$$\begin{aligned} \psi_{11} &= \left( \pm 2\sqrt{-30}(-\sqrt{-\gamma}\tanh(\sqrt{-\gamma}z))^2 \right. \\ & \left. \pm \sqrt{-40\gamma - 4\nu p_1 + 12p_1^2 + 2\phi(x) + 2p_2}(-\sqrt{-\gamma}\tanh(\sqrt{-\gamma}z))^{-1} \right) e^{ip}, \end{aligned}$$

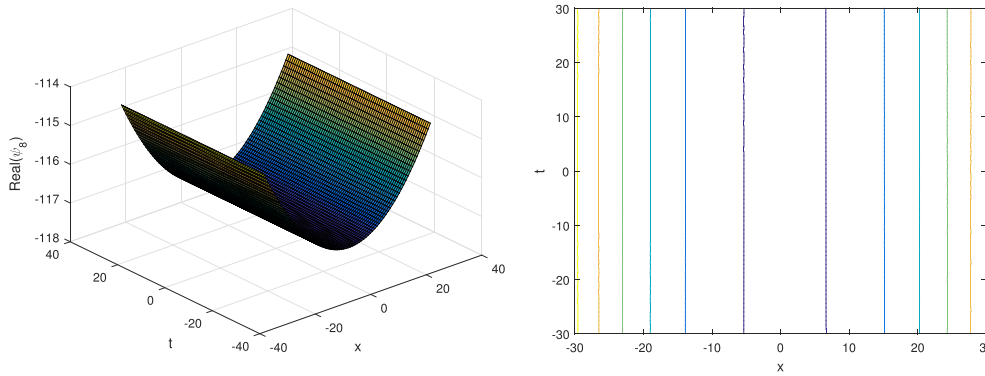


Fig. 7. The 3D plot and the corresponding contour representation of the traveling wave solutions  $\psi_8(x, t)$  for different values of parameters.

or

$$\psi_{12} = \left( \pm 2\sqrt{-30} (-\sqrt{-\gamma} \coth(\sqrt{-\gamma} z))^2 \pm \sqrt{-40\gamma - 4vp_1 + 12p_1^2 + 2\phi(x) + 2p_2 (-\sqrt{-\gamma} \coth(\sqrt{-\gamma} z))^{-1}} \right) e^{ip}$$

For  $\gamma > 0$ , thus following periodic wave solutions are obtained.

$$\psi_{13} = \left( \pm 2\sqrt{-30} (\sqrt{\gamma} \tan(\sqrt{\gamma} z))^2 \pm \sqrt{-40\gamma - 4vp_1 + 12p_1^2 + 2\phi(x) + 2p_2 (\sqrt{\gamma} \tan(\sqrt{\gamma} z))^{-1}} \right) e^{ip}$$

$$16p_1^9 \gamma v = -\phi(x)p_1^{10} + 136\gamma^2 p_1^8 - 8\phi(x)\gamma p_1^8 - 8p_2\gamma p_1^8 + p_2 p_1^{10} + 6\sqrt{-30} p_1^8 + p_1^{12} - 48p_1^{10}\gamma + 48(\phi'(x)(2\sqrt{-30}))^2$$

or

$$\psi_{14} = \left( \pm 2\sqrt{-30} (-\sqrt{\gamma} \cot(\sqrt{\gamma} z))^2 \pm \sqrt{-40\gamma - 4vp_1 + 12p_1^2 + 2\phi(x) + 2p_2 (-\sqrt{\gamma} \cot(\sqrt{\gamma} z))^{-1}} \right) e^{ip}$$

For  $\gamma = 0$ , thus following plane solutions are obtained.

$$\psi_{15} = \left( \pm 2\sqrt{-30} z^{-2} \pm \sqrt{-40\gamma - 4vp_1 + 12p_1^2 + 2\phi(x) + 2p_2 (-z^{-1})} \right) e^{ip}$$

The graph and the corresponding contour plot of the solutions  $\psi_{11}(x, t)$  are shown in Fig. 8, for the values of parameters  $\rho_1 = 1.1, \rho_2 = 0.0002, \gamma = 0.001, \phi(x) = \sin(x)$  and  $\nu = 0.001$

The graph and the corresponding contour plot of the solutions  $\psi_{12}(x, t)$  are shown in Fig. 9, for the values of parameters  $\rho_1 = 1.1, \rho_2 = 0.0002, \gamma = 0.001, \phi(x) = \sin(x)$  and  $\nu = 0.001$ .

The graph and its corresponding contour plot of the solutions  $\psi_{14}(x,$

$t)$  are shown in Fig. 10, for the values of parameters  $\rho_1 = 1.1, \rho_2 = 0.0002, \gamma = 0.001, \phi(x) = \sin(x)$  and  $\nu = 0.001$ .

The graph and its corresponding contour plot of the solutions  $\psi_{15}(x, t)$  are shown in Fig. 11, for the values of parameters  $\rho_1 = 1.1, \rho_2 = 0.0002, \gamma = 0.001, \phi(x) = \sin(x)$  and  $\nu = 0.001$ .

For Case-4: The  $\beta_0, \beta_2, \gamma_2, v$ , are free parameters while along with  $\gamma_1 = 0, \beta_1 = 0, \beta_2 = \pm 2\sqrt{-30}, \beta_0 = \pm 4 \frac{\phi'(x)(2\sqrt{-30})}{p_1^4}, \gamma_2 = \frac{\gamma(\phi(x)p_1^{10} - 184\gamma^2 p_1^8 + p_2 p_1^{10} + 6\sqrt{-30} p_1^8 + p_1^{12} + 48(\phi'(x)(2\sqrt{-30}))^2)}{16(p_1^4 \phi'(x)(2\sqrt{-30}))}, \gamma = \frac{-6p_1^6 - \phi(x)p_1^4 - p_2 p_1^4 + 4\sqrt{-30}\phi'(x)(2\sqrt{-30}) + 2vp_1^5}{40p_1^4}$

The constraint condition holds for the following results, when vis putting in it,

$$\gamma = \frac{(8p_1^6 - 6p_1^4 p_2 - 6p_1^4 \phi(x) - \sqrt{\alpha_1 + \alpha_2})}{68p_1^4}$$

where

$$\alpha_1 = -30p_1^{12} + 130p_2 p_1^{10} + 130\phi(x)p_1^{10} - 36p_1^8 p_2^2 - 72p_1^8 p_2 \phi(x) - 36p_1^8 \phi^2(x)$$

$$\alpha_2 = 204\sqrt{-30} p_1^8 + 1632(\phi(x)(2\sqrt{-30}))^2$$

For  $\gamma < 0$ , thus following families of hyperbolic wave solutions are obtained.

$$\psi_{16} = \left( \pm \frac{8\phi'(x)(\sqrt{-30})}{p_1^4} \pm 2\sqrt{-30} (-\sqrt{-\gamma} \tanh(\sqrt{-\gamma} z))^2 \pm \frac{\gamma(\phi(x)p_1^{10} - 184\gamma^2 p_1^8 + p_2 p_1^{10} + 6sqrt - 30p_1^8 + p_1^{12} + 48(\phi'(x)(2\sqrt{-30}))^2)}{16(p_1^4 \phi'(x)(2\sqrt{-30}))} (-\sqrt{-\gamma} \tanh(\sqrt{-\gamma} z))^{-2} \right) e^{ip}$$

or

$$\psi_{17} = \left( \pm \frac{8\phi'(x)(\sqrt{-30})}{p_1^4} \pm 2\sqrt{-30}(-\sqrt{-\gamma} \coth(\sqrt{-\gamma}z))^2 \pm \frac{\gamma(\phi(x)p_1^{10} - 184\gamma^2 p_1^8 + p_2 p_1^{10} + 6sqrt - 30p_1^8 + p_1^{12} + 48(\phi'(x)(2\sqrt{-30}))^2)}{16(p_1^4 \phi'(x)(2\sqrt{-30}))} (-\sqrt{-\gamma} \coth \sqrt{-\gamma} z)^{-2} \right) e^{ip}$$

For  $\gamma > 0$ , thus following periodic wave solutions are obtained.

$$\psi_{18} = \left( \pm \frac{8\phi'(x)(\sqrt{-30})}{p_1^4} \pm 2\sqrt{-30}(-\sqrt{-\gamma} \tan(\sqrt{-\gamma}z))^2 \pm \frac{\gamma(\phi(x)p_1^{10} - 184\gamma^2 p_1^8 + p_2 p_1^{10} + 6sqrt - 30p_1^8 + p_1^{12} + 48(\phi'(x)(2\sqrt{-30}))^2)}{16(p_1^4 \phi'(x)(2\sqrt{-30}))} (-\sqrt{-\gamma} \tan \sqrt{-\gamma} z)^{-2} \right) e^{ip}$$

or

$$\psi_{19} = \left( \pm \frac{8\phi'(x)(\sqrt{-30})}{p_1^4} \pm 2\sqrt{-30}(-\sqrt{-\gamma} \cot(\sqrt{-\gamma}z))^2 \pm \frac{\gamma(\phi(x)p_1^{10} - 184\gamma^2 p_1^8 + p_2 p_1^{10} + 6sqrt - 30p_1^8 + p_1^{12} + 48(\phi'(x)(2\sqrt{-30}))^2)}{16(p_1^4 \phi'(x)(2\sqrt{-30}))} (-\sqrt{-\gamma} \cot \sqrt{-\gamma} z)^{-2} \right) e^{ip}$$

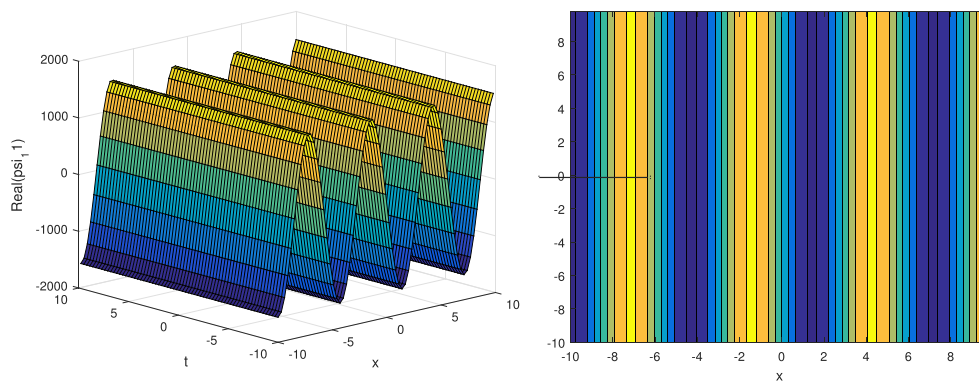


Fig. 8. The 3D plot and the corresponding contour representation of the traveling wave solutions  $\psi_{11}(x, t)$  for different values of parameters.

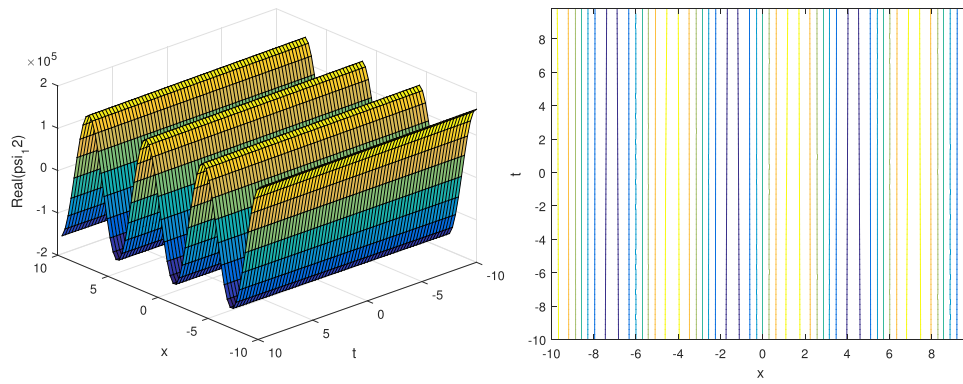


Fig. 9. The 3D plot and the corresponding contour representation of the traveling wave solutions  $\psi_{12}(x, t)$  for different values of parameters.

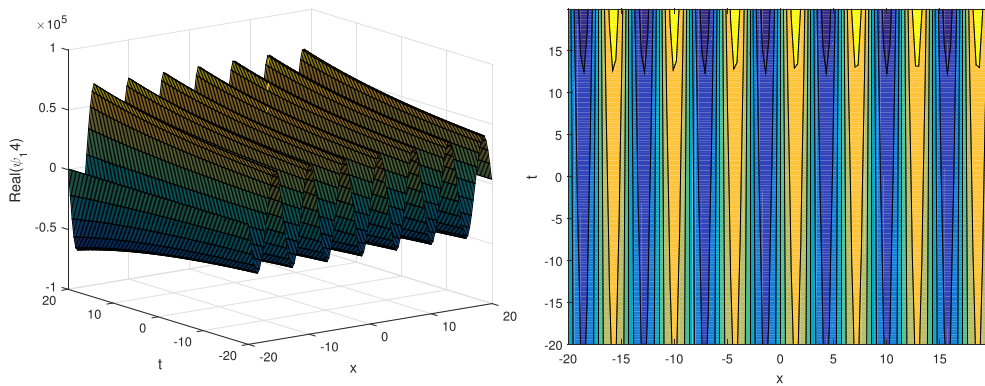


Fig. 10. The 3D plot and the corresponding contour representation of the traveling wave solutions  $\psi_{14}(x, t)$  for different values of parameters.

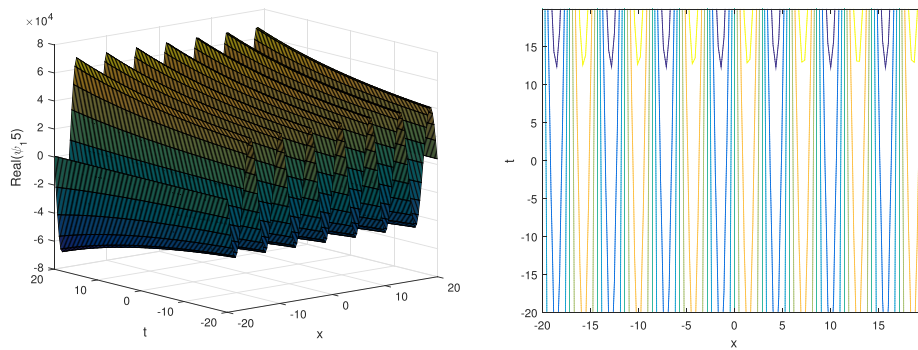


Fig. 11. The 3D plot and the corresponding contour representation of the traveling wave solutions  $\psi_{15}(x, t)$  for different values of parameters.

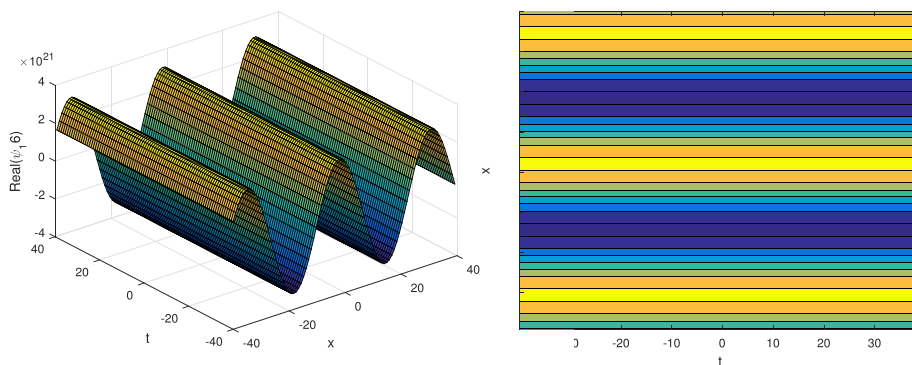


Fig. 12. The 3D plot and the corresponding contour representation of the traveling wave solutions  $\psi_{16}(x, t)$  for different values of parameters.

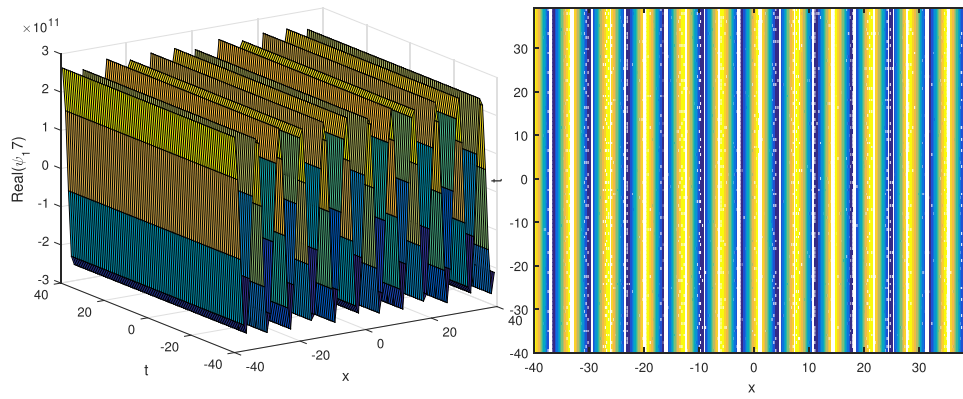


Fig. 13. The 3D plot and the corresponding contour representation of the traveling wave solutions  $\psi_{17}(x, t)$  for different values of parameters.

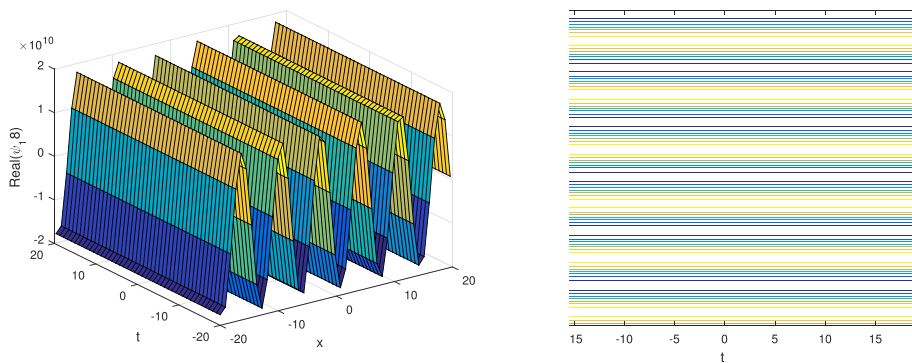


Fig. 14. The 3D plot and the corresponding contour representation of the traveling wave solutions  $\psi_{18}(x, t)$  for different values of parameters.

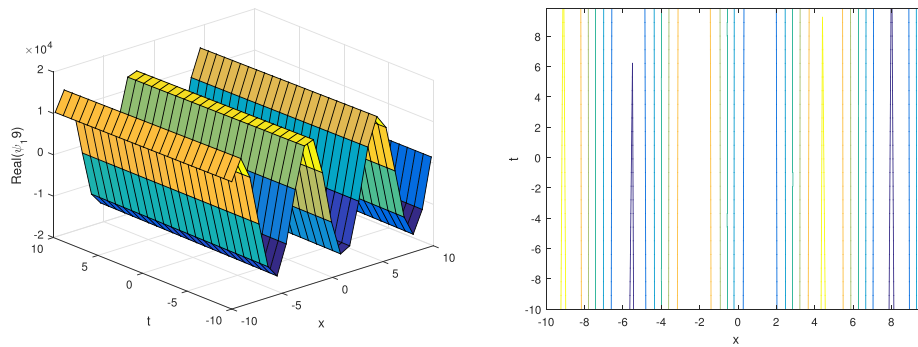


Fig. 15. The 3D plot and the corresponding contour representation of the traveling wave solutions  $\psi_{19}(x, t)$  for different values of parameters.

For  $\gamma = 0$ , thus following plane solutions are obtained.

The graph and its corresponding contour plot of the solutions  $\psi_{17}(x,$

$$\psi_{20} = \left( \pm \frac{8\phi'(x)(\sqrt{-30})}{p_1^4} \pm 2\sqrt{-30}z^{-2} \pm \frac{\gamma(\phi(x)p_1^{10} - 184\gamma^2p_1^8 + p_2p_1^{10} + 6sqrt - 30p_1^8 + p_1^{12} + 48(\phi'(x)(2\sqrt{-30}))^2)}{16(p_1^4\phi'(x)(2\sqrt{-30}))} z^{-2} \right) e^{ip}$$

The graph and its corresponding contour plot of the solutions  $\psi_{16}(x, t)$  are shown in Fig. 12, for the values of parameters  $\rho_1 = 0.9190, \rho_2 = 1000.08, \gamma = 0.005, \phi(x) = e^x$  and  $\nu = 0.005$

$t)$  are shown in Fig. 13, for the values of parameters  $\rho_1 = 0.9190, \rho_2 = 1000.08, \gamma = 0.005, \phi(x) = e^x$  and  $\nu = 0.005$ .

The graph and its corresponding contour plot of the solutions  $\psi_{18}(x,$



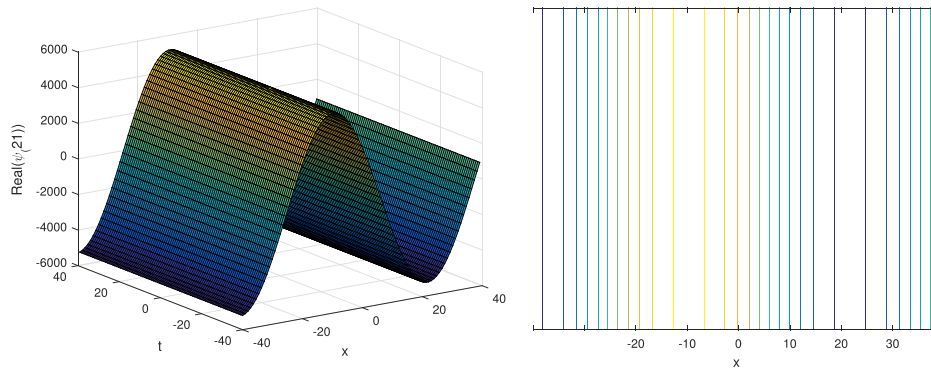


Fig. 16. The 3D plot and the corresponding contour representation of the traveling wave solutions  $\psi_{21}(x, t)$  for different values of parameters.

$t$ ) are shown in Fig. 14, for the values of parameters  $\rho_1 = 0.9190, \rho_2 = 1000.08, \gamma = 0.005, \phi(x) = e^x$  and  $\nu = 0.005$ .

The graph and its corresponding contour plot of the solutions  $\psi_{19}(x, t)$  are shown in Fig. 15, for the values of parameters  $\rho_1 = 0.9190, \rho_2 = 1000.08, \gamma = 0.005, \phi(x) = e^x$  and  $\nu = 0.005$ .

For Case-5: The  $\gamma_2, \beta_0$ , v are free parameters while along with  $\beta_1 = 0, \beta_2 = 0, \gamma_1 = 0$ ,

$$\gamma_2 = \frac{1\sqrt{24\phi(x)\gamma - 3p_1^4 - 48vp_1\gamma + 144p_1^2\gamma - 408\gamma^2 + 24p_2\gamma - 3\phi(x)p_1^2 - 3p_2p_1^2} \times \alpha_1}{-2vp_1 + 6p_1^2 + \phi(x) + p_2 - 8\gamma},$$

where

$$\alpha_1 = (4\phi(x)\gamma + p_1^4 - 8vp_1\gamma + 24p_1^2\gamma - 68\gamma^2 + 4p_2\gamma + \phi(x)p_1^2 + p_2p_1^2)$$

$$\gamma_2 = \frac{1\sqrt{24\phi(x)\gamma - 3p_1^4 - 48vp_1\gamma + 144p_1^2\gamma - 408\gamma^2 + 24p_2\gamma - 3\phi(x)p_1^2 - 3p_2p_1^2} \times \alpha_1}{-2vp_1 + 6p_1^2 + \phi(x) + p_2 - 8\gamma},$$

where

$$\alpha_1 = (4\phi(x)\gamma + p_1^4 - 8vp_1\gamma + 24p_1^2\gamma - 68\gamma^2 + 4p_2\gamma + \phi(x)p_1^2 + p_2p_1^2)$$

$$\beta_0 = 1 / \sqrt{3\sqrt{24\phi(x)\gamma - 3p_1^4 - 48vp_1\gamma + 144p_1^2\gamma - 408\gamma^2 + 24p_2\gamma - 3\phi(x)p_1^2 - 3p_2p_1^2}},$$

$$v = 1 / \sqrt{14p_1^4 + 3p_1^2\gamma + 1 / \sqrt{14p_2p_1^2 + 1 / \sqrt{14\phi(x)p_1^2 - 4\gamma^2 + 1 / \sqrt{2\phi(x)\gamma + 1 / \sqrt{2p_2\gamma - \frac{3}{28p_1\gamma} \{56\gamma^2\phi(x)p_1^2 + 56\gamma^2p_1^4 + 4\phi(x)p_1^6 + 4p_2p_1^4\phi(x) + 2(\phi(x))^2p_1^4 + 4p_2p_1^6 + 2p_2^2p_1^4 - 4032\gamma^4 + 2p_1^8 + 56\gamma^2p_2p_1^2\}}^{1/2}}}}}}$$

$$\beta_0 = 1 / \sqrt{3\sqrt{24\phi(x)\gamma - 3p_1^4 - 48vp_1\gamma + 144p_1^2\gamma - 408\gamma^2 + 24p_2\gamma - 3\phi(x)p_1^2 - 3p_2p_1^2}},$$

For Case-5: The  $\gamma_2, \beta_0$ , v are free parameters while along with  $\beta_1 = 0, \beta_2 = 0, \gamma_1 = 0$ ,

For  $\gamma < 0$ , thus following families of hyperbolic wave solutions are obtained.

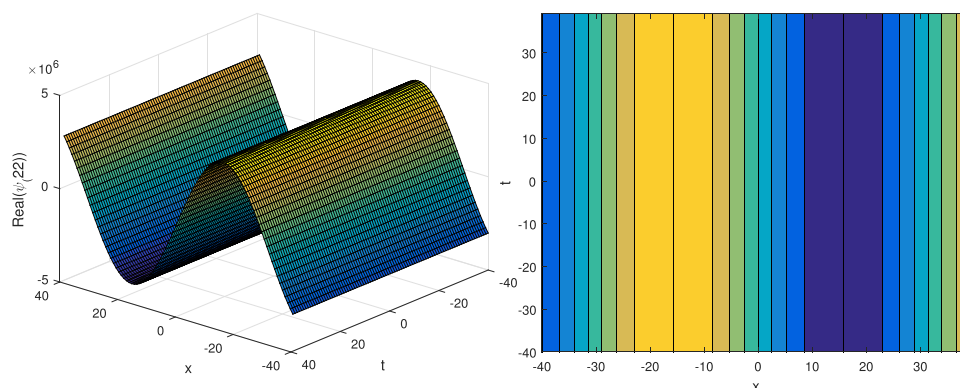


Fig. 17. The 3D plot and the corresponding contour representation of the traveling wave solutions  $\psi_{22}(x, t)$  for different values of parameters.

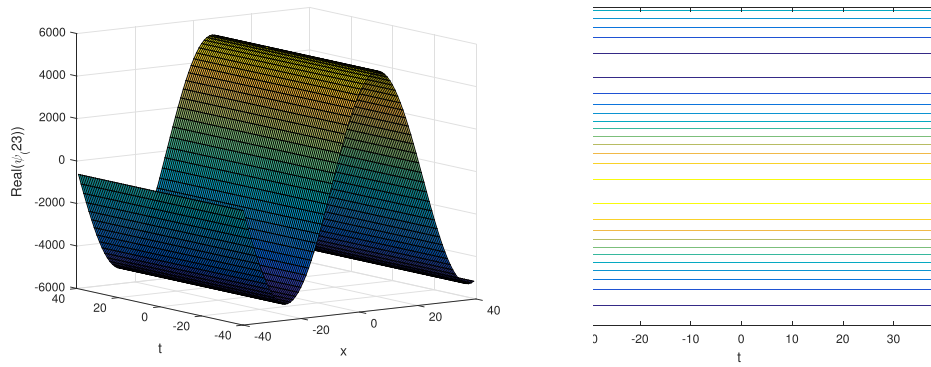


Fig. 18. The 3D plot and the corresponding contour representation of the traveling wave solutions  $\psi_{23}(x, t)$  for different values of parameters.

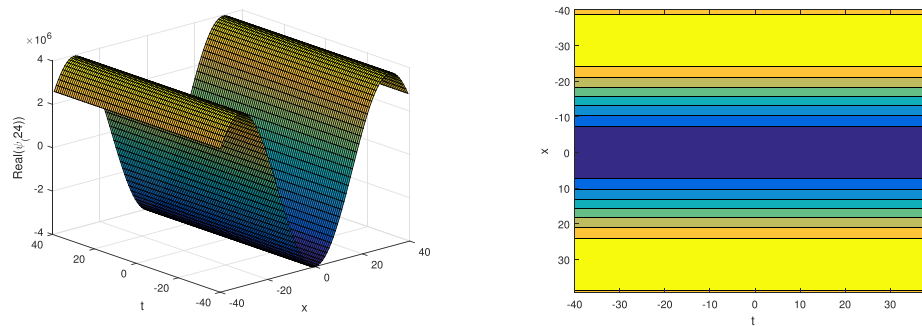


Fig. 19. The 3D plot and the corresponding contour representation of the traveling wave solutions  $\psi_{24}(x, t)$  for different values of parameters.

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$$\psi_{21} = \left\{ \pm 1 \sqrt{3 \sqrt{24\phi(x)\gamma - 3p_1^4 - 48vp_1\gamma + 144p_1^2\gamma - 408\gamma^2 + 24p_2\gamma - 3\phi(x)p_1^2 - 3p_2p_1^2}} \right. \\ \left. \pm \frac{1}{9} \frac{\sqrt{24\phi(x)\gamma - 3p_1^4 - 48vp_1\gamma + 144p_1^2\gamma - 408\gamma^2 + 24p_2\gamma - 3\phi(x)p_1^2 - 3p_2p_1^2} \times \alpha_1 (-\sqrt{-\gamma} \tanh(\sqrt{-\gamma}z))^{-2}}{-2vp_1 + 6p_1^2 + \phi(x) + p_2 - 8\gamma} \right\} e^{ip},$$

or

---


$$\psi_{22} = \left\{ \pm 1 \sqrt{3 \sqrt{24\phi(x)\gamma - 3p_1^4 - 48vp_1\gamma + 144p_1^2\gamma - 408\gamma^2 + 24p_2\gamma - 3\phi(x)p_1^2 - 3p_2p_1^2}} \right. \\ \left. \pm \frac{1}{9} \frac{\sqrt{24\phi(x)\gamma - 3p_1^4 - 48vp_1\gamma + 144p_1^2\gamma - 408\gamma^2 + 24p_2\gamma - 3\phi(x)p_1^2 - 3p_2p_1^2} \times \alpha_1 (-\sqrt{-\gamma} \coth(\sqrt{-\gamma}z))^{-2}}{-2vp_1 + 6p_1^2 + \phi(x) + p_2 - 8\gamma} \right\} e^{ip}.$$


---

For  $\gamma > 0$ , thus following periodic wave solutions are obtained.

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$$\psi_{23} = \left\{ \pm 1 \sqrt{3 \sqrt{24\phi(x)\gamma - 3p_1^4 - 48vp_1\gamma + 144p_1^2\gamma - 408\gamma^2 + 24p_2\gamma - 3\phi(x)p_1^2 - 3p_2p_1^2}} \right. \\ \left. \pm \frac{1}{9} \frac{\sqrt{24\phi(x)\gamma - 3p_1^4 - 48vp_1\gamma + 144p_1^2\gamma - 408\gamma^2 + 24p_2\gamma - 3\phi(x)p_1^2 - 3p_2p_1^2} \times \alpha_1 (-\sqrt{-\gamma} \tan(\sqrt{-\gamma}z))^{-2}}{-2vp_1 + 6p_1^2 + \phi(x) + p_2 - 8\gamma} \right\} e^{ip},$$


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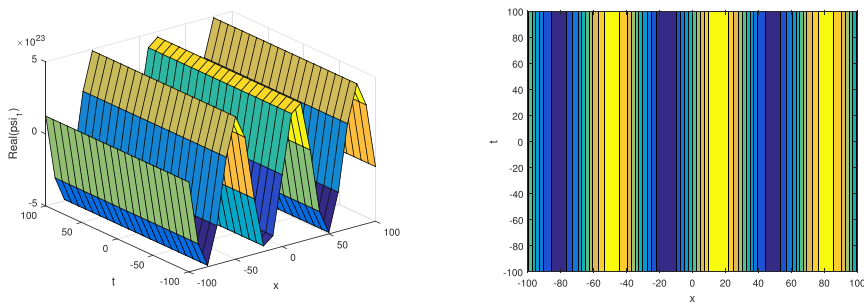


Fig. 20. The 3D plot and the corresponding contour representation of the traveling wave solutions  $\psi_1(x, t)$  for different values of parameters.

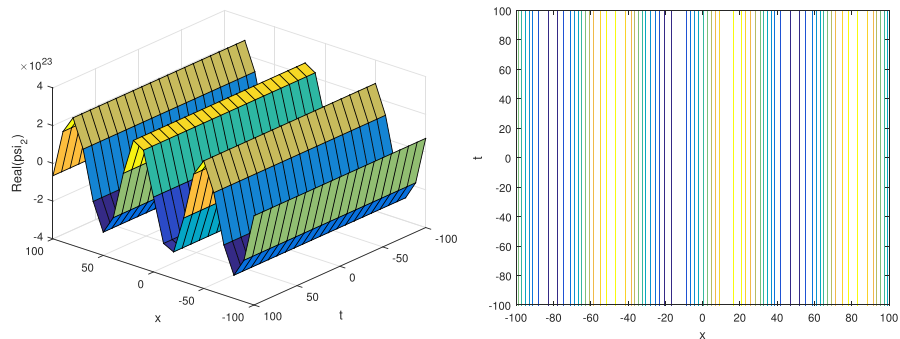


Fig. 21. The 3D plot and the corresponding contour representation of the traveling wave solutions  $\psi_2(x, t)$  for different values of parameters.

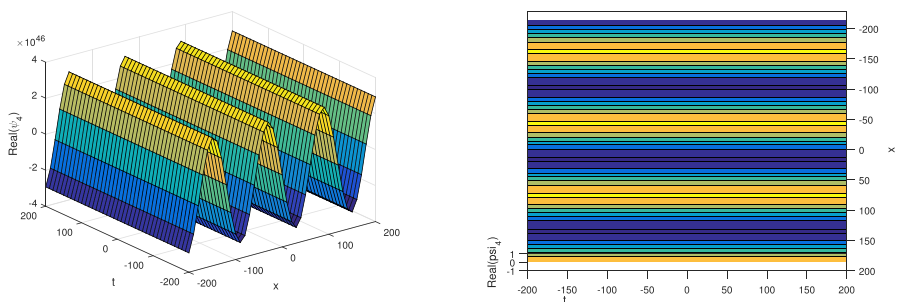


Fig. 22. The 3D plot and the corresponding contour representation of the traveling wave solutions  $\psi_4(x, t)$  for different values of parameters.

or

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$$\psi_{24} = \left\{ \pm 1 \sqrt{3 \sqrt{24\phi(x)\gamma - 3p_1^4 - 48vp_1\gamma + 144p_1^2\gamma - 408\gamma^2 + 24p_2\gamma - 3\phi(x)p_1^2 - 3p_2p_1^2}} \right. \\ \left. \pm \frac{1}{9} \frac{\sqrt{24\phi(x)\gamma - 3p_1^4 - 48vp_1\gamma + 144p_1^2\gamma - 408\gamma^2 + 24p_2\gamma - 3\phi(x)p_1^2 - 3p_2p_1^2} \times \alpha_1 (-\sqrt{-\gamma} \cot(\sqrt{-\gamma}z))^{-2}}{-2vp_1 + 6p_1^2 + \phi(x) + p_2 - 8\gamma} \right\} e^{ip}$$


---

For  $\gamma = 0$ , thus following plane solutions are obtained.

$$\psi_{25} = \left\{ \pm 1 \left/ 3 \sqrt{24\phi(x)\gamma - 3p_1^4 - 48vp_1\gamma + 144p_1^2\gamma - 408\gamma^2 + 24p_2\gamma - 3\phi(x)p_1^2 - 3p_2p_1^2} \right. \right. \\ \left. \left. \pm \frac{1}{9} \frac{\sqrt{24\phi(x)\gamma - 3p_1^4 - 48vp_1\gamma + 144p_1^2\gamma - 408\gamma^2 + 24p_2\gamma - 3\phi(x)p_1^2 - 3p_2p_1^2} \times \alpha_1}{-2vp_1 + 6p_1^2 + \phi(x) + p_2 - 8\gamma} (-z^{-2}) \right\} e^{ip}$$

The graph and its corresponding contour plot of the solutions  $\psi_{21}(x, t)$  are shown in Fig. 16, for the values of parameters  $\rho_1 = 0.1, \rho_2 = 2, \gamma = 0.005, \phi(x) = \sin(x)$  and  $\nu = 0.009$

The graph and its corresponding contour plot of the solutions  $\psi_{22}(x, t)$  are shown in Fig. 17, for the values of parameters  $\rho_1 = 0.1, \rho_2 = 2, \gamma = 0.005, \phi(x) = \sin(x)$  and  $\nu = 0.009$ .

The graph and its corresponding contour plot of the solutions  $\psi_{23}(x, t)$  are shown in Fig. 18, for the values of parameters  $\rho_1 = 0.1, \rho_2 = 2, \gamma = 0.005, \phi(x) = \sin(x)$  and  $\nu = 0.009$ .

The graph and its corresponding contour plot of the solutions  $\psi_{24}(x, t)$  are shown in Fig. 19, for the values of parameters  $\rho_1 = 0.1, \rho_2 = 2, \gamma = 0.005, \phi(x) = \sin(x)$  and  $\nu = 0.009$ .

Using  $\left(\frac{G'}{G}\right)$ -expansion method

$$\alpha_1 = 1728\phi(x)\mu p_1^2 + 324p_2^2\mu - 54p_1^6\phi(x)^2 - 144p_1^6 + 54p_2^2p_1^2 - 12(\phi'(x))^2 - 144p_1^2 + 324(\phi(x))^2\mu - 198p_1^4\phi(x)$$

The solution of Eq. (5) has the formal solutions of the following form.

$$U(z) = \beta_0 + \beta_1 \left(\frac{G'(z)}{G(z)}\right) + \beta_2 \left(\frac{G'(z)}{G(z)}\right)^2, \tag{7}$$

Differentiating Eq. (7) w.r.t  $z$  and then putting  $G''(z) + \lambda G'(z) + \mu G(z) =$

$$\psi_1(x, t) = \left\{ \frac{\sqrt{-2\phi(x) + 4vp_1 - 2p_2 + 12p_1^2} (3\phi(x)\lambda - 18p_1^2\lambda - 6vp_1\lambda + 3p_2\lambda - 2\phi'(x))}{6\phi(x) - 12vp_1 + 6p_2 - 36p_1^2} + \sqrt{-2\phi(x) + 4vp_1 - 2p_2 + 12p_1^2} \times \left(-\frac{\lambda}{2}\right) \right. \\ \left. + \frac{\sqrt{\lambda^2 - 4\mu}}{2} \left( \frac{\text{Acosh}\left(z/2\sqrt{\lambda^2 - 4\mu}\right) + \text{Bsinh}\left(z/2\sqrt{\lambda^2 - 4\mu}\right)}{\text{Asinh}\left(z/2\sqrt{\lambda^2 - 4\mu}\right) + \text{Bcosh}\left(z/2\sqrt{\lambda^2 - 4\mu}\right)} \right) \right\} \times e^{ip} \tag{8}$$

0.

After equating each coefficient equal to zero, this yields a set of following algebraic equations. For Case-1: The vis free parameter while along with  $\beta_2 = 0, \beta_1 = \pm \sqrt{-2\phi(x) + 4vp_1 - 2p_2 + 12p_1^2}$ ,

$$\beta_0 = \pm \frac{\sqrt{-2\phi(x) + 4vp_1 - 2p_2 + 12p_1^2} (3\phi(x)\lambda - 18p_1^2\lambda - 6vp_1\lambda + 3p_2\lambda - 2\phi'(x))}{6\phi(x) - 12vp_1 + 6p_2 - 36p_1^2}$$

$$v = \frac{1}{6(4\mu - \lambda^2)p_1} \left( 12\phi(x)\mu - 3\phi(x)\lambda^2 + 18p_1^2\lambda^2 - 3p_2p_1^2 + 12p_2\mu - 3p_2\lambda^2 \right. \\ \left. + 3 - 72p_2^2\mu + 3p_1^4 + 3\phi(x)p_1^2 \right. \\ \left. - (9 - 18p_2p_1^2 + 18\phi(x)p_1^2 + 48\mu(\phi'(x))^2 - 12\lambda^2(\phi'(x))^2 - 18p_1^6p_2 \right. \\ \left. + 9p_1^4(\phi(x))^2 + 18p_1^6\phi(x) + 9p_1^4p_2^2 + 9p_1^8 - 18p_1^4\phi(x)p_2 + 18p_1^4) \right)^{1/2}$$

The constraint condition holds for the following results, when vis putting in it,

$$\lambda = -\frac{\sqrt{1728p_1^2p_2\mu - 54p_2 + 2304p_1^4\mu + 90p_1^4p_2 - 54\phi(x) + \alpha_1 + 648\phi(x)\mu p_2}}{3\phi(x) + 8p_1^2 + 3p_2}$$

where

- (i) When  $\lambda^2 - 4\mu > 0$  thus following hyperbolic solutions are obtained.
- (ii) When  $\lambda^2 - 4\mu < 0$  thus following trigonometric solutions are obtained.

$$\psi_2(x, t) = \left\{ \frac{\sqrt{-2\phi(x) + 4vp_1 - 2p_2 + 12p_1^2} (3\phi(x)\lambda - 18p_1^2\lambda - 6vp_1\lambda + 3p_2\lambda - 2\phi'(x))}{6\phi(x) - 12vp_1 + 6p_2 - 36p_1^2} + \sqrt{-2\phi(x) + 4vp_1 - 2p_2 + 12p_1^2} \times \left(-\frac{\lambda}{2}\right) \right. \\ \left. + \frac{\sqrt{\lambda^2 - 4\mu}}{2} \left( \frac{-\text{Asin}\left(z/2\sqrt{\lambda^2 - 4\mu}\right) + \text{Bcos}\left(z/2\sqrt{\lambda^2 - 4\mu}\right)}{\text{Acos}\left(z/2\sqrt{\lambda^2 - 4\mu}\right) + \text{Bsin}\left(z/2\sqrt{\lambda^2 - 4\mu}\right)} \right) \right\} \times e^{ip} \tag{9}$$

- (iii) When  $\lambda^2 - 4\mu = 0$  thus following plane solutions are obtained.

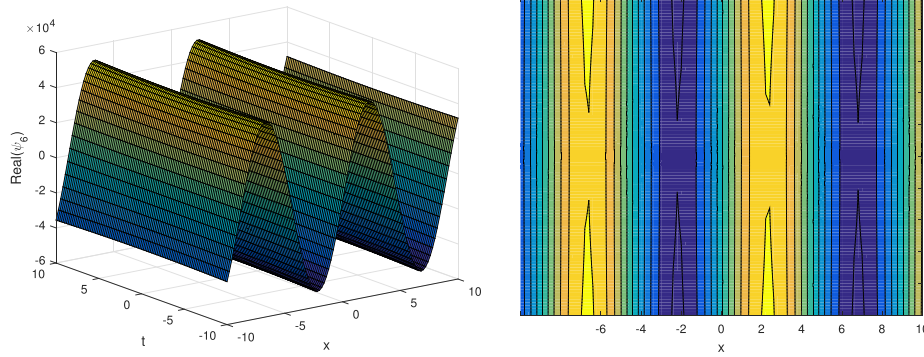


Fig. 23. The 3D plot and the corresponding contour representation of the traveling wave solutions  $\psi_6(x, t)$  for different values of parameters.

$$\psi_3(x, t) = \left( -\frac{\lambda}{2} + \frac{B}{A + Bz} \right) \times e^{ip} \tag{10}$$

where

In particular, if  $A \rightarrow 0$  and  $B \neq 0$  and  $\lambda > 0$  then the solution Eq. (8)

The graph and its corresponding contour plot of the solutions  $\psi_1(x, t)$  are shown in Fig. 20, for the values of parameters  $\rho_1 = 700, \rho_2 = 0.0008, \lambda = 0.05, \mu = 0.001, A = 0.001, B = 50, \phi(x) = \sin(x)$  and  $\nu = 0.02$

The graph and its corresponding contour plot of the solutions  $\psi_2(x, t)$

$$z = x - \frac{1}{6(4\mu - \lambda^2)p_1} \left( 12\phi(x)\mu - 3\phi(x)\lambda^2 + 18\rho_1^2\lambda^2 - 3\rho_2p_1^2 + 12p_2\mu - 3\rho_2\lambda^2 + 3 - 72\rho_1^2\mu + 3\rho_1^4 + 3\phi(x)p_1^2 - (9 - 18\rho_2p_1^2 + 18\phi(x)p_1^2 + 48\mu(\phi'(x))^2 - 12\lambda^2(\phi'(x))^2 - 18\rho_1^6p_2 + 9\rho_1^4(\phi(x))^2 + 18\rho_1^6\phi(x) + 9\rho_1^4p_2^2 + 9\rho_1^8 - 18\rho_1^4\phi(x)p_2 + 18\rho_1^4)^{1/2} \right) t.$$

converges to

$$\psi_4(x, t) = \left\{ \frac{\sqrt{-2\phi(x) + 4vp_1 - 2p_2 + 12p_1^2} (3\phi(x)\lambda - 18\rho_1^2\lambda - 6vp_1\lambda + 3p_2\lambda - 2\phi'(x))}{6\phi(x) - 12vp_1 + 6p_2 - 36\rho_1^2} + \sqrt{-2\phi(x) + 4vp_1 - 2p_2 + 12p_1^2} \left( \frac{\lambda}{2} + \frac{\sqrt{\lambda^2 - 4\mu}}{2} \tanh\left(\frac{z}{2}\sqrt{\lambda^2 - 4\mu}\right) \right) \right\} \times e^{ip}. \tag{11}$$

Also, if  $B \rightarrow 0$  and  $A \neq 0$  and  $\lambda > 0$  then the solution Eq. (8) converges to

$$\psi_5(x, t) = \left\{ \frac{\sqrt{-2\phi(x) + 4vp_1 - 2p_2 + 12p_1^2} (3\phi(x)\lambda - 18\rho_1^2\lambda - 6vp_1\lambda + 3p_2\lambda - 2\phi'(x))}{6\phi(x) - 12vp_1 + 6p_2 - 36\rho_1^2} + \sqrt{-2\phi(x) + 4vp_1 - 2p_2 + 12p_1^2} \left( \frac{\lambda}{2} + \frac{\sqrt{\lambda^2 - 4\mu}}{2} \coth\left(\frac{z}{2}\sqrt{\lambda^2 - 4\mu}\right) \right) \right\} \times e^{ip}. \tag{12}$$

are shown in Fig. 21, for the values of parameters  $\rho_1 = 700, \rho_2 = 0.0008, \lambda = 0.05, \mu = 0.001, A = 0.001, B = 50, \phi(x) = \sin(x)$  and  $\nu = 0.02$ .

The graph and its corresponding contour plot of the solutions  $\psi_4(x, t)$  are shown in Fig. 22, for the values of parameters  $\rho_1 = 700, \rho_2 = 0.0008, \lambda = 0.05, \mu = 0.001, A = 0.001, B = 50, \phi(x) = \sin(x)$  and  $\nu = 0.02$ .

For Case-2; The  $\nu$  is free parameter while along with  $\beta_0 = 0, \beta_1 =$

$$\frac{1}{5} \sqrt{-20 \frac{\phi'(x)}{\lambda}}, \beta_2 = \frac{4\phi'(x)\sqrt{-20 \frac{\phi'(x)}{\lambda}}}{((-60\lambda^2 + 80\mu)\phi'(x) + 25\lambda(p_1^4 + 1 - p_2p_1^2 + \phi(x)p_1^2))},$$

$$\nu = \frac{-30\rho_1^2\lambda + 5\phi(x)\lambda + 5p_2\lambda - 2\phi'(x)}{10\rho_1\lambda},$$

The constraint condition holds for the following results, when  $\nu$  is putting in it,

$$\lambda = -\frac{2\phi'(x)}{40\rho_1^2 + 15\phi(x) + 15p_2}$$

(i) When  $\lambda^2 - 4\mu > 0$  thus following families of hyperbolic solutions

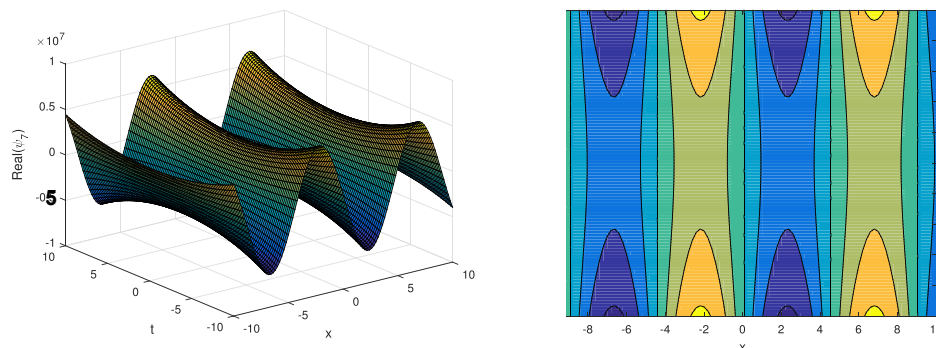


Fig. 24. The 3D plot and the corresponding contour representation of the traveling wave solutions  $\psi_7(x, t)$  for different values of parameters.

are obtained.

$$\psi_5(x, t) = \left\{ \frac{2}{5} \sqrt{\frac{-5\phi'(x)}{\lambda}} \left( -\frac{\lambda}{2} + \frac{\sqrt{\lambda^2 - 4\mu}}{2} \left( \frac{\text{Acosh}\left(\frac{1}{2z}\sqrt{\lambda^2 - 4\mu}\right) + \text{Bsinh}\left(\frac{1}{2z}\sqrt{\lambda^2 - 4\mu}\right)}{\text{Asinh}\left(\frac{1}{2z}\sqrt{\lambda^2 - 4\mu}\right) + \text{Bcosh}\left(\frac{1}{2z}\sqrt{\lambda^2 - 4\mu}\right)} \right) \right) \right. \\ \left. + 8\phi'(x) \sqrt{\frac{-5\phi'(x)}{\lambda}} \lambda \left( -\frac{\lambda}{2} + \frac{\sqrt{\lambda^2 - 4\mu}}{2} \left( \frac{\text{Acosh}\left(\frac{1}{2z}\sqrt{\lambda^2 - 4\mu}\right) + \text{Bsinh}\left(\frac{1}{2z}\sqrt{\lambda^2 - 4\mu}\right)}{\text{Asinh}\left(\frac{1}{2z}\sqrt{\lambda^2 - 4\mu}\right) + \text{Bcosh}\left(\frac{1}{2z}\sqrt{\lambda^2 - 4\mu}\right)} \right) \right)^2 \right. \\ \left. \left( (-60\lambda^2 + 80\mu)\phi'(x) + 25\lambda(p_1^4 + 1 - p_2p_1^2 + \phi(x)p_1^2) \right)^{-1} \right\} \times e^{ip}. \tag{13}$$

If  $A \rightarrow 0$  and  $B \neq 0$  and  $\lambda > 0$  then the solution Eq. (13) converges to

$$\psi_6(x, t) = \left\{ \frac{2}{5} \sqrt{\frac{-5\phi'(x)}{\lambda}} \left( -\frac{\lambda}{2} + \frac{\sqrt{\lambda^2 - 4\mu}}{2} \tanh\left(\frac{1}{2}z\sqrt{\lambda^2 - 4\mu}\right) \right) \right\} + \\ 8\phi'(x) \sqrt{\frac{-5\phi'(x)}{\lambda}} \lambda \left( -\frac{\lambda}{2} + \frac{\sqrt{\lambda^2 - 4\mu}}{2} \tanh\left(\frac{1}{2}z\sqrt{\lambda^2 - 4\mu}\right) \right)^2 \\ \left( (-60\lambda^2 + 80\mu)\phi'(x) + 25\lambda(p_1^4 + 1 - p_2p_1^2 + \phi(x)p_1^2) \right)^{-1} \right\} \times e^{ip}. \tag{14}$$

Also, if  $B \rightarrow 0$  and  $A \neq 0$  and  $\lambda > 0$  then the solution Eq. (13) converges to

$$\psi_7(x, t) = \left\{ \frac{2}{5} \sqrt{\frac{-5\phi'(x)}{\lambda}} \left( -\frac{\lambda}{2} + \frac{\sqrt{\lambda^2 - 4\mu}}{2} \tanh\left(\frac{1}{2}z\sqrt{\lambda^2 - 4\mu}\right) \right) \right\} + \\ 8\phi'(x) \sqrt{\frac{-5\phi'(x)}{\lambda}} \lambda \left( -\frac{\lambda}{2} + \frac{\sqrt{\lambda^2 - 4\mu}}{2} \coth\left(\frac{1}{2}z\sqrt{\lambda^2 - 4\mu}\right) \right)^2 \\ \left( (-60\lambda^2 + 80\mu)\phi'(x) + 25\lambda(p_1^4 + 1 - p_2p_1^2 + \phi(x)p_1^2) \right)^{-1} \right\} \times e^{ip}. \tag{15}$$

(ii) When  $\lambda^2 - 4\mu < 0$  thus following trigonometric wave solution are obtained.

$$\psi_8(x, t) = \left\{ \frac{2}{5} \sqrt{\frac{-5\phi'(x)}{\lambda}} \left( -\frac{\lambda}{2} + \sqrt{4\mu - \lambda^2} \left( \frac{-\text{Asin}\left(\frac{1}{2z}\sqrt{4\mu - \lambda^2}\right) + \text{Bcos}\left(\frac{1}{2z}\sqrt{4\mu - \lambda^2}\right)}{\text{Acos}\left(\frac{1}{2z}\sqrt{4\mu - \lambda^2}\right) + \text{Bsin}\left(\frac{1}{2z}\sqrt{4\mu - \lambda^2}\right)} \right) \right) \right. \\ \left. + 8\phi'(x) \sqrt{\frac{-5\phi'(x)}{\lambda}} \lambda \left( -\frac{\lambda}{2} + \sqrt{4\mu - \lambda^2} \left( \frac{-\text{Asin}\left(\frac{1}{2z}\sqrt{4\mu - \lambda^2}\right) + \text{Bcos}\left(\frac{1}{2z}\sqrt{4\mu - \lambda^2}\right)}{\text{Acos}\left(\frac{1}{2z}\sqrt{4\mu - \lambda^2}\right) + \text{Bsin}\left(\frac{1}{2z}\sqrt{4\mu - \lambda^2}\right)} \right) \right)^2 \right. \\ \left. + 1 - p_2p_1^2 + \phi(x)p_1^2 \right)^{-1} \right\} \times e^{ip}. \tag{16}$$

(iii) When  $\lambda^2 - 4\mu = 0$  thus plane solution are obtained.

$$\psi_9(x, t) = \left( -\frac{\lambda}{2} + \frac{B}{A + Bz} \right) \times e^{ip} \tag{17}$$

where

$$z = x - \left( \frac{-30p_1^2\lambda + 5\phi(x)\lambda + 5p_2\lambda - 2\phi'(x)}{10p_1\lambda} \right) t \tag{18}$$

The graph and its corresponding contour plot of the solutions  $\psi_6(x, t)$  are shown in Fig. 23, for the values of parameters  $\rho_1 = 0.7, \rho_2 = 0.008, \lambda = 0.05, \mu = 0.09, A = 10, B = 0.005, \phi(x) = \sin(x)$  and  $\nu = 0.2$

The graph and its corresponding contour plot of the solutions  $\psi_7(x, t)$  are shown in Fig. 24, for the values of parameters  $\rho_1 = 0.7, \rho_2 = 0.008, \lambda = 0.05, \mu = 0.09, A = 10, B = 0.005, \phi(x) = \sin(x)$  and  $\nu = 0.2$ .

For Case-3: The  $\beta_1$  is free parameter while along with  $\beta_0 = \beta_0, \beta_2 = 0, \beta_1 = -2 \frac{\phi(x) - 2\nu\rho_1 - 6\rho_1^2 + \rho_2}{\beta_0}$ ,

$$\nu = -\frac{1}{24\lambda\mu\rho_1} \left( 36\lambda^2\mu^2(\phi(x))^2 - 432\lambda \left( (-1/12\rho_1 - 1/6\rho_2 + \rho_1^2)\lambda - 1/36\rho_1 \right) \mu - 1/12\rho_1\beta_0 \right) \mu \phi(x) - 48\lambda\mu\rho_1\phi'(x)\beta_0 \\ + 1296 \left( (-1/12\rho_1 - 1/6\rho_2 + \rho_1^2)\lambda - 1/36\rho_1 \right) \mu - 1/12\rho_1\beta_0 \right)^2 \\ + 6\lambda\mu\phi(x) + \left( (-3\rho_1 + 6\rho_2 - 36\rho_1^2)\lambda - \rho_1 \right) \mu - 3\rho_1\beta_0 \right)^{1/2}.$$

The constraint condition holds for the following results, when vis

putting in it,

$$\lambda = -\frac{p_1(8\mu p_1^2 + 9\phi(x)\beta_0 + 2\phi'(x)\beta_0 + 24p_1^2\beta_0 + 9p_2\beta_0 + 3\mu\phi(x) + 3\mu p_2)}{3\mu(8p_1^3 + 24p_2^2 + 32p_1^4 + 24(\phi(x))^2 + 76p_2p_1^2 + 3\phi(x)p_1 + 48\phi(x)p_2 + 76\phi(x)p_1^2 + 3p_1p_2)}.$$

(i) When  $\lambda^2 - 4\mu > 0$  thus following hyperbolic solutions are obtained.

$$\psi_{10}(x, t) = \left\{ \beta_0 - \frac{2(\phi(x) - 2vp_1 - 6p_1^2 + p_2) \left( -\frac{\lambda}{2} + \frac{\sqrt{\lambda^2 - 4\mu}}{2} \left( \frac{\text{Acosh}(1/2z\sqrt{\lambda^2 - 4\mu}) + B\sinh(1/2z\sqrt{\lambda^2 - 4\mu})}{\text{Asinh}(1/2z\sqrt{\lambda^2 - 4\mu}) + B\cosh(1/2z\sqrt{\lambda^2 - 4\mu})} \right) \right)^2}{\beta_0} \right\} \times e^{ip}. \tag{19}$$

If  $A \rightarrow 0$  and  $B \neq 0$  and  $\lambda > 0$ , then Eq. (19) becomes

$$\psi_{11}(x, t) = \left\{ \beta_0 - \frac{2(\phi(x) - 2vp_1 - 6p_1^2 + p_2) \left( -\frac{\lambda}{2} + \frac{\sqrt{\lambda^2 - 4\mu}}{2} \left( \tanh\left(\frac{1}{2}z\sqrt{\lambda^2 - 4\mu}\right) \right) \right)^2}{\beta_0} \right\} \times e^{ip}. \tag{20}$$

Also, if  $B \rightarrow 0$  and  $A \neq 0$  and  $\lambda > 0$ , then Eq. (19) becomes

$$\psi_{12}(x, t) = \left\{ \beta_0 - \frac{2(\phi(x) - 2vp_1 - 6p_1^2 + p_2) \left( -\frac{\lambda}{2} + \frac{\sqrt{\lambda^2 - 4\mu}}{2} \left( \coth\left(\frac{1}{2}z\sqrt{\lambda^2 - 4\mu}\right) \right) \right)^2}{\beta_0} \right\} \times e^{ip}. \tag{21}$$

(ii) When  $\lambda^2 - 4\mu < 0$  thus following traveling wave solutions are obtained.

$$\psi_{13}(x, t) = \left\{ \beta_0 - \frac{2(\phi(x) - 2vp_1 - 6p_1^2 + p_2) \left( -\frac{\lambda}{2} + \sqrt{4\mu - \lambda^2} \left( \frac{-A\sin(1/2z\sqrt{4\mu - \lambda^2}) + B\cos(1/2z\sqrt{4\mu - \lambda^2})}{A\cos(1/2z\sqrt{4\mu - \lambda^2}) + B\sin(1/2z\sqrt{4\mu - \lambda^2})} \right) \right)^2}{\beta_0} \right\} \times e^{ip}. \tag{22}$$

(iii) When  $\lambda^2 - 4\mu = 0$  thus plane solutions are obtained.

$$\psi_{14}(x, t) = \left( -\frac{\lambda}{2} + \frac{B}{A + Bz} \right) \times e^{ip}. \tag{23}$$

where

$$z = x - \frac{1}{24\lambda\mu p_1} \left( 36\lambda^2\mu^2(\phi(x))^2 - 432\lambda \left( (-1/12p_1 - 1/6p_2 + p_1^2)\lambda - 1/36p_1 \right) \mu - 1/12p_1\beta_0 \right) \mu \phi(x) - 48\lambda\mu p_1 \phi'(x)\beta_0 + 1296 \left( (-1/12p_1 - 1/6p_2 + p_1^2)\lambda - 1/36p_1 \right) \mu - 1/12p_1\beta_0)^2 + 6\lambda\mu\phi(x) + \left( (-3p_1 + 6p_2 - 36p_1^2)\lambda - p_1 \right) \mu - 3p_1\beta_0)^{1/2} \times t.$$

#### Remark

Since the relation between  $\phi$  and  $\psi$  is the defined in Eq. (1), so using the definition of  $\Delta$  operator. The value of  $\phi$  can be obtained by twice integrating to  $|\psi|$  w.r.t space variable for above each solutions  $\psi_i$ , where  $i = 1, 2, 3, \dots, 25$ .

#### Conclusion

In this article, the diverse exact traveling wave solutions are obtained in the form of hyperbolic, trigonometric and plane wave families. The system under investigation is the Schrodinger–Poisson system which has applications in gravity’s role of quantum state. The solutions are constructed using two norms of integration MEDA and (G’/G)-expansion techniques. The constraints conditions for the existence of solutions also emerged during the derivation of solutions. The 3D graphs and their corresponding contour graphs are also being plotted by the proper choice of parameters. The results are new and helpful to compare with numerical results.

#### CRediT authorship contribution statement

**M. Younis:** Software, Validation, Supervision. **A.R. Seadawy:** Conceptualization, Methodology. **M.Z. Baber:** Writing - review & editing. **S. Husain:** Writing - review & editing. **M.S. Iqbal:** Data curation, Writing - original draft. **S.T. R. Rizvi:** Visualization, Investigation. **Dumitru Baleanu:** Writing - review & editing.

#### Declaration of Competing Interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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